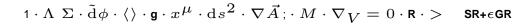
Special relativity and steps towards general relativity: ϵ GR

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= vector space (e.g. 4-momentum vectors)

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- = dual vector space (think: contour map, gradients)

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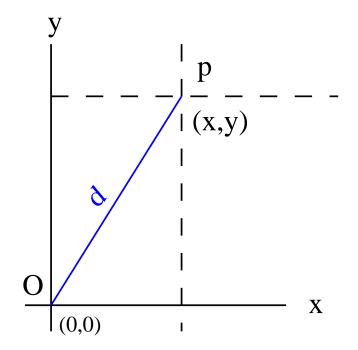
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duality in a basis of $T_{\bf x}M$ and a basis of $T_{\bf x}^*M$ usually defined using $\delta^{\mu}_{\ \nu}$

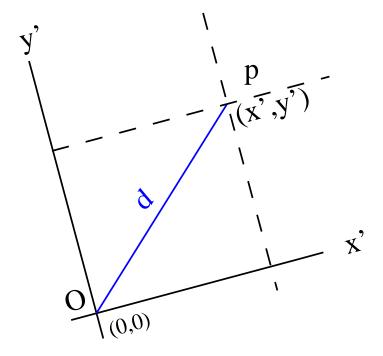
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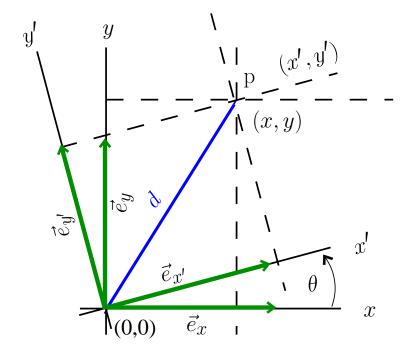
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- 5. metric \leftarrow Einstein field equations



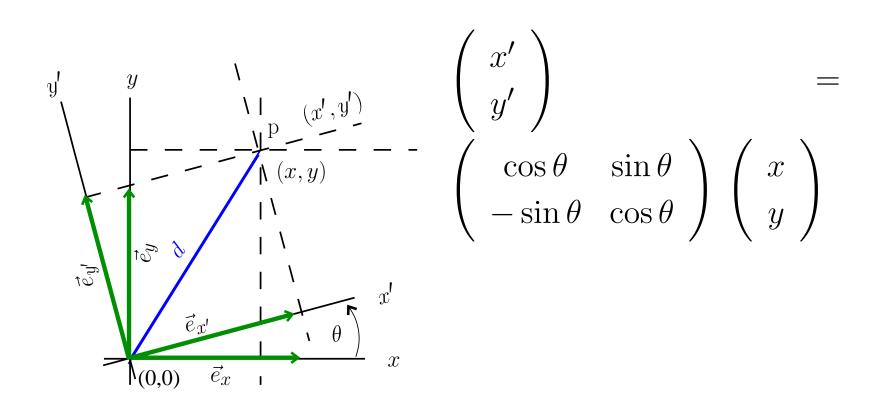
 $\mathbf{1}\cdot\Lambda\ \boldsymbol{\Sigma}\cdot\tilde{\mathbf{d}}\phi\cdot\langle\rangle\cdot\mathbf{g}\cdot\boldsymbol{x}^{\mu}\cdot\mathbf{d}\boldsymbol{s}^{2}\cdot\boldsymbol{\nabla}\vec{A}\,;\,\boldsymbol{M}\cdot\boldsymbol{\nabla}_{V}=\boldsymbol{0}\cdot\mathbf{R}\cdot\boldsymbol{>}\quad\mathbf{SR+}\epsilon\mathbf{GR}$

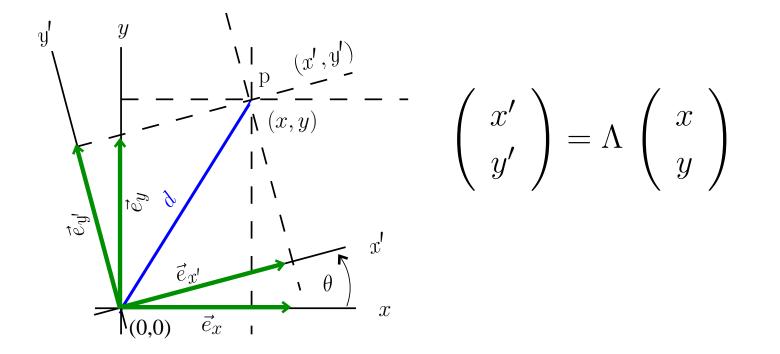


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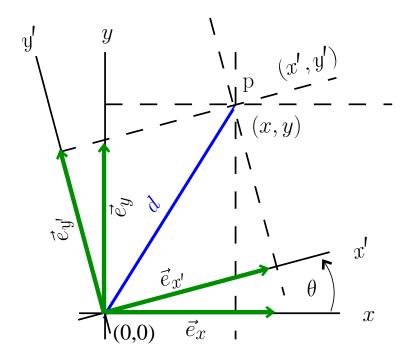


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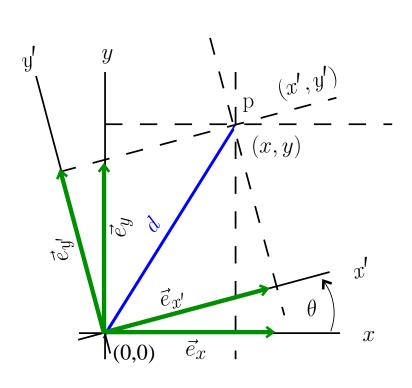


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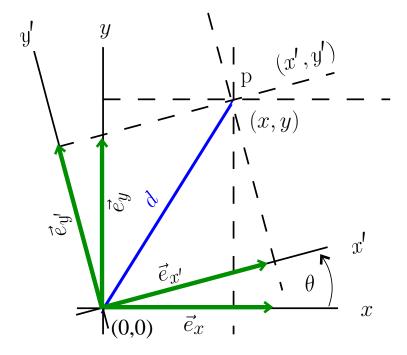
but
$$\begin{pmatrix} \cos \theta \\ \sin \theta \end{pmatrix} = \begin{pmatrix} \cos \theta & -\sin \theta \\ \sin \theta & \cos \theta \end{pmatrix} \begin{pmatrix} 1 \\ 0 \end{pmatrix}$$

 $1\cdot\Lambda\ \Sigma\cdot\tilde{\mathrm{d}}\phi\cdot\langle\rangle\cdot\mathbf{g}\cdot x^{\mu}\cdot\mathrm{d}s^{2}\cdot\nabla\vec{A}\,;\,\cdot\,M\cdot\nabla_{V}=0\cdot\mathbf{R}\cdot\rangle\quad\,\mathrm{SR+}\epsilon\mathrm{GR}$



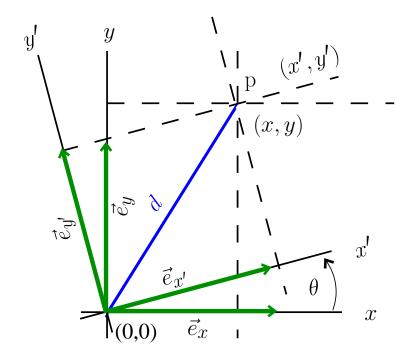
$$\begin{pmatrix} \cos \theta \\ \sin \theta \end{pmatrix} = \\ \begin{pmatrix} \cos \theta & -\sin \theta \\ \sin \theta & \cos \theta \end{pmatrix} \begin{pmatrix} 1 \\ 0 \end{pmatrix} + \\ \begin{pmatrix} \cos \theta & -\sin \theta \\ \sin \theta & \cos \theta \end{pmatrix} \begin{pmatrix} 0 \\ 1 \end{pmatrix}$$

 $\mathbf{1}\cdot\Lambda\ \boldsymbol{\Sigma}\cdot\tilde{\mathbf{d}}\boldsymbol{\phi}\cdot\boldsymbol{\langle}\boldsymbol{\rangle}\cdot\mathbf{g}\cdot\boldsymbol{x}^{\mu}\cdot\mathbf{d}\boldsymbol{s}^{2}\cdot\boldsymbol{\nabla}\vec{A}\,;\,\cdot\ \boldsymbol{M}\cdot\boldsymbol{\nabla}_{V}=\boldsymbol{0}\cdot\mathbf{R}\cdot\boldsymbol{\boldsymbol{\rangle}}\quad \mathbf{SR+}\epsilon\mathbf{GR}$



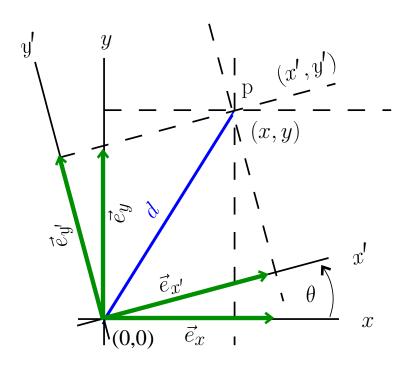
$$\vec{e}_{x'} = \begin{pmatrix} \cos \theta & -\sin \theta \\ \sin \theta & \cos \theta \end{pmatrix} \vec{e}_x + \\ \begin{pmatrix} \cos \theta & -\sin \theta \\ \sin \theta & \cos \theta \end{pmatrix} \vec{e}_y$$

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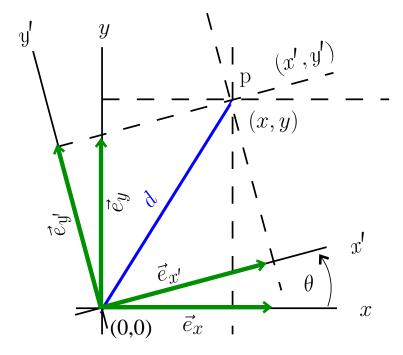


$$\vec{e}_{x'} = \Lambda^x_{x'} \, \vec{e}_x + \Lambda^y_{x'} \, \vec{e}_y$$

 $\mathbf{1}\cdot\Lambda\ \boldsymbol{\Sigma}\cdot\tilde{\mathbf{d}}\boldsymbol{\phi}\cdot\boldsymbol{\langle}\boldsymbol{\rangle}\cdot\mathbf{g}\cdot\boldsymbol{x}^{\mu}\cdot\mathbf{d}\boldsymbol{s}^{2}\cdot\boldsymbol{\nabla}\vec{A}\,;\ \boldsymbol{M}\cdot\boldsymbol{\nabla}_{V}=\boldsymbol{0}\cdot\mathbf{R}\cdot\boldsymbol{\boldsymbol{\rangle}}\quad \mathbf{SR+}\epsilon\mathbf{GR}$

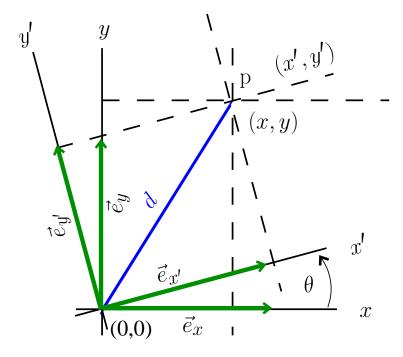


also: $\begin{pmatrix} -\sin\theta\\\cos\theta \end{pmatrix}$ $\begin{pmatrix} \cos\theta & -\sin\theta \\ \sin\theta & \cos\theta \end{pmatrix} \begin{pmatrix} 1 \\ 0 \end{pmatrix} +$ $\begin{pmatrix} \cos\theta & -\sin\theta \\ \sin\theta & \cos\theta \end{pmatrix} \begin{pmatrix} 0 \\ 1 \end{pmatrix}$



$$\vec{e}_{y'} = \begin{pmatrix} \cos\theta & -\sin\theta \\ \sin\theta & \cos\theta \end{pmatrix} \vec{e}_x + \\ \begin{pmatrix} \cos\theta & -\sin\theta \\ \sin\theta & \cos\theta \end{pmatrix} \vec{e}_y$$

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summary: $\vec{e}_{x'} = \Lambda_{x'}^x \vec{e}_x + \Lambda_{x'}^y \vec{e}_y,$ $\vec{e}_{y'} = \Lambda_{y'}^x \vec{e}_x + \Lambda_{y'}^y \vec{e}_y,$ where $\Lambda_{\beta'}^{\alpha} :=$ element of inverse of $\Lambda_{\beta}^{\alpha'},$ $\begin{pmatrix} x'\\ y' \end{pmatrix} = \Lambda \begin{pmatrix} x\\ y \end{pmatrix}$

 $1\cdot\Lambda \ \Sigma\cdot \widetilde{\mathrm{d}}\phi\cdot\langle\rangle\cdot\mathbf{g}\cdot x^{\mu}\cdot\mathrm{d}s^{2}\cdot\nablaec{A};\cdot M\cdot\nabla_{V}=0\cdot\mathbf{R}\cdot\rangle$ SR+ ϵ GR

$$\vec{e}_{x'} = \Lambda_{x'}^x \vec{e}_x + \Lambda_{x'}^y \vec{e}_y, \qquad \vec{p} \to \mathcal{O}' \begin{pmatrix} x' \\ y' \end{pmatrix} = \Lambda \begin{pmatrix} x \\ y \end{pmatrix}$$

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 $\vec{p} = \sum_{i} p^{i} \vec{e}_{i}$

$$\vec{e}_{x'} = \Lambda_{x'}^x \vec{e}_x + \Lambda_{x'}^y \vec{e}_y, \qquad \vec{p} \to \mathcal{O}' \left(\begin{array}{c} x' \\ y' \end{array} \right) = \Lambda \left(\begin{array}{c} x \\ y \end{array} \right)$$

 $\vec{p} = p^i \vec{e_i}$ (w:Einstein summation)

$$\vec{e}_{x'} = \Lambda_{x'}^x \vec{e}_x + \Lambda_{x'}^y \vec{e}_y, \qquad \vec{p} \to_{\mathcal{O}'} \begin{pmatrix} x' \\ y' \end{pmatrix} = \Lambda \begin{pmatrix} x \\ y \end{pmatrix}$$

 $\vec{p} = p^i \vec{e_i}$ (w:Einstein summation) Einstein summation:

• coordinates like r, θ, x, y : not a sum: $\Lambda^x_{y'} \vec{e}_x$

• repeated up-down coordinate indices like $i, j \in \{0, 1, 2\}$ or $\alpha, \beta, \gamma, \lambda, \mu, \nu \in \{0, 1, 2, 3\}$:

sum: $\Lambda_{j'}^i \vec{e_i} := \Lambda_{y'}^x \vec{e_x} + \Lambda_{y'}^y \vec{e_y}$ for a 2D manifold, coords x, y

$$\vec{e}_{x'} = \Lambda_{x'}^x \vec{e}_x + \Lambda_{x'}^y \vec{e}_y, \qquad \vec{p} \to \mathcal{O}' \left(\begin{array}{c} x' \\ y' \end{array} \right) = \Lambda \left(\begin{array}{c} x \\ y \end{array} \right)$$

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new basis vectors = sum of inverse $\Lambda \times$ old vectors

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 $\vec{p} = p^i \vec{e_i}$ (w:Einstein summation) new basis vectors = sum of inverse $\Lambda \times$ old vectors $\vec{e_{\mu'}} = \sum_{\nu} \Lambda^{\nu}_{\mu'} \vec{e_{\nu}}$

$$\vec{e}_{x'} = \Lambda_{x'}^x \vec{e}_x + \Lambda_{x'}^y \vec{e}_y, \qquad \vec{p} \to \mathcal{O}' \left(\begin{array}{c} x' \\ y' \end{array} \right) = \Lambda \left(\begin{array}{c} x \\ y \end{array} \right)$$

 $\vec{p} = p^i \vec{e}_i$ (w:Einstein summation) new basis vectors = sum of inverse $\Lambda \times$ old vectors $\vec{e}_{\mu'} = \Lambda^{\nu}_{\ \mu'} \vec{e}_{\nu}$

 $1 \cdot \Lambda \ \Sigma \cdot \tilde{\mathrm{d}} \phi \cdot \langle \rangle \cdot \mathbf{g} \cdot x^{\mu} \cdot \mathrm{d} s^2 \cdot \nabla \vec{A}; \cdot M \cdot \nabla_V = 0 \cdot \mathbf{R} \cdot \rangle \quad SR+\epsilon GR$

$$\vec{e}_{x'} = \Lambda_{x'}^x \vec{e}_x + \Lambda_{x'}^y \vec{e}_y, \qquad \vec{p} \to \mathcal{O}' \left(\begin{array}{c} x' \\ y' \end{array} \right) = \Lambda \left(\begin{array}{c} x \\ y \end{array} \right)$$

 $\vec{p} = p^i \vec{e_i}$ (w:Einstein summation) new basis vectors = sum of inverse $\Lambda \times$ old vectors new coords of vector $\vec{p} = \Lambda \times$ old coords of same vector \vec{p}

$$\vec{e}_{x'} = \Lambda_{x'}^x \vec{e}_x + \Lambda_{x'}^y \vec{e}_y, \qquad \vec{p} \to_{\mathcal{O}'} \begin{pmatrix} x' \\ y' \end{pmatrix} = \Lambda \begin{pmatrix} x \\ y \end{pmatrix}$$

 $\vec{p} = p^i \vec{e_i}$ (w:Einstein summation)

new basis vectors = sum of inverse $\Lambda \times$ old vectors

vector invariance requires contravariance of its coords "contra" = inverse of change of basis vectors

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 $\vec{p} = p^i \vec{e_i}$ (w:Einstein summation)

new basis vectors = sum of inverse $\Lambda \times$ old vectors

vector invariance requires contravariance of its coords "contra" = inverse of change of basis vectors

- \vec{p} is invariant: no dependence on coords
- \vec{p} is contravariant: p^i change inversely to $\vec{e_i}$

 $1 \cdot \Lambda \ \Sigma \cdot \tilde{\mathrm{d}} \phi \cdot \langle \rangle \cdot \mathbf{g} \cdot x^{\mu} \cdot \mathrm{d} s^2 \cdot \nabla \vec{A}; \quad M \cdot \nabla_V = 0 \cdot \mathbf{R} \cdot \rangle \quad \mathsf{SR}_{\mathsf{f}} \epsilon \mathsf{GR}$

GR: coord. transf.: 1-forms

$$\phi = \text{scalar field} = \phi(x, y) \equiv \phi(x', y')$$

write $\phi_{,x} := \frac{\partial \phi}{\partial x} =: (\tilde{d}\phi)_x$

 $\mathbf{1}\cdot\Lambda\ \boldsymbol{\Sigma}\cdot\tilde{\mathbf{d}}\boldsymbol{\phi}\cdot\boldsymbol{\langle}\boldsymbol{\rangle}\cdot\mathbf{g}\cdot\boldsymbol{x}^{\mu}\cdot\mathbf{d}\boldsymbol{s}^{2}\cdot\boldsymbol{\nabla}\vec{A}\,;\,\boldsymbol{M}\cdot\boldsymbol{\nabla}_{V}=\boldsymbol{0}\cdot\mathbf{R}\cdot\boldsymbol{\boldsymbol{\rangle}}\quad \mathbf{SR+}\epsilon\mathbf{GR}$

GR: coord. transf.: 1-forms

 $\phi = \text{scalar field} = \phi(x, y) \equiv \phi(x', y')$ write $\phi_{,x} := \frac{\partial \phi}{\partial x} =: (\tilde{d}\phi)_x$ What is the relation between $(\phi_{,x'}, \phi_{,y'})$ and $(\phi_{,x}, \phi_{,y})$?

 $\phi = \text{scalar field} = \phi(x, y) \equiv \phi(x', y')$ write $\phi_{,x} := \frac{\partial \phi}{\partial x} =: (\tilde{d}\phi)_x$

 ϕ depends either on x and y, or on x' and y'

 $\Rightarrow \frac{\partial \phi}{\partial x'} = \frac{\partial \phi}{\partial x} \frac{\partial x}{\partial x'} + \frac{\partial \phi}{\partial y} \frac{\partial y}{\partial x'}$

 $\phi = \text{scalar field} = \phi(x, y) \equiv \phi(x', y')$ write $\phi_{,x} := \frac{\partial \phi}{\partial x} =: (\tilde{d}\phi)_x$ ϕ depends either on x and y, or on x' and y' $\Rightarrow \phi_{,x'} = \phi_{,x} x_{,x'} + \phi_{,y} y_{,x'}$

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 $(\phi_{,x'},\phi_{,y'}) =$

 $1\cdot\Lambda\ \Sigma\cdot\tilde{\mathrm{d}}\phi\cdot\langle\rangle\cdot\mathbf{g}\cdot x^{\mu}\cdot\mathrm{d}s^{2}\cdot\nabla\vec{A}\,;\,\cdot\,M\cdot\nabla_{V}=0\cdot\mathbf{R}\cdot\rangle\quad\,\mathrm{SR+}\epsilon\mathrm{GR}$

$$\begin{split} \phi &= \text{scalar field} = \phi(x, y) \equiv \phi(x', y') \\ \text{write } \phi_{,x} &:= \frac{\partial \phi}{\partial x} =: (\tilde{d}\phi)_x \\ \phi \text{ depends either on } x \text{ and } y, \text{ or on } x' \text{ and } y' \\ \Rightarrow \phi_{,x'} &= \phi_{,x} x_{,x'} + \phi_{,y} y_{,x'} \\ (\phi_{,x'}, \phi_{,y'}) &= (\phi_{,x} x_{,x'} + \phi_{,y} y_{,x'}, \phi_{,x} x_{,y'} + \phi_{,y} y_{,y'}) \end{split}$$

 $1 \cdot \Lambda \ \Sigma \cdot \tilde{\mathrm{d}} \phi \cdot \langle \rangle \cdot \mathbf{g} \cdot x^{\mu} \cdot \mathrm{d} s^2 \cdot \nabla \vec{A}; \quad M \cdot \nabla_V = 0 \cdot \mathbf{R} \cdot \rangle \quad \mathsf{SR}_{\mathsf{f}} \epsilon \mathsf{GR}$

y'

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 $\mathbf{1}\cdot\Lambda\ \boldsymbol{\Sigma}\cdot\tilde{\mathbf{d}}\phi\cdot\langle\rangle\cdot\mathbf{g}\cdot\boldsymbol{x}^{\mu}\cdot\mathbf{d}\boldsymbol{s}^{2}\cdot\boldsymbol{\nabla}\vec{A}\,;\,\cdot\,\boldsymbol{M}\cdot\boldsymbol{\nabla}_{V}=\boldsymbol{0}\cdot\mathbf{R}\cdot\boldsymbol{>}\quad\mathbf{SR+}\epsilon\mathbf{GR}$

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 $(\phi_{,x'}, \phi_{,y'}) = (\phi_{,x}, \phi_{,y}) \begin{pmatrix} x_{,x'} & x_{,y'} \\ y_{,x'} & y_{,y'} \end{pmatrix}$
 $\begin{pmatrix} x \\ y \end{pmatrix} = \begin{pmatrix} \cos \theta & -\sin \theta \\ \sin \theta & \cos \theta \end{pmatrix} \begin{pmatrix} x' \\ y' \end{pmatrix}$ (example: rotation)
 $x_{,x'} = \frac{\partial x}{\partial x'} = \cos \theta$
 $x_{,y'} = \frac{\partial x}{\partial y'} = -\sin \theta \dots$

 $\mathbf{1}\cdot\Lambda\ \boldsymbol{\Sigma}\cdot\tilde{\mathbf{d}}\phi\cdot\langle\rangle\cdot\mathbf{g}\cdot\boldsymbol{x}^{\mu}\cdot\mathbf{d}\boldsymbol{s}^{2}\cdot\boldsymbol{\nabla}\vec{A}\,;\,\boldsymbol{M}\cdot\boldsymbol{\nabla}_{V}=\boldsymbol{0}\cdot\mathbf{R}\cdot\boldsymbol{>}\quad\mathbf{SR+}\epsilon\mathbf{GR}$

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 $\begin{pmatrix} x \\ y \end{pmatrix} = \begin{pmatrix} x_{,x'} & x_{,y'} \\ y_{,x'} & y_{,y'} \end{pmatrix} \begin{pmatrix} x' \\ y' \end{pmatrix}$ (general)

 $\mathbf{1}\cdot\Lambda\ \boldsymbol{\Sigma}\cdot\tilde{\mathbf{d}}\phi\cdot\langle\rangle\cdot\mathbf{g}\cdot\boldsymbol{x}^{\mu}\cdot\mathbf{d}\boldsymbol{s}^{2}\cdot\boldsymbol{\nabla}\vec{A}\,;\,\cdot\,\boldsymbol{M}\cdot\boldsymbol{\nabla}_{V}=\boldsymbol{0}\cdot\mathbf{R}\cdot\boldsymbol{>}\quad\mathbf{SR+}\epsilon\mathbf{GR}$

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 $\begin{pmatrix} x \\ y \end{pmatrix} = \Lambda^{-1} \begin{pmatrix} x' \\ y' \end{pmatrix}$ (general)

 $\mathbf{1}\cdot\Lambda\ \boldsymbol{\Sigma}\cdot\tilde{\mathbf{d}}\boldsymbol{\phi}\cdot\boldsymbol{\langle}\boldsymbol{\rangle}\cdot\mathbf{g}\cdot\boldsymbol{x}^{\boldsymbol{\mu}}\cdot\mathbf{d}\boldsymbol{s}^{2}\cdot\boldsymbol{\nabla}\vec{A}\,;\,\cdot\ \boldsymbol{M}\cdot\boldsymbol{\nabla}_{V}=\boldsymbol{0}\cdot\mathbf{R}\cdot\boldsymbol{\boldsymbol{\rangle}}\quad \mathbf{SR+}\epsilon\mathbf{GR}$

y'

$$\begin{split} \phi &= \text{scalar field} = \phi(x, y) \equiv \phi(x', y') \\ \text{write } \phi_{,x} &:= \frac{\partial \phi}{\partial x} =: (\tilde{d}\phi)_{x} \\ \phi \text{ depends either on } x \text{ and } y, \text{ or on } x' \text{ and} \\ \Rightarrow \phi_{,x'} &= \phi_{,x} x_{,x'} + \phi_{,y} y_{,x'} \\ (\phi_{,x'}, \phi_{,y'}) &= (\phi_{,x}, \phi_{,y}) \begin{pmatrix} x_{,x'} & x_{,y'} \\ y_{,x'} & y_{,y'} \end{pmatrix} \\ \begin{pmatrix} x \\ y \end{pmatrix} &= \Lambda^{-1} \begin{pmatrix} x' \\ y' \end{pmatrix} \text{ (general)} \\ \Rightarrow (\phi_{,x'}, \phi_{,y'}) &= (\phi_{,x}, \phi_{,y}) \Lambda^{-1} \end{split}$$

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basis vectors of different bases: $\vec{e}_{\mu'} = \Lambda^{\nu}_{\ \mu'} \vec{e}_{\nu}$ same vector: $(\vec{p})^{\mu'} = \Lambda^{\mu'}_{\ \nu} (\vec{p})^{\nu}$

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w:Covariance and contravariance of vectors

GR tensors: two different scalar products

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$$\langle \vec{p}, \tilde{q} \rangle = \sum_{\mu} p^{\mu} q_{\mu}$$

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can be called I with components δ^{μ}_{ν} in a coordinate basis

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think: vector \rightarrow column vector 1-form \rightarrow row vector

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 $\langle , \rangle = (1,1)$ -tensor = "row-column" matrix I with $I^{\mu}_{\nu} = \delta^{\mu}_{\nu}$

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GR tensors: two different scalar products

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ordinary linear algebra: column vectors, row vectors, matrices

GR tensors: two different scalar products

(m, n)-tensor algebra: m column n row m + n-arrays

GR tensors: two different scalar products

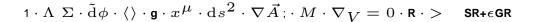
(m, n)-tensor algebra: m column n row m + n-arrays

e.g.: (0,2)-tensor: metric $g_{\mu\nu}$

GR tensors: two different scalar products

(m, n)-tensor algebra: m column n row m + n-arrays

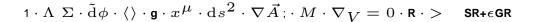
using \langle , \rangle , (1,0)-tensor = vector = function of 1-forms



GR tensors: two different scalar products

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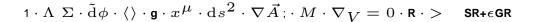
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GR tensors: two different scalar products

(m, n)-tensor algebra: m column n row m + n-arrays

(m, n)-tensor = function of m 1-forms and n vectors



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(m, n)-tensor algebra: m column n row m + n-arrays

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V = space of vectors $\vec{p} = p^{\mu} \vec{e}_{\mu}$

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loosely speaking, the second \otimes means "function of two vectors" (or 1-forms, or a vector and a 1-form) in *that particular left-to-right order*

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order of $V^* \otimes V^* = 2$

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warning: the "rank" of tensors has two different meanings: w:Tensor_(intrinsic_definition)#Tensor_rank

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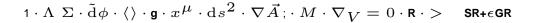
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dimension of $V^* \otimes V^* = 16$ (for V = spacetime)

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$$= \left[\begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} \begin{pmatrix} A^x \\ A^y \end{pmatrix} \right]^{\mathrm{T}} \begin{pmatrix} B^x \\ B^y \end{pmatrix}$$

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g can be applied to basis vectors \vec{e}_{μ}

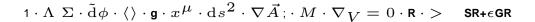
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 $\mathbf{g}(\vec{e}_r, \vec{e}_r) = g_{rr} \times 1 \times 1 + g_{\theta\theta} \times 0 \times 0$ by duality through scalar product \langle , \rangle

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 $\mathbf{g}(\vec{e_r}, \vec{e_r}) = g_{rr}$ self-consistent definition

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lower an index: $g_{\mu\nu}A^{\mu} = A_{\nu}$

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a coordinate, e.g. x^0 or x^1 is a scalar field on the 4-manifold

a coordinate system x^{μ} = set of four scalar fields on the 4-manifold

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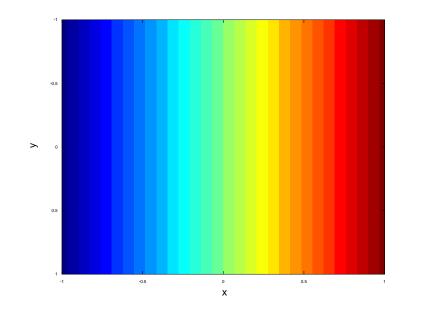
(Bertschinger writes $x_{\mathbf{x}}^{\mu}$ to show dependence on position \mathbf{x} in manifold \neq vector space)

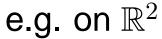
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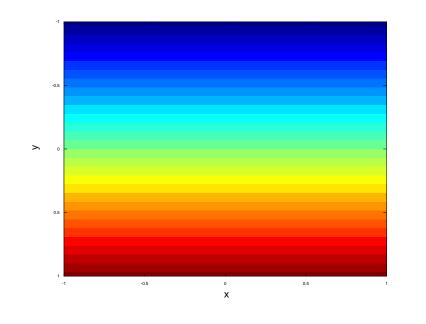




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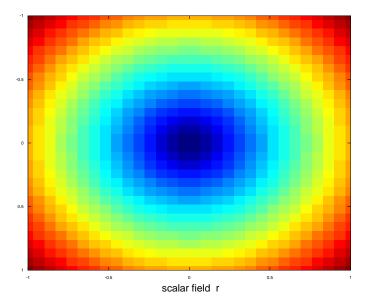
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e.g. on \mathbb{R}^2

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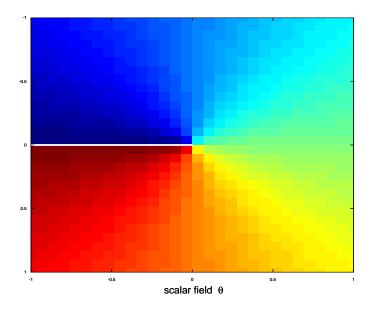
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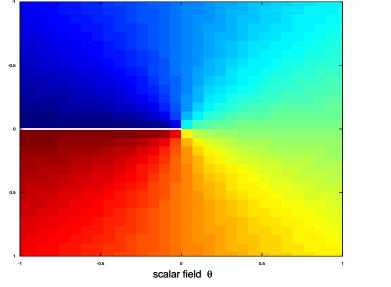
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coordinate singularity \neq singularity in manifold

coordinate basis: \vec{e}_{μ} , \tilde{e}^{ν} chosen so that:

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where $\tilde{d} = \tilde{e}^{\mu} \partial_{\mu}$ in a coordinate basis

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 $df = \langle \tilde{d}f, d\vec{x} \rangle \text{ for any scalar field } f \text{ coordinate-free}$ where $\tilde{d} = \tilde{e}^{\mu} \partial_{\mu}$ in a coordinate basis check: $df = \langle \tilde{d}f, d\vec{x} \rangle$ $= \langle \tilde{e}^{\mu} \partial_{\mu} f, dx^{\nu} \vec{e_{\nu}} \rangle$

 $1 \cdot \Lambda \ \Sigma \cdot \tilde{\mathrm{d}} \phi \cdot \langle \rangle \cdot \mathbf{g} \cdot x^{\mu} \cdot \mathrm{d} s^2 \cdot \nabla \vec{A}; \quad M \cdot \nabla_V = 0 \cdot \mathbf{R} \cdot > \quad \mathsf{SR}_{\mathbf{f}} \epsilon \mathsf{GR}$

coordinate basis: \vec{e}_{μ} , \tilde{e}^{ν} chosen so that:

$$\mathrm{d}\vec{x} = \mathrm{d}x^{\mu}\vec{e_{\mu}}$$
 and

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 $1 \cdot \Lambda \ \Sigma \cdot \tilde{\mathrm{d}} \phi \cdot \langle \rangle \cdot \mathbf{g} \cdot x^{\mu} \cdot \mathrm{d} s^2 \cdot \nabla \vec{A}; \quad M \cdot \nabla_V = 0 \cdot \mathbf{R} \cdot \rangle \quad \mathsf{SR}_{\mathsf{f}} \epsilon \mathsf{GR}$

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 $1 \cdot \Lambda \ \Sigma \cdot \tilde{\mathrm{d}} \phi \cdot \langle \rangle \cdot \mathbf{g} \cdot x^{\mu} \cdot \mathrm{d} s^2 \cdot \nabla \vec{A}; \quad M \cdot \nabla_V = 0 \cdot \mathbf{R} \cdot \rangle \quad \mathsf{SR}_{\mathsf{f}} \epsilon \mathsf{GR}$

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 $1 \cdot \Lambda \ \Sigma \cdot \tilde{\mathrm{d}} \phi \cdot \langle \rangle \cdot \mathbf{g} \cdot x^{\mu} \cdot \mathrm{d} s^2 \cdot \nabla \vec{A}; \quad M \cdot \nabla_V = 0 \cdot \mathbf{R} \cdot \rangle \quad SR+\epsilon GR$

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 $1 \cdot \Lambda \ \Sigma \cdot \tilde{\mathrm{d}} \phi \cdot \langle \rangle \cdot \mathbf{g} \cdot x^{\mu} \cdot \mathrm{d} s^2 \cdot \nabla \vec{A}; \quad M \cdot \nabla_V = 0 \cdot \mathbf{R} \cdot \rangle \quad \mathrm{SR}_{+} \epsilon \mathrm{GR}$

coordinate basis: \vec{e}_{μ} , \tilde{e}^{ν} chosen so that:

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 $1 \cdot \Lambda \ \Sigma \cdot \tilde{\mathrm{d}} \phi \cdot \langle \rangle \cdot \mathbf{g} \cdot x^{\mu} \cdot \mathrm{d} s^2 \cdot \nabla \vec{A}; \quad M \cdot \nabla_V = 0 \cdot \mathbf{R} \cdot \rangle \quad SR+\epsilon GR$

coordinate basis: \vec{e}_{μ} , \tilde{e}^{ν} chosen so that:

$$\mathrm{d}\vec{x} = \mathrm{d}x^{\mu}\vec{e_{\mu}}$$
 and

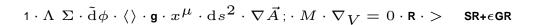
 $df = \langle \tilde{d}f, d\vec{x} \rangle \text{ for any scalar field } f \text{ coordinate-free}$ where $\tilde{d} = \tilde{e}^{\mu} \partial_{\mu}$ in a coordinate basis check: $df = \langle \tilde{d}f, d\vec{x} \rangle$ $= \langle \tilde{e}^{\mu} \partial_{\mu} f, dx^{\nu} \vec{e}_{\nu} \rangle$ $= (\partial_{\mu} f) dx^{\nu} \langle \tilde{e}^{\mu}, \vec{e}_{\nu} \rangle$ since scalars commute i.e. $df = (\partial_{\mu} f) dx^{\mu}$ check: $\tilde{d}x^{\mu} = \tilde{e}^{\nu} \partial_{\nu} x^{\mu}$ $= \tilde{e}^{\mu}$

 $1 \cdot \Lambda \Sigma \cdot \tilde{\mathrm{d}} \phi \cdot \langle \rangle \cdot \mathbf{g} \cdot x^{\mu} \cdot \mathrm{d} s^2 \cdot \nabla \vec{A}; \quad M \cdot \nabla_V = 0 \cdot \mathbf{R} \cdot \rangle \quad SR+\epsilon GR$

we now have

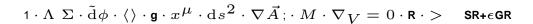
we now have

 $\mathrm{d}s^2 := |\mathrm{d}\vec{x}|^2$



we now have

 $\mathrm{d}s^2 := |\mathrm{d}\vec{x}|^2 = \mathbf{g}(\mathrm{d}\vec{x}, \mathrm{d}\vec{x})$



we now have

 $ds^2 := |d\vec{x}|^2 = \mathbf{g}(d\vec{x}, d\vec{x}) = d\vec{x} \cdot d\vec{x}$ coordinate-free

 $1\cdot\Lambda\ \Sigma\cdot\tilde{\mathrm{d}}\phi\cdot\langle\rangle\cdot\mathbf{g}\cdot x^{\mu}\cdot\mathrm{d}s^{2}\cdot\nabla\vec{A}\,;\,\cdot\,M\cdot\nabla_{V}=0\cdot\mathbf{R}\cdot\rangle\quad\mathrm{SR+}\epsilon\mathrm{GR}$

we now have

 $ds^2 := |d\vec{x}|^2 = \mathbf{g}(d\vec{x}, d\vec{x}) = d\vec{x} \cdot d\vec{x}$ coordinate-free

 $ds^2 = g_{\mu\nu} dx^{\mu} x^{\nu}$ if x^{μ} are a coordinate basis

 $g_{r\theta}$ and g_{xy}

 $\mathrm{d}s^2 = \mathrm{d}x^2 + \mathrm{d}y^2 = \mathrm{d}r^2 + r^2\mathrm{d}\theta^2$

 $1\cdot\Lambda\ \Sigma\cdot\tilde{\mathrm{d}}\phi\cdot\langle\rangle\cdot\mathbf{g}\cdot x^{\mu}\cdot\mathrm{d}s^{2}\cdot\nabla\vec{A}\,;\,M\cdot\nabla_{V}=0\cdot\mathbf{R}\cdot\rangle\quad \mathrm{SR+}\epsilon\mathrm{GR}$

 $g_{r\theta}$ and g_{xy}

- $\mathrm{d}s^2 = \mathrm{d}x^2 + \mathrm{d}y^2 = \mathrm{d}r^2 + r^2\mathrm{d}\theta^2$
- $\vec{e}_x \cdot \vec{e}_x = 1 = \vec{e}_y \cdot \vec{e}_y$, others zero

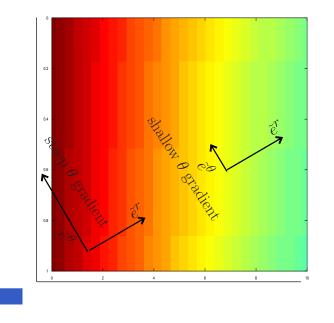
 $g_{r\theta}$ and g_{xy} $\mathrm{d}s^2 = \mathrm{d}x^2 + \mathrm{d}y^2 = \mathrm{d}r^2 + r^2\mathrm{d}\theta^2$ $\vec{e}_x \cdot \vec{e}_x = 1 = \vec{e}_y \cdot \vec{e}_y$, others zero $\vec{e_r} \cdot \vec{e_r} = 1, \ \vec{e_\theta} \cdot \vec{e_\theta} = r^2 \neq 1$ yB, x

 $1 \cdot \Lambda \ \Sigma \cdot \tilde{\mathrm{d}} \phi \cdot \langle \rangle \cdot \mathbf{g} \cdot x^{\mu} \cdot \mathrm{d} s^2 \cdot \nabla \vec{A}; \quad M \cdot \nabla_V = 0 \cdot \mathbf{R} \cdot \rangle \quad \mathrm{SR}_{\mathbf{f}} \epsilon \mathrm{GR}$

 $g_{r\theta} \text{ and } g_{xy}$ $ds^{2} = dx^{2} + dy^{2} = dr^{2} + r^{2}d\theta^{2}$ $\vec{e}_{x} \cdot \vec{e}_{x} = 1 = \vec{e}_{y} \cdot \vec{e}_{y}, \text{ others zero}$ $\vec{e}_{r} \cdot \vec{e}_{r} = 1, \vec{e}_{\theta} \cdot \vec{e}_{\theta} = r^{2} \neq 1$ $g^{\mu\alpha}g_{\alpha\nu} = \delta^{\mu}_{\nu} \Rightarrow g^{xx} = 1 = g^{yy}, g^{xy} = 0 = g^{yx}$

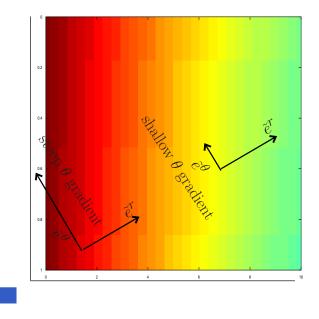
 $g_{r\theta} \text{ and } g_{xy}$ $ds^{2} = dx^{2} + dy^{2} = dr^{2} + r^{2}d\theta^{2}$ $\vec{e}_{x} \cdot \vec{e}_{x} = 1 = \vec{e}_{y} \cdot \vec{e}_{y}, \text{ others zero}$ $\vec{e}_{r} \cdot \vec{e}_{r} = 1, \vec{e}_{\theta} \cdot \vec{e}_{\theta} = r^{2} \neq 1$ $g^{\mu\alpha}g_{\alpha\nu} = \delta^{\mu}_{\nu} \Rightarrow g^{xx} = 1 = g^{yy}, g^{xy} = 0 = g^{yx}$ but $g^{rr} = 1 \neq g^{\theta\theta} = r^{-2}, g^{r\theta} = 0 = g^{\theta r}$

 $g_{r\theta} \text{ and } g_{xy}$ $ds^{2} = dx^{2} + dy^{2} = dr^{2} + r^{2}d\theta^{2}$ $\vec{e}_{x} \cdot \vec{e}_{x} = 1 = \vec{e}_{y} \cdot \vec{e}_{y}, \text{ others zero}$ $\vec{e}_{r} \cdot \vec{e}_{r} = 1, \vec{e}_{\theta} \cdot \vec{e}_{\theta} = r^{2} \neq 1$ $g^{\mu\alpha}g_{\alpha\nu} = \delta^{\mu}_{\nu} \Rightarrow g^{xx} = 1 = g^{yy}, g^{xy} = 0 = g^{yx}$ but $g^{rr} = 1 \neq g^{\theta\theta} = r^{-2}, g^{r\theta} = 0 = g^{\theta r}$



 $1 \cdot \Lambda \ \Sigma \cdot \tilde{\mathrm{d}} \phi \cdot \langle \rangle \cdot \mathbf{g} \cdot x^{\mu} \cdot \mathrm{d} s^2 \cdot \nabla \vec{A}; \quad M \cdot \nabla_V = 0 \cdot \mathbf{R} \cdot \rangle \quad \mathrm{SR}_{\mathbf{f}} \epsilon \mathrm{GR}$

 $g_{r\theta} \text{ and } g_{xy}$ $ds^{2} = dx^{2} + dy^{2} = dr^{2} + r^{2}d\theta^{2}$ $\vec{e}_{x} \cdot \vec{e}_{x} = 1 = \vec{e}_{y} \cdot \vec{e}_{y}, \text{ others zero}$ $\vec{e}_{r} \cdot \vec{e}_{r} = 1, \vec{e}_{\theta} \cdot \vec{e}_{\theta} = r^{2} \neq 1$ $g^{\mu\alpha}g_{\alpha\nu} = \delta^{\mu}_{\nu} \Rightarrow g^{xx} = 1 = g^{yy}, g^{xy} = 0 = g^{yx}$ but $g^{rr} = 1 \neq g^{\theta\theta} = r^{-2}, g^{r\theta} = 0 = g^{\theta r}$



SO
$$ilde{e}^r \cdot ilde{e}^r = 1$$
, $ilde{e}^ heta \cdot ilde{e}^ heta = r^{-2}
eq 1$

 $1 \cdot \Lambda \ \Sigma \cdot \tilde{\mathrm{d}} \phi \cdot \langle \rangle \cdot \mathbf{g} \cdot x^{\mu} \cdot \mathrm{d} s^2 \cdot \nabla \vec{A}; \quad M \cdot \nabla_V = 0 \cdot \mathbf{R} \cdot \rangle \quad \mathrm{SR}_{\mathbf{f}} \epsilon \mathrm{GR}$

gradient of scalar field: $\tilde{d}\phi \equiv \widetilde{\nabla}\phi$

 $1 \cdot \Lambda \ \Sigma \cdot \tilde{\mathrm{d}} \phi \cdot \langle \rangle \cdot \mathbf{g} \cdot x^{\mu} \cdot \mathrm{d} s^2 \cdot \nabla \vec{A}; \quad M \cdot \nabla_V = 0 \cdot \mathbf{R} \cdot > \quad \mathsf{SR+}\epsilon \mathsf{GR}$

what is gradient of vector field $\widetilde{\nabla} \vec{A}$?

GR: gradient of a vector: $\nabla \vec{A}$ $\tilde{\nabla} \vec{A} = \tilde{\nabla} (A^{\nu} \vec{e}_{\nu})$

 $1 \cdot \Lambda \ \Sigma \cdot \tilde{\mathrm{d}} \phi \cdot \langle \rangle \cdot \mathbf{g} \cdot x^{\mu} \cdot \mathrm{d} s^2 \cdot \nabla \vec{A}; \quad M \cdot \nabla_V = 0 \cdot \mathbf{R} \cdot \rangle \quad \mathrm{SR}_{+} \epsilon \mathrm{GR}$

 $\widetilde{\nabla} \vec{A} = \widetilde{\nabla} (A^{\nu} \vec{e}_{\nu})$ $= \tilde{e}^{\mu} \partial_{\mu} (A^{\nu} \vec{e}_{\nu})$

- $\widetilde{\nabla} \vec{A} = \widetilde{\nabla} (A^{\nu} \vec{e}_{\nu})$
- $= \tilde{e}^{\mu} \partial_{\mu} (A^{\nu} \vec{e}_{\nu})$

 $= \tilde{e}^{\mu} \otimes [(\partial_{\mu}A^{\nu})\vec{e}_{\nu} + A^{\nu}\partial_{\mu}\vec{e}_{\nu}]$ by product rule and linearity

- $\widetilde{\nabla} \vec{A} = \widetilde{\nabla} (A^{\nu} \vec{e}_{\nu})$
- $= \tilde{e}^{\mu} \partial_{\mu} (A^{\nu} \vec{e}_{\nu})$
- $= \partial_{\mu}A^{\nu}\tilde{e}^{\mu}\otimes\vec{e}_{\nu} + A^{\nu}\tilde{e}^{\mu}\otimes\partial_{\mu}\vec{e}_{\nu}$

- $\widetilde{\nabla} \vec{A} = \widetilde{\nabla} (A^{\nu} \vec{e}_{\nu})$
- $= \tilde{e}^{\mu} \partial_{\mu} (A^{\nu} \vec{e}_{\nu})$
- $= \partial_{\mu}A^{\nu}\tilde{e}^{\mu}\otimes\vec{e}_{\nu} + A^{\nu}\tilde{e}^{\mu}\otimes\partial_{\mu}\vec{e}_{\nu}$

give a name to the second part: it must be a linear combination of basis vectors \vec{e}_{λ}

- $\widetilde{\nabla} \vec{A} = \widetilde{\nabla} (A^{\nu} \vec{e}_{\nu})$
- $= \tilde{e}^{\mu} \partial_{\mu} (A^{\nu} \vec{e}_{\nu})$
- $= \partial_{\mu}A^{\nu}\tilde{e}^{\mu}\otimes\vec{e}_{\nu} + A^{\nu}\tilde{e}^{\mu}\otimes\partial_{\mu}\vec{e}_{\nu}$

define $\Gamma^{\lambda}_{\nu\mu}\vec{e}_{\lambda} := \partial_{\mu}\vec{e}_{\nu}$ Christoffel symbols of second kind (symmetric defn)

- $\widetilde{\nabla} \vec{A} = \widetilde{\nabla} (A^{\nu} \vec{e}_{\nu})$
- $= \tilde{e}^{\mu} \partial_{\mu} (A^{\nu} \vec{e}_{\nu})$
- $= \partial_{\mu}A^{\nu}\tilde{e}^{\mu}\otimes\vec{e}_{\nu} + A^{\nu}\tilde{e}^{\mu}\otimes\partial_{\mu}\vec{e}_{\nu}$
- define $\Gamma^{\lambda}_{\nu\mu}\vec{e}_{\lambda} := \partial_{\mu}\vec{e}_{\nu}$ Christoffel symbols of second kind (symmetric defn)
- so $\widetilde{\nabla} \vec{A} = \partial_{\mu} A^{\nu} \tilde{e}^{\mu} \otimes \vec{e}_{\nu} + A^{\nu} \tilde{e}^{\mu} \otimes \Gamma^{\lambda}_{\ \nu\mu} \vec{e}_{\lambda}$

- $\widetilde{\nabla} \vec{A} = \widetilde{\nabla} (A^{\nu} \vec{e}_{\nu})$
- $= \tilde{e}^{\mu} \partial_{\mu} (A^{\nu} \vec{e}_{\nu})$
- $= \partial_{\mu}A^{\nu}\tilde{e}^{\mu}\otimes\vec{e}_{\nu} + A^{\nu}\tilde{e}^{\mu}\otimes\partial_{\mu}\vec{e}_{\nu}$

define $\Gamma^{\lambda}_{\nu\mu}\vec{e}_{\lambda} := \partial_{\mu}\vec{e}_{\nu}$ Christoffel symbols of second kind (symmetric defn)

SO $\widetilde{\nabla}\vec{A} = \partial_{\mu}A^{\nu}\tilde{e}^{\mu}\otimes\vec{e}_{\nu} + A^{\nu}\tilde{e}^{\mu}\otimes\Gamma^{\lambda}_{\ \nu\mu}\vec{e}_{\lambda}$

 $= \partial_{\mu}A^{\nu}\tilde{e}^{\mu} \otimes \vec{e}_{\nu} + A^{\nu}\Gamma^{\lambda}_{\ \nu\mu}\tilde{e}^{\mu} \otimes \vec{e}_{\lambda} \text{ since any } \Gamma^{\lambda}_{\ \nu\mu} \text{ is a scalar}$

- $\widetilde{\nabla}\vec{A} = \widetilde{\nabla}(A^{\nu}\vec{e}_{\nu})$
- $= \tilde{e}^{\mu} \partial_{\mu} (A^{\nu} \vec{e}_{\nu})$
- $= \partial_{\mu}A^{\nu}\tilde{e}^{\mu}\otimes\vec{e}_{\nu} + A^{\nu}\tilde{e}^{\mu}\otimes\partial_{\mu}\vec{e}_{\nu}$

define $\Gamma^{\lambda}_{\nu\mu}\vec{e}_{\lambda} := \partial_{\mu}\vec{e}_{\nu}$ Christoffel symbols of second kind (symmetric defn)

$$\mathbf{SO} \ \widetilde{\nabla} \vec{A} = \partial_{\mu} A^{\nu} \tilde{e}^{\mu} \otimes \vec{e}_{\nu} + A^{\nu} \tilde{e}^{\mu} \otimes \Gamma^{\lambda}_{\nu\mu} \vec{e}_{\lambda}$$
$$= \partial_{\mu} A^{\nu} \tilde{e}^{\mu} \otimes \vec{e}_{\nu} + A^{\lambda} \Gamma^{\nu}_{\lambda\mu} \tilde{e}^{\mu} \otimes \vec{e}_{\nu}$$

since name of summation index is arbitrary, e.g. $\sum_{\lambda} x^{-2\lambda} = \sum_{\mu} x^{-2\mu} = \sum_{\nu} x^{-2\nu}$

 $1 \cdot \Lambda \ \Sigma \cdot \tilde{\mathrm{d}} \phi \cdot \langle \rangle \cdot \mathbf{g} \cdot x^{\mu} \cdot \mathrm{d} s^2 \cdot \nabla \vec{A}; \quad M \cdot \nabla_V = 0 \cdot \mathbf{R} \cdot \rangle \quad \mathsf{SR}_{\mathsf{f}} \epsilon \mathsf{GR}$

- $\widetilde{\nabla} \vec{A} = \widetilde{\nabla} (A^{\nu} \vec{e}_{\nu})$
- $= \tilde{e}^{\mu} \partial_{\mu} (A^{\nu} \vec{e}_{\nu})$
- $= \partial_{\mu}A^{\nu}\tilde{e}^{\mu}\otimes\vec{e}_{\nu} + A^{\nu}\tilde{e}^{\mu}\otimes\partial_{\mu}\vec{e}_{\nu}$

define $\Gamma^{\lambda}_{\nu\mu}\vec{e}_{\lambda} := \partial_{\mu}\vec{e}_{\nu}$ Christoffel symbols of second kind (symmetric defn)

$$\begin{aligned} \mathbf{SO} \ \widetilde{\nabla} \vec{A} &= \partial_{\mu} A^{\nu} \tilde{e}^{\mu} \otimes \vec{e}_{\nu} + A^{\nu} \tilde{e}^{\mu} \otimes \Gamma^{\lambda}_{\nu\mu} \vec{e}_{\lambda} \\ &= \partial_{\mu} A^{\nu} \tilde{e}^{\mu} \otimes \vec{e}_{\nu} + A^{\lambda} \Gamma^{\nu}_{\lambda\mu} \tilde{e}^{\mu} \otimes \vec{e}_{\nu} \\ &= (\partial_{\mu} A^{\nu} + A^{\lambda} \Gamma^{\nu}_{\lambda\mu}) \tilde{e}^{\mu} \otimes \vec{e}_{\nu} \end{aligned}$$

 $1 \cdot \Lambda \ \Sigma \cdot \tilde{\mathrm{d}} \phi \cdot \langle \rangle \cdot \mathbf{g} \cdot x^{\mu} \cdot \mathrm{d} s^2 \cdot \nabla \vec{A}; \quad M \cdot \nabla_V = 0 \cdot \mathbf{R} \cdot \rangle \quad \mathrm{SR}_{\mathbf{f}} \epsilon \mathrm{GR}$

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$$\nabla_{\mu}A^{\nu} := A^{\nu}_{;\mu} := \partial_{\mu}A^{\nu} + A^{\lambda}\Gamma^{\nu}_{\lambda\mu}$$

w:covariant derivative of vector

 $1\cdot\Lambda\ \Sigma\cdot\tilde{\mathrm{d}}\phi\cdot\langle\rangle\cdot\mathbf{g}\cdot x^{\mu}\cdot\mathrm{d}s^{2}\cdot\nabla\vec{A}\,;\,M\cdot\nabla_{V}=0\cdot\mathbf{R}\cdot\rangle\quad \mathrm{SR+}\epsilon\mathrm{GR}$

- $\widetilde{\nabla} \vec{A} = \widetilde{\nabla} (A^{\nu} \vec{e}_{\nu})$
- $= \tilde{e}^{\mu} \partial_{\mu} (A^{\nu} \vec{e}_{\nu})$
- $= \partial_{\mu}A^{\nu}\tilde{e}^{\mu}\otimes\vec{e}_{\nu} + A^{\nu}\tilde{e}^{\mu}\otimes\partial_{\mu}\vec{e}_{\nu}$
- define $\Gamma^{\lambda}_{\nu\mu}\vec{e}_{\lambda} := \partial_{\mu}\vec{e}_{\nu}$ Christoffel symbols of second kind (symmetric defn)
- $\mathbf{SO} \ \widetilde{\nabla} \vec{A} = \partial_{\mu} A^{\nu} \tilde{e}^{\mu} \otimes \vec{e}_{\nu} + A^{\nu} \tilde{e}^{\mu} \otimes \Gamma^{\lambda}_{\nu\mu} \vec{e}_{\lambda}$ $= \partial_{\mu} A^{\nu} \tilde{e}^{\mu} \otimes \vec{e}_{\nu} + A^{\lambda} \Gamma^{\nu}_{\lambda\mu} \tilde{e}^{\mu} \otimes \vec{e}_{\nu}$ $= (\partial_{\mu} A^{\nu} + A^{\lambda} \Gamma^{\nu}_{\lambda\mu}) \tilde{e}^{\mu} \otimes \vec{e}_{\nu}$

$$\nabla_{\mu}A^{\nu} := A^{\nu}_{;\mu} := A^{\nu}_{,\mu} + A^{\lambda}\Gamma^{\nu}_{\lambda\mu}$$

w:covariant derivative of vector

 $1 \cdot \Lambda \ \Sigma \cdot \tilde{\mathrm{d}} \phi \cdot \langle \rangle \cdot \mathbf{g} \cdot x^{\mu} \cdot \mathrm{d} s^2 \cdot \nabla \vec{A}; \quad M \cdot \nabla_V = 0 \cdot \mathbf{R} \cdot \rangle \quad \mathrm{SR}_{\mathbf{f}} \epsilon \mathrm{GR}$

mathematically deeper: $\widetilde{\nabla}$, usually written just as ∇ , is the <u>w:Levi-Civita connection</u>

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so far we showed how $\widetilde{\nabla}$ applied to a (0,0)-tensor field = scalar field ϕ gives a (0,1)-tensor field = one-form field = $(\widetilde{d}\phi)_{\mu}\tilde{e}^{\mu}$

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angle\cdot \mathsf{g}\cdot x^{\mu}\cdot\mathrm{d}s^{2}\cdot
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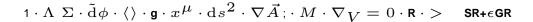
• not components of tensor: $\Gamma^{\nu}_{\lambda\mu}$

how does a one-form change with position? $\widetilde{\nabla} \tilde{A} = ?$

evaluating $\widetilde{\nabla} \widetilde{A}$ as we did $\widetilde{\nabla} \overrightarrow{A}$ shows that we again need $\partial_{\mu} \widetilde{e}^{\nu} = F^{\nu}_{\lambda\mu} \widetilde{e}^{\lambda}$ for some coefficients $F^{\nu}_{\lambda\mu}$

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relation between vectors and one-forms: $\langle \tilde{e}^{\nu}, \vec{e}_{\lambda} \rangle = \delta^{\nu}_{\lambda}$ $\partial_{\mu} \delta^{\nu}_{\lambda} = 0$ (obviously)

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$$0 = \partial_{\mu} \delta^{\nu}_{\ \lambda} = \partial_{\mu} \left(\left\langle \tilde{e}^{\nu}, \vec{e}_{\lambda} \right\rangle \right)$$

GR: gradient of one-form $\widetilde{\nabla} \widetilde{A}$

evaluating $\widetilde{\nabla} \widetilde{A}$ as we did $\widetilde{\nabla} \overrightarrow{A}$ shows that we again need $\partial_{\mu} \widetilde{e}^{\nu} = F^{\nu}_{\lambda\mu} \widetilde{e}^{\lambda}$ for some coefficients $F^{\nu}_{\lambda\mu}$

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$$0 = \partial_{\mu} \delta^{\nu}_{\ \lambda} = \partial_{\mu} \left(\left\langle \tilde{e}^{\nu}, \vec{e}_{\lambda} \right\rangle \right)$$

can we use the product rule with this scalar product? $\partial_{\mu}\left(\left\langle \tilde{A}, \vec{B} \right\rangle\right) = ?$

evaluating $\widetilde{\nabla} \widetilde{A}$ as we did $\widetilde{\nabla} \vec{A}$ shows that we again need $\partial_{\mu} \tilde{e}^{\nu} = F^{\nu}_{\lambda\mu} \tilde{e}^{\lambda}$ for some coefficients $F^{\nu}_{\lambda\mu}$

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can we use the product rule with this scalar product? $\partial_{\mu}\left(\left\langle \tilde{A}, \vec{B} \right\rangle\right) = \partial_{\mu}\left(A_{\nu}B^{\nu}\right)$ in some coordinate basis

evaluating $\widetilde{\nabla} \widetilde{A}$ as we did $\widetilde{\nabla} \vec{A}$ shows that we again need $\partial_{\mu} \tilde{e}^{\nu} = F^{\nu}_{\lambda\mu} \tilde{e}^{\lambda}$ for some coefficients $F^{\nu}_{\lambda\mu}$

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can we use the product rule with this scalar product?

$$\partial_{\mu}\left(\left\langle \tilde{A}, \vec{B} \right\rangle\right) = \partial_{\mu}\left(A_{\nu}B^{\nu}\right)$$

 $= (\partial_{\mu}A_{\nu})B^{\nu} + A_{\nu}(\partial_{\mu}B^{\nu})$ by product rule on functions

evaluating $\widetilde{\nabla} \widetilde{A}$ as we did $\widetilde{\nabla} \overrightarrow{A}$ shows that we again need $\partial_{\mu} \widetilde{e}^{\nu} = F^{\nu}_{\lambda\mu} \widetilde{e}^{\lambda}$ for some coefficients $F^{\nu}_{\lambda\mu}$

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$$\partial_{\mu} \left(\left\langle \tilde{A}, \vec{B} \right\rangle \right) = \partial_{\mu} \left(A_{\nu} B^{\nu} \right)$$
$$= \left(\partial_{\mu} A_{\nu} \right) B^{\nu} + A_{\nu} \left(\partial_{\mu} B^{\nu} \right)$$
$$= \left\langle \partial_{\mu} \tilde{A}, \vec{B} \right\rangle + \left\langle \tilde{A}, \partial_{\mu} \vec{B} \right\rangle \text{ since }$$
$$\partial_{\mu} \tilde{A} = \left(\partial_{\mu} A_{0}, \partial_{\mu} A_{1}, \partial_{\mu} A_{2}, \partial_{\mu} A_{3} \right)$$

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so $0 = \left\langle \partial_{\mu} \tilde{e}^{\nu}, \vec{e}_{\lambda} \right\rangle + \left\langle \tilde{e}^{\nu}, \partial_{\mu} \vec{e}_{\lambda} \right\rangle$

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 $1\cdot\Lambda\ \Sigma\cdot\tilde{\mathrm{d}}\phi\cdot\langle\rangle\cdot\mathbf{g}\cdot x^{\mu}\cdot\mathrm{d}s^{2}\cdot\nabla\vec{A}\,;\,\cdot\,M\cdot\nabla_{V}=0\cdot\mathbf{R}\cdot\rangle\quad\mathrm{SR+}\epsilon\mathrm{GR}$

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similarly, we can write the (0,3)-tensor $\widetilde{\nabla} \mathbf{g} = (\nabla_{\lambda} g_{\mu\nu}) \tilde{e}^{\lambda} \otimes \tilde{e}^{\mu} \otimes \tilde{e}^{\nu}$ giving $\nabla_{\lambda} g_{\mu\nu} = \partial_{\lambda} g_{\mu\nu} - \Gamma^{\kappa}_{\ \mu\lambda} g_{\kappa\nu} - \Gamma^{\kappa}_{\ \nu\lambda} g_{\mu\kappa}$

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Do we know anything interesting about $\widetilde{\nabla} \mathbf{g}$ for the manifolds of interest to GR?

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Do we know anything interesting about $\widetilde{\nabla} \mathbf{g}$ for the manifolds of interest to GR?

First, we need a rough description of the manifolds we need for GR.

topological manifold *M* w:Manifold#Mathematical_definition

only topological properties needed

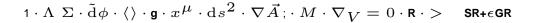
topological manifold *M* w:Manifold#Mathematical_definition

- only topological properties needed
- no differentiability, no metric needed

topological manifold *M* w:Manifold#Mathematical_definition

• only topological properties needed

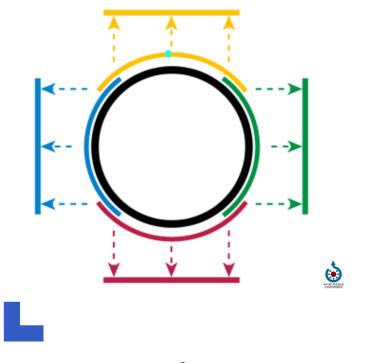
next: relation with \mathbb{R}^4 (or M^4)



topological manifold *M* w:Manifold#Mathematical_definition

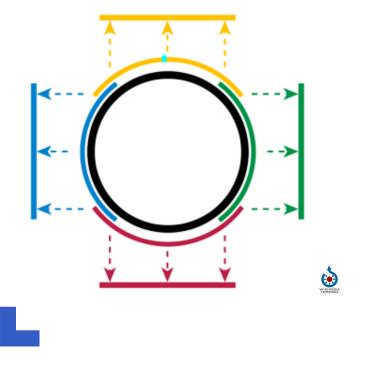
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topological manifold *M* w:Manifold#Mathematical_definition

 \bullet only topological properties needed next: relation with \mathbb{R}^4 (or $M^4)$

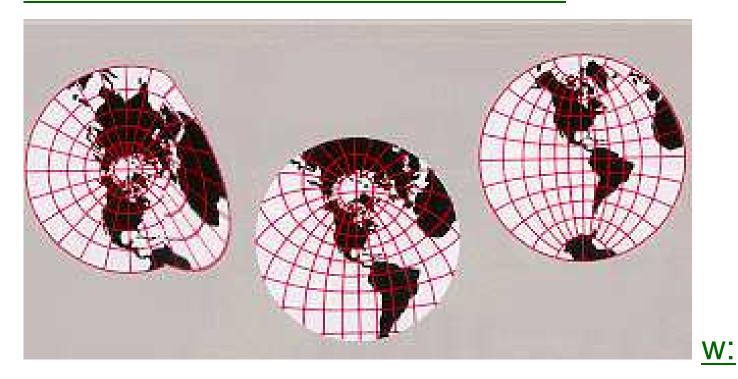


w:Manifold

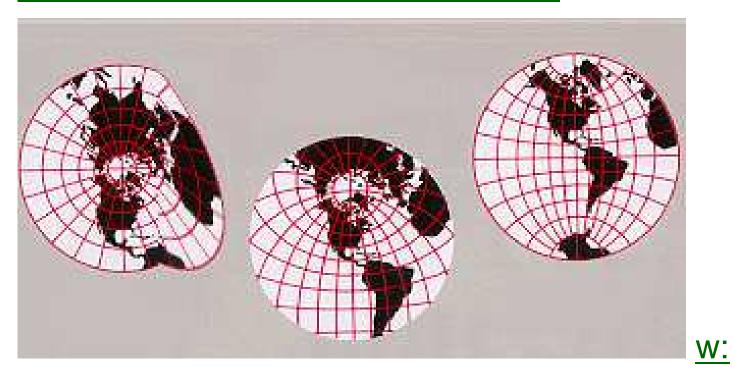
- chart := function ϕ_{α} from part of pseudo-4manifold M to part of M⁴ (Minkowski)
- atlas := set of overlapping charts that cover M

if every transition chart $:= \phi_{\beta} \circ \phi_{\alpha}^{-1}$ in an atlas for M is differentiable on \mathbb{R}^4 (or M^4), then M is a w:differentiable 4-(pseudo-)manifold

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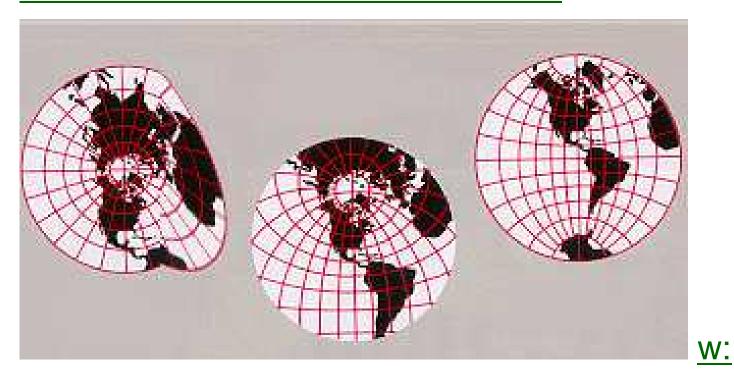


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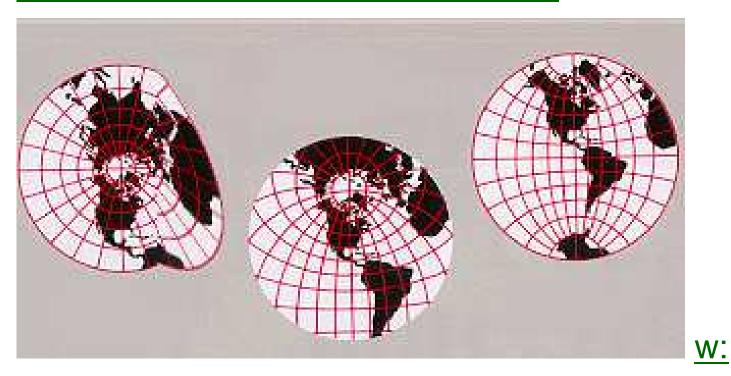
projections (left-to-right) ϕ_1 , ϕ_2 , ϕ_3 from S^2 to \mathbb{R}^2

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 ϕ_1 is not differentiable, so $\phi_1 \circ \phi_2^{-1}$ is not differentiable

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atlas not enough to show that $S^2 = differentiable$ 2-manifold

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if a (pseudo-)<u>w:Riemannian metric</u> g can be added to M, then (M,g) is a (pseudo-)Riemannian 4-manifold if g has signature (1, n - 1) (i.e. (-, +, +, +) or (+, -, -, -), etc.), then (M,g) is a Lorentzian *n*-manifold

topological manifolds

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differentiable (pseudo-)manifolds

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smooth (pseudo-)manifolds

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_orentzian manifolds

Lorentzian 4-manifolds

GR: smooth manifold and $\widetilde{\nabla}\mathbf{g}$

topological manifolds

differentiable (pseudo-)manifolds

smooth (pseudo-)manifolds

(pseudo-)Riemannian manifolds

Lorentzian manifolds

Lorentzian 4-manifolds

GR: assume that spacetime is a Lorentzian 4-manifold

GR: smooth manifold and $\widetilde{\nabla}\mathbf{g}$

from above:

$$\nabla_{\lambda}g_{\mu\nu} = \partial_{\lambda}g_{\mu\nu} - \Gamma^{\kappa}_{\ \mu\lambda}g_{\kappa\nu} - \Gamma^{\kappa}_{\ \nu\lambda}g_{\mu\kappa}$$

 $1\cdot\Lambda\ \Sigma\cdot\tilde{\mathrm{d}}\phi\cdot\langle\rangle\cdot\mathbf{g}\cdot x^{\mu}\cdot\mathrm{d}s^{2}\cdot\nabla\vec{A}\,;\,\cdot\,M\cdot\nabla_{V}=0\cdot\mathbf{R}\cdot\rangle\quad\mathrm{SR+}\epsilon\mathrm{GR}$

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in the tangent space at \mathbf{x} , \exists coordinate basis $\vec{e}_{\bar{\mu}}$ with $g_{\bar{\mu}\bar{\nu}} = \eta_{\bar{\mu}\bar{\nu}} = \text{diag}(-1, 1, 1, 1) = g^{\bar{\mu}\bar{\nu}}$

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so $\widetilde{\nabla} \mathbf{g} = \mathbf{0}$ (also $\widetilde{\nabla} \mathbf{g}^{-1} = 0$) on the tangent space, since if true in one coord system, also true in others

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warning: $\{x^{\mu}(\lambda)\}\$ at some λ on the manifold is a point on the manifold but NOT a vector; while $d\vec{x}$ — in the tangent space — IS a vector

using $\vec{V}(\lambda) := \frac{d\vec{x}}{d\lambda}$, project covariant derivative to curve using scalar product \langle , \rangle

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or in a coordinate basis...

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 ∇_V written by Bertschinger without \neg or \sim because ∇_V T of tensor T has the same tensor order as T

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for a vector field:

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$$\begin{split} \frac{\mathrm{d}\phi}{\mathrm{d}\lambda} &\equiv \nabla_V \phi := \left\langle \widetilde{\nabla}\phi, \vec{V} \right\rangle \\ &= V^{\mu} \partial_{\mu} \phi \\ \frac{\mathrm{d}\vec{A}}{\mathrm{d}\lambda} &\equiv \nabla_V \vec{A} := \left\langle \widetilde{\nabla}\vec{A}, \vec{V} \right\rangle \\ &= V^{\mu} (\nabla_{\mu} A^{\nu}) \vec{e}_{\nu} \\ &= V^{\mu} (A^{\nu}_{,\mu} + A^{\kappa} \Gamma^{\nu}_{\,\kappa\mu}) \vec{e}_{\nu} \end{split}$$

so in a coord basis,

$$\nabla_V \vec{A} = \left(\frac{\mathrm{d}A^{\nu}}{\mathrm{d}\lambda} + V^{\mu} A^{\kappa} \Gamma^{\nu}_{\ \kappa\mu}\right) \vec{e_{\nu}}$$

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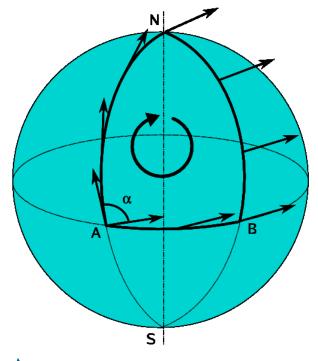
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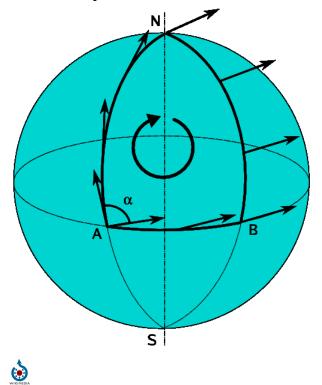
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example:



on S², parallel transport of \vec{A} around a closed loop does not conserve \vec{A} 's direction

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- tensorial definition independent of coordinate basis
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cf w:Euler-Lagrange equation

parallel transport around "small" parallelogram in two directions $d\vec{x}_1$, $d\vec{x}_2$,

("1" and "2" are not component indices here)

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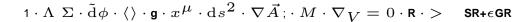
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- so **R** has second order partial derivatives of $g_{\nu\kappa}, \ldots$

• first order ∂ :

(pseudo-)manifold locally like \mathbb{R}^3 (M^4), \exists coords where $\Gamma^{\mu}_{~\nu\beta}=0$ locally

• first order ∂ :

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• second order ∂ :

(pseudo-)manifold globally like \mathbb{R}^3 (M⁴) $\Leftrightarrow R^{\mu}_{\nu\alpha\beta}(\mathbf{x}) = 0 \ \forall \mathbf{x}$

... second Bianchi identity:

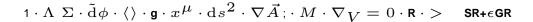
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warning: "R" written coordinate-free may mean:

- an order 4, dimension 64 tensor R;
- an order 2, dimension 16 tensor **R** or *R*; or
- an order 0, dimension 1 tensor \equiv scalar R
- all three are fields over a spacetime 4-manifold

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w:List of formulas in Riemannian geometry

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can be thought of as a *consequence* of the model



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Cactus - http://cactuscode.org