

Averaging of disturbed Friedmann–Lemaitre cosmological model

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Outline

- The model and its properties
- Averaging with weak limit
- Construction of the background

Motivation

- Being inspired by the perturbation theory, the aim is to construct a disturbed version of the Friedmann–Lemaître cosmological model by simple generalization of its metric.
- We are interested in a model with the following properties
 - metric is lorentzian
 - energy density is positive
 - growing inhomogeneities are allowed
 - inhomogeneities are finite
- What can and what can not be achieved by performing such a construction?
- Is a model, even unphysical, with such properties possible? If it exists, what are its limitations?
- Explicit models can be used for testing different averaging strategies.

The model

- The metric, the matter four-velocity and the Einstein equations are

$$g_{\mu\nu} = \begin{pmatrix} -b^2 e^{ce} & 0 & 0 & 0 \\ 0 & a^2 e^{cd(x)} & 0 & 0 \\ 0 & 0 & a^2 e^{cd(y)} & 0 \\ 0 & 0 & 0 & a^2 e^{cd(z)} \end{pmatrix}, \quad u_\nu = \begin{pmatrix} -b e^{\frac{1}{2}ce} \\ 0 \\ 0 \\ 0 \end{pmatrix},$$

$$\kappa T_{\mu\nu} = G_{\mu\nu},$$

where

$$a = t^\alpha, \quad b = \varpi t^{\frac{3}{2}\alpha\gamma-1}, \quad c = \varepsilon t^\delta, \quad d(z) = \sin^{2\beta} \left(\frac{z}{\lambda} \right) + \alpha,$$

$$e = \frac{3}{2}(\varsigma_1 + \alpha)F, \quad F = f + \frac{2}{3} \frac{\varsigma_2 - \varsigma_1^2}{(\varsigma_1 + \alpha)^2}, \quad \varsigma_s = \frac{(2s\beta)!}{2^{2s\beta} ((s\beta)!)^2}.$$

- Spacetime of the model is lorentzian, inhomogeneous and does not satisfy the weak energy condition.

The model

- Energy density in the model is non-negative, infinite in zero and finite in infinity when

$$\delta < 0 \text{ and}$$

$$\left(\varkappa < 0 \wedge \gamma \leq 0 \wedge F \leq 0 \wedge (\varepsilon < 0 \wedge \alpha \leq -1 \vee 0 < \varepsilon \wedge 0 \leq \alpha) \right) \text{ or}$$

$$0 < \varkappa \wedge 0 \leq \gamma \wedge 0 \leq F \wedge (\varepsilon < 0 \wedge 0 \leq \alpha \vee 0 < \varepsilon \wedge \alpha \leq -1) \Big) \text{ or}$$

$$0 < \delta \text{ and}$$

$$\left(\varkappa < 0 \wedge \gamma < 0 \wedge (F < 0 \vee F = 0 \wedge 2\delta \leq 3\varkappa\gamma) \wedge (\varepsilon < 0 \wedge 0 \leq \alpha \vee 0 < \varepsilon \wedge \alpha \leq -1) \right) \text{ or}$$

$$0 < \varkappa \wedge 0 < \gamma \wedge (0 < F \vee F = 0 \wedge 2\delta \leq 3\varkappa\gamma) \wedge (\varepsilon < 0 \wedge \alpha \leq -1 \vee 0 < \varepsilon \wedge 0 \leq \alpha) \Big).$$

- Quantities, ${}^3R_{\mu\nu}$, $H_{\mu\nu}$, η_ν , $\omega_{\mu\nu}$, q_ν vanish. The anisotropic stress tensor is given by

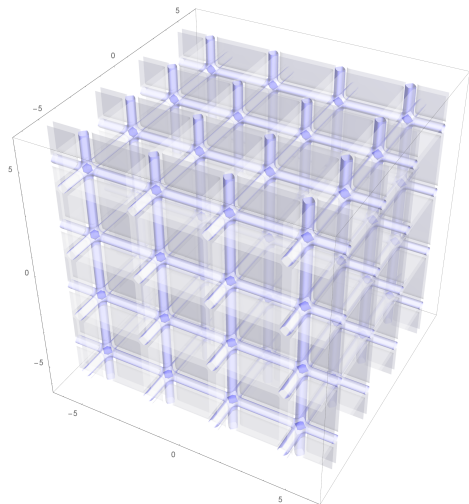
$$\kappa\pi_{\mu\nu} = u^\alpha \nabla_\alpha \sigma_{\mu\nu} + \theta \sigma_{\mu\nu},$$

and the energy density

$$\kappa\epsilon = \frac{1}{3}\theta^2 - \frac{1}{2}\sigma^{\beta\alpha}\sigma_{\beta\alpha}.$$

The model

- Isodensity surfaces



The model

- There are nine parameters in the model, $\alpha, \beta, \gamma, \delta, \varepsilon, f, \varkappa, \varpi, \lambda$.
- Parameters ε, ϖ are irrelevant due to scaling symmetry in the model. Physical quantities (e.g. age) are invariant under the transformation

$$\delta \rightarrow \nu \mu^{-\frac{2}{3} \frac{\gamma}{\varkappa}} \delta, \quad \varepsilon \rightarrow \mu \varepsilon, \quad \varkappa \rightarrow \nu \mu^{-\frac{2}{3} \frac{\gamma}{\varkappa}} \varkappa, \quad \varpi \rightarrow \nu \varpi.$$

- The parameter λ scales spatial coordinates. The distance between overdensities is $\lambda \pi$.
- The parameter δ controls the growth rate of inhomogeneities. They grow when it is positive.

The model

- The inhomogeneities are convex when α is positive. This parameter controls also the density contrast, $\frac{\epsilon|_{x,y,z=\lambda\pi/2}}{\epsilon|_{x,y,z=0}} - 1$. It grows from 0 at zero time to $\frac{1}{\alpha}(2 + \frac{1}{\alpha})$ at infinity and is finite.
- The parameter β controls the width of the inhomogeneities. Their width at half maximum decreases from $\frac{2}{\pi} \arccos(\frac{1}{2})^{\frac{1}{2\beta}}$ at zero to $\frac{2}{\pi} \arccos(\sqrt{\frac{1}{2} + \alpha + \alpha^2} - \alpha)^{\frac{1}{2\beta}}$ at infinity.
- The matter in the model evolves in time. There is a limited control over its behavior. The adiabatic index of the cosmic fluid, $\frac{p}{\epsilon} + 1$, changes from γ to f . The quotient $\frac{\delta}{\varkappa}$ decides on the character of the phase transition (smooth or violent).
- The shear viscosity coefficient, the ratio of $\pi^\mu{}_\nu$ to $-2\theta\sigma^\mu{}_\nu$, changes from $\frac{1}{4}(\gamma - \frac{2}{3}\frac{\delta}{\varkappa} - 2)$ to $\frac{1}{4}(f + \frac{2}{3}\frac{s_2 - s_1^2}{(s_1 + \alpha)^2} - 2)$. For reasonable values of parameters, it is negative at zero time.
- It is possible to match the basic property of the Universe that at the Hubble constant near 70 km/Mpc/s its age is about 13 Gyr.

Averaging of the model

- We average the model by averaging individual fields in it. We use the weak limit as a tool (Burnett (1989), Green and Wald (2011))

$$\lim_{\lambda \rightarrow 0} \int f^{a_1 \dots a_n} \left(A_{a_1 \dots a_n}(\lambda) - \text{w-lim}_{\lambda \rightarrow 0} A_{a_1 \dots a_n}(\lambda) \right) = 0.$$

For example, $\text{w-lim}_{\lambda \rightarrow 0} \sin^2\left(\frac{z}{\lambda}\right) = \frac{1}{2}$.

- Here, it is needed the following

$$\begin{aligned} & \text{w-lim}_{\lambda \rightarrow 0} \sin^{2s\beta} \left(\frac{z}{\lambda} \right) \exp \left(t \sin^{2\beta} \left(\frac{z}{\lambda} \right) \right) \\ &= \varsigma_s F \left(\left\{ \frac{2n-1}{2\beta} + s : n = 1, 2, \dots, \beta \right\}, \left\{ \frac{n}{\beta} + s : n = 1, 2, \dots, \beta \right\}, t \right). \end{aligned}$$

Averaging of the model

- The metric does not converge uniformly as $\lambda \rightarrow 0$. We need the weak limit to average it. Due to this, the Green–Wald theorems do not apply here.
- For averaged scalar quantities there occurs that

$$\kappa\langle\epsilon\rangle = \frac{1}{3}\langle\theta^2\rangle - \frac{1}{2}\langle\sigma^{\beta\alpha}\sigma_{\beta\alpha}\rangle = \frac{1}{3}\langle\theta\rangle^2.$$

- The averaging of tensor quantities reveals that

$$\langle\sigma_{\mu\nu}\rangle \neq 0, \quad \text{but} \quad \langle\sigma^\mu{}_\nu\rangle = 0, \quad \langle\pi_{\mu\nu}\rangle \neq 0, \quad \text{but} \quad \langle\pi^\mu{}_\nu\rangle = 0.$$

The weak limit depends on the valence of the averaged tensor.

- According to this, we should question that $\langle g_{\mu\nu}\rangle$ can be taken as the proper background metric.

The background model

- The Green–Wald approach

$$\langle G_{\mu\nu}(g) \rangle = G_{\mu\nu}(\bar{g}) - \kappa t_{\mu\nu}.$$

Taking $\langle g_{\mu\nu} \rangle$ as the background metric is unjustified. When \bar{g} and $t_{\mu\nu}$ are free this equation could be satisfied by infinitely many ways.

- For our model, there is one distinguished solution which could serve as the background model. It is given by the metric which reproduces the averaged mixed energy-momentum tensor alone

$$\langle G^{\mu}{}_{\nu}(g) \rangle = G^{\mu}{}_{\nu}(\bar{g}).$$

We assume here that the weak limit acting on a tensor with mixed indices gives reliable results.

The background model

- It is equivalent to searching for a metric reproducing averaged energy density $\langle \epsilon \rangle$ and pressure $\langle p \rangle$.
- We assume that the background metric has the Robertson–Walker symmetry

$$\bar{g}_{\mu\nu} = \begin{pmatrix} -B^2 & 0 & 0 & 0 \\ 0 & A^2 & 0 & 0 \\ 0 & 0 & A^2 & 0 \\ 0 & 0 & 0 & A^2 \end{pmatrix}.$$

- The equations are as follows

$$\frac{d}{dt} \ln A = -\frac{1}{3} \frac{\frac{d}{dt} \langle \epsilon \rangle}{\langle \epsilon \rangle + \langle p \rangle}, \quad B = -\frac{2}{3} \frac{\frac{d}{dt} \sqrt{\langle \epsilon \rangle}}{\langle \epsilon \rangle + \langle p \rangle}.$$

Right hand side of the differential equation is a rational function of time so it can be integrated explicitly.

- It appears that the expansion scalar is also reproduced, $\langle \theta(g) \rangle = \theta(\bar{g})$.

Outlook

- What are optical properties of the background model in comparison to the original model?
- What criterion should be applied to select the background model properly?
- How such an unphysical model could resemble so well the true Universe?