



# Monograph: Shape of the Universe

B.F. Roukema

(c) CC-BY-SA-3.0 and/or GPLv3 (details in progress!)

10 April 2014



# Introduction

- basic cosmology model:
  - ◆ GR: curvature = matter-energy content

# Introduction

- basic cosmology model:
  - ◆ GR: curvature = matter-energy content
  - ◆ verbal averaging: homogeneous, isotropic spatial slices

# Introduction

- basic cosmology model:
  - ◆ GR: curvature = matter-energy content
  - ◆ verbal averaging: homogeneous, isotropic spatial slices
  - ◆ shape of space (curvature + topology)

# Introduction

## ■ basic cosmology model:

- ◆ GR: curvature = matter-energy content
  - ◆ verbal averaging: homogeneous, isotropic spatial slices
  - ◆ shape of space (curvature + topology)
- + observations of expansion  $\Rightarrow$  hot big bang

# Introduction

## ■ basic cosmology model:

- ◆ GR: curvature = matter-energy content
  - ◆ verbal averaging: homogeneous, isotropic spatial slices
  - ◆ shape of space (curvature + topology)
- + observations of expansion  $\Rightarrow$  hot big bang
- $\Rightarrow$  black body radiation, nucleosynthesis

# Introduction

## ■ basic cosmology model:

- ◆ GR: curvature = matter-energy content
  - ◆ verbal averaging: homogeneous, isotropic spatial slices
  - ◆ shape of space (curvature + topology)
- + observations of expansion  $\Rightarrow$  hot big bang
- $\Rightarrow$  **black body radiation, nucleosynthesis**
- ◆ but:  $\exists$  galaxies  $\Rightarrow$  inhomogeneous, anisotropic spatial slices

# Introduction

## ■ basic cosmology model:

◆ GR: curvature = matter-energy content

◆ verbal averaging: homogeneous, isotropic spatial slices

◆ shape of space (curvature + topology)

+ observations of expansion  $\Rightarrow$  hot big bang

$\Rightarrow$  black body radiation, nucleosynthesis

◆ but:  $\exists$  galaxies  $\Rightarrow$  inhomogeneous, anisotropic spatial slices

◆ standard model: density perturbations (anisotropy)



# Introduction

## ■ basic cosmology model:

◆ GR: curvature = matter-energy content

◆ verbal averaging: homogeneous, isotropic spatial slices

◆ shape of space (curvature + topology)

+ observations of expansion  $\Rightarrow$  hot big bang

$\Rightarrow$  black body radiation, nucleosynthesis

◆ but:  $\exists$  galaxies  $\Rightarrow$  inhomogeneous, anisotropic spatial slices

◆ standard model: density perturbations (anisotropy)

◆ verbal averaging: can we do better?

# Introduction

## ■ basic cosmology model:

◆ GR: curvature = matter-energy content

◆ verbal averaging: homogeneous, isotropic spatial slices

◆ shape of space (curvature + topology)

+ observations of expansion  $\Rightarrow$  hot big bang

$\Rightarrow$  black body radiation, nucleosynthesis

◆ but:  $\exists$  galaxies  $\Rightarrow$  inhomogeneous, anisotropic spatial slices

◆ standard model: density perturbations (anisotropy)

◆ ~~verbal averaging~~

# Introduction

## ■ basic cosmology model:

- ◆ GR: curvature = matter-energy content

- ◆ verbal averaging: homogeneous, isotropic spatial slices

- ◆ shape of space (curvature + topology)

+ observations of expansion  $\Rightarrow$  hot big bang

$\Rightarrow$  **black body radiation, nucleosynthesis**

- ◆ but:  $\exists$  galaxies  $\Rightarrow$  inhomogeneous, anisotropic spatial slices

- ◆ standard model: **density perturbations (anisotropy)**

- ◆ scalar (GR) averaging: statistically homogeneous spatial slices

## ■ within this model, what is the shape of the Universe?

# verbal averaging

- w:Cosmological principle

# verbal averaging

- w:Cosmological principle
- practical meaning:

# verbal averaging

- w:Cosmological principle

- practical meaning:

1. assume homogeneity and isotropy

# verbal averaging

- w:Cosmological principle

- practical meaning:

1. assume homogeneity and isotropy
2. find the (differential 4-pseudo-manifold, metric) pairs  $(M, \mathbf{g})$  that solve  $\mathbf{G} = 8\pi\mathbf{T}$

# verbal averaging

- w:Cosmological principle

- practical meaning:

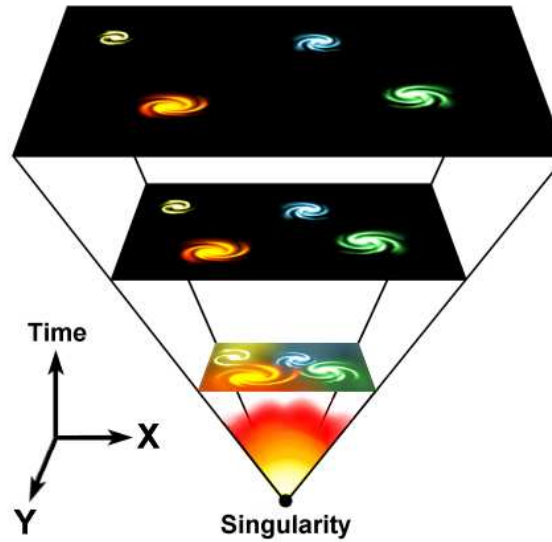
1. assume homogeneity and isotropy
2. find the (differential 4-pseudo-manifold, metric) pairs  $(M, \mathbf{g})$  that solve  $\mathbf{G} = 8\pi\mathbf{T}$
3. assume that  $(M, \mathbf{g})$  remains unchanged if we add density perturbations to an early time slice



# verbal averaging

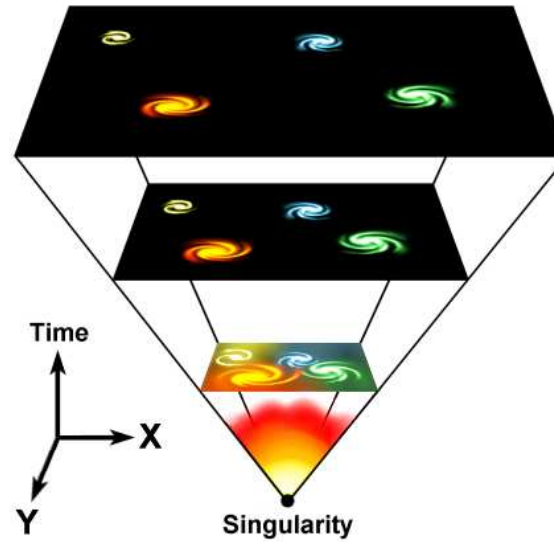
- w: Comoving coordinates

# verbal averaging



- w: Comoving coordinates

# verbal averaging

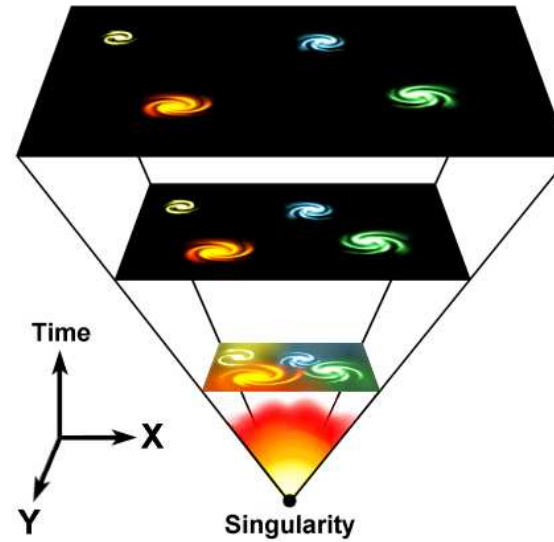


- w: Comoving coordinates



$$\Delta x(t) = a(t) \Delta r$$

# verbal averaging



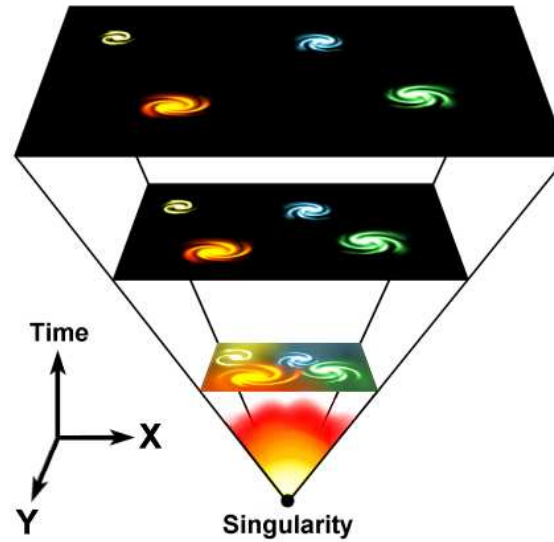
- w: Comoving coordinates



$$\Delta x(t) = a(t) \Delta r$$

- spherical coordinates for spatial slice

# verbal averaging

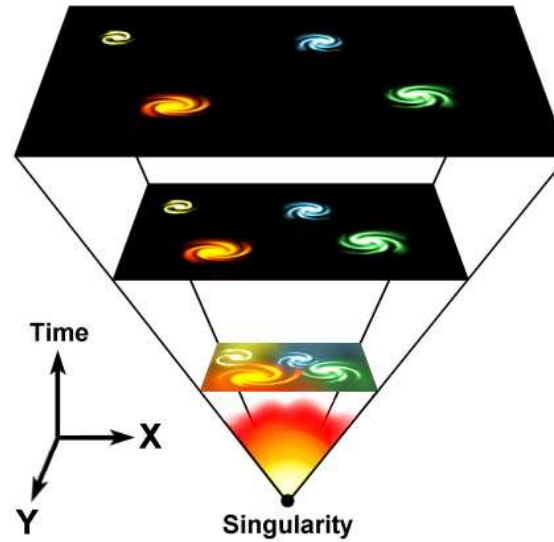


- w: Comoving coordinates



$$\int_{(t,r_1,\theta,\phi)}^{(t,r_2,\theta,\phi)} ds = a(t)\Delta r = a(t)|r_2 - r_1|$$

# verbal averaging



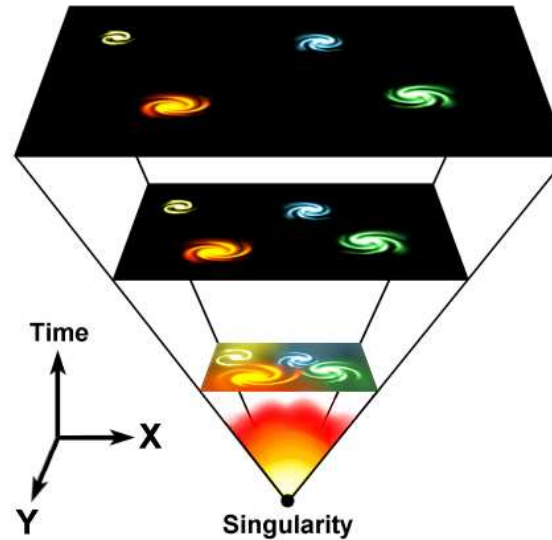
- w: Comoving coordinates

- 

$$\int_{(t,r_1,\theta,\phi)}^{(t,r_2,\theta,\phi)} ds = a(t)\Delta r = a(t)|r_2 - r_1|$$

where all expansion/contraction  $\rightarrow$  w: scale factor  $a(t)$

# verbal averaging



- w: Comoving coordinates



$$\int_{(t,r_1,\theta,\phi)}^{(t,r_2,\theta,\phi)} ds = a(t)\Delta r = a(t)|r_2 - r_1|$$

where all expansion/contraction  $\rightarrow$  w: scale factor  $a(t)$

- universe is static in comoving coordinates  $(r, \theta, \phi)$

# FLRW metric

- w:Friedmann–Lemaître–Robertson–Walker metric



# FLRW metric

- w:Friedmann–Lemaître–Robertson–Walker metric



- w: *A. P. Robertson* w:

w:Howard Percy Robertson

w:Arthur Geoffrey Walker

# FLRW metric

$$ds^2 = -dt^2 + \dots$$

# FLRW metric

$$ds^2 = -dt^2 + a^2(t) [\dots]$$

# FLRW metric

$$ds^2 = -dt^2 + a^2(t) [dr^2 + \dots]$$

# FLRW metric

$$ds^2 = -dt^2 + a^2(t) [dr^2 + r_{\perp}^2(\dots)]$$

# FLRW metric

$$ds^2 = -dt^2 + a^2(t) [dr^2 + r_{\perp}^2 (d\theta^2 + \cos^2 \theta d\phi^2)]$$

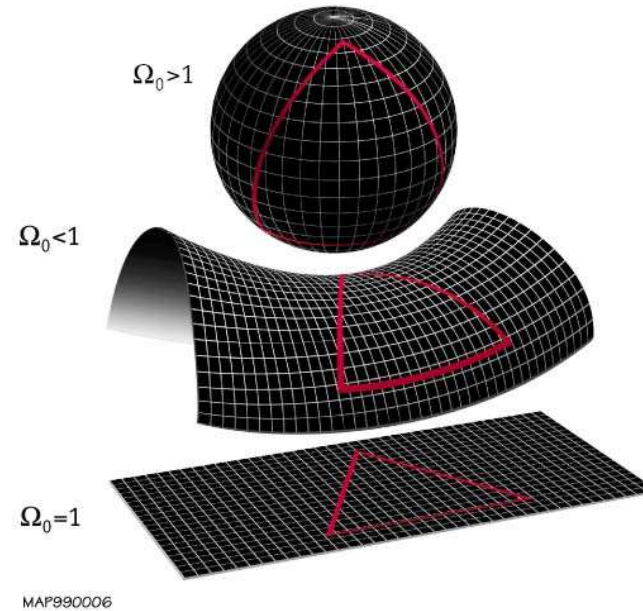
# FLRW metric

$$ds^2 = -dt^2 + a^2(t) [dr^2 + r_{\perp}^2 (d\theta^2 + \cos^2 \theta d\phi^2)]$$

$$\text{where } r_{\perp} := \begin{cases} R_C \sinh \frac{r}{R_C} & k < 0 \\ r & k = 0 \\ R_C \sin \frac{r}{R_C} & k > 0 \end{cases}$$

# FLRW metric

$$ds^2 = -dt^2 + a^2(t) [dr^2 + r_{\perp}^2 (d\theta^2 + \cos^2 \theta d\phi^2)]$$



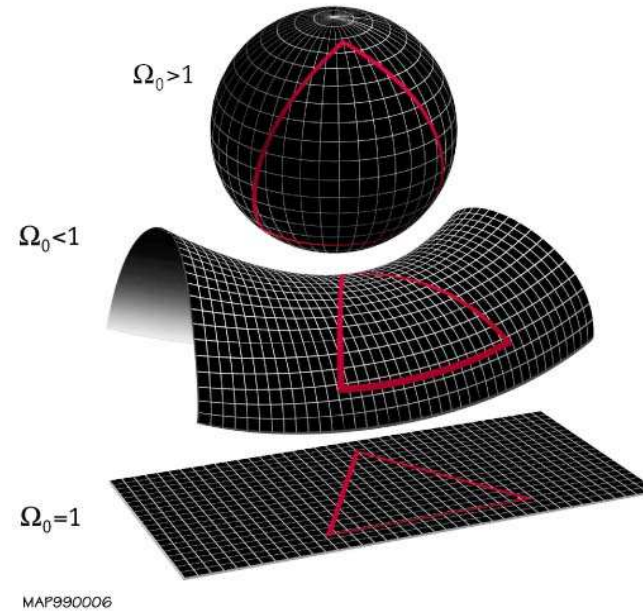
where  $r_{\perp} :=$

$$\begin{cases} R_C \sinh \frac{r}{R_C} & k < 0 \\ r & k = 0 \\ R_C \sin \frac{r}{R_C} & k > 0 \end{cases}$$



# FLRW metric

$$ds^2 = -dt^2 + a^2(t) [dr^2 + r_{\perp}^2 (d\theta^2 + \cos^2 \theta d\phi^2)]$$



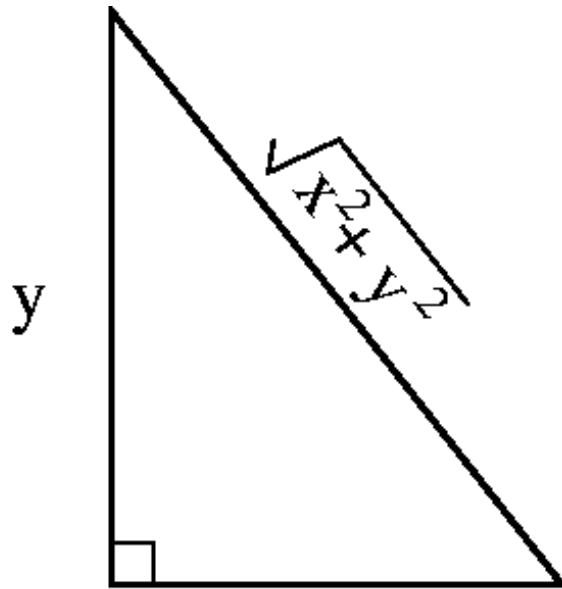
where  $r_{\perp} :=$

$$\begin{cases} R_C \sinh \frac{r}{R_C} & k < 0 \\ r & k = 0 \\ R_C \sin \frac{r}{R_C} & k > 0 \end{cases}$$

for a comoving radius of curvature  $R_C$  and curvature of sign  $k$

# curvature

- on a spatial slice (fixed value of  $t$ ):

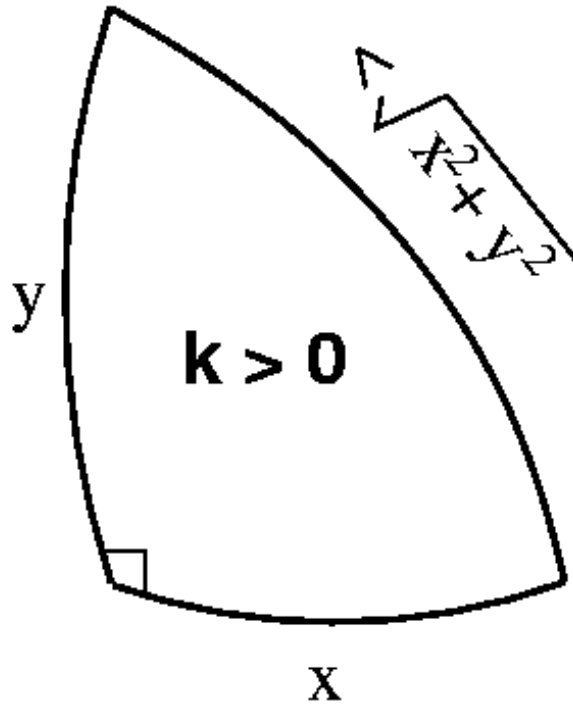


$x$

$$k = 0$$

# curvature

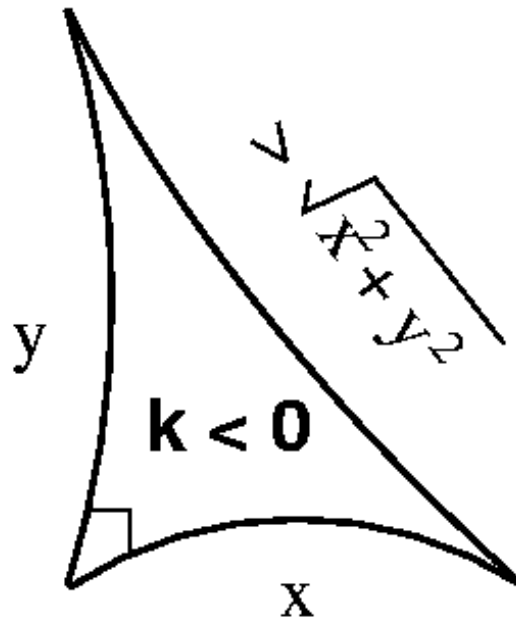
- on a spatial slice (fixed value of  $t$ ):



$$k > 0$$

# curvature

- on a spatial slice (fixed value of  $t$ ):



$$k < 0$$

# 2D curvature intuition: $k > 0$

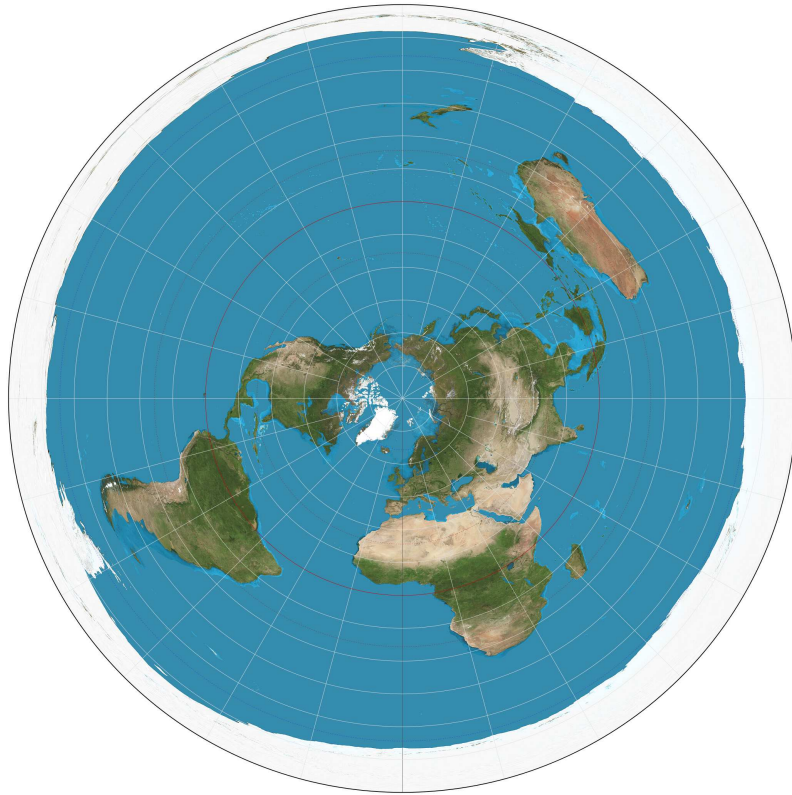
$$ds^2|_{\phi=\text{const}, a=1} = dr^2 + r_{\perp}^2 d\theta^2,$$

# 2D curvature intuition: $k > 0$

$$ds^2|_{\phi=\text{const}, a=1} = dr^2 + r_{\perp}^2 d\theta^2, \text{ where } r_{\perp} := R_C \sin \frac{r}{R_C}$$

# 2D curvature intuition: $k > 0$

$$ds^2|_{\phi=\text{const}, a=1} = dr^2 + r_{\perp}^2 d\theta^2, \text{ where } r_{\perp} := R_C \sin \frac{r}{R_C}$$

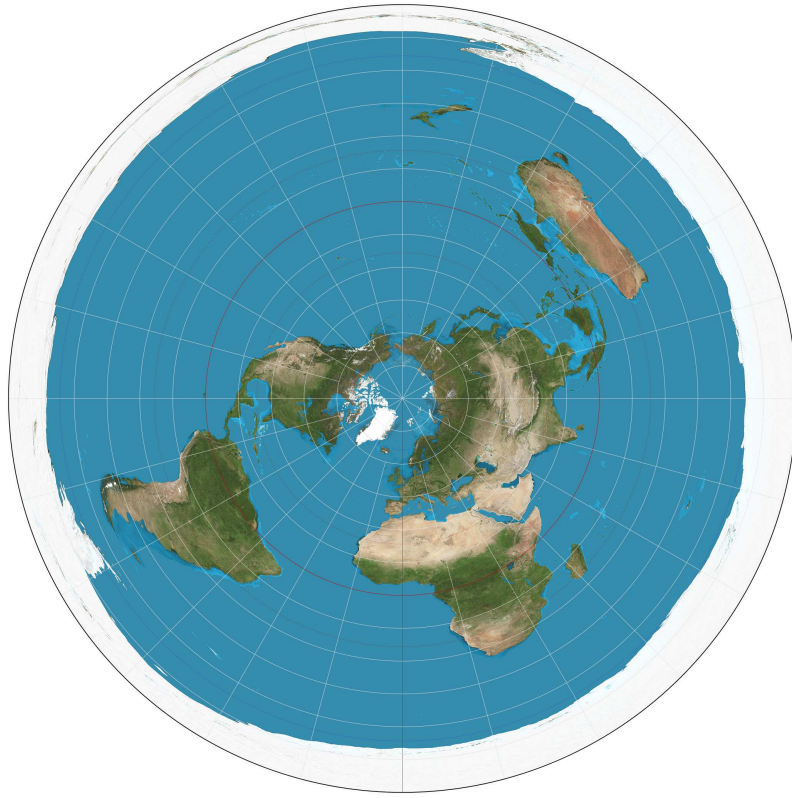


w:

(al-Biruni, c. 1000 CE)

# 2D curvature intuition: $k > 0$

$$ds^2|_{\phi=\text{const}, a=1} = dr^2 + r_{\perp}^2 d\theta^2, \text{ where } r_{\perp} := R_C \sin \frac{r}{R_C}$$



w:

(al-Biruni, c. 1000 CE)

■ intuition switch:  $S^2$  easier vs  $S^3$  more physical



# 2D topology intuition ( $k = 0$ )



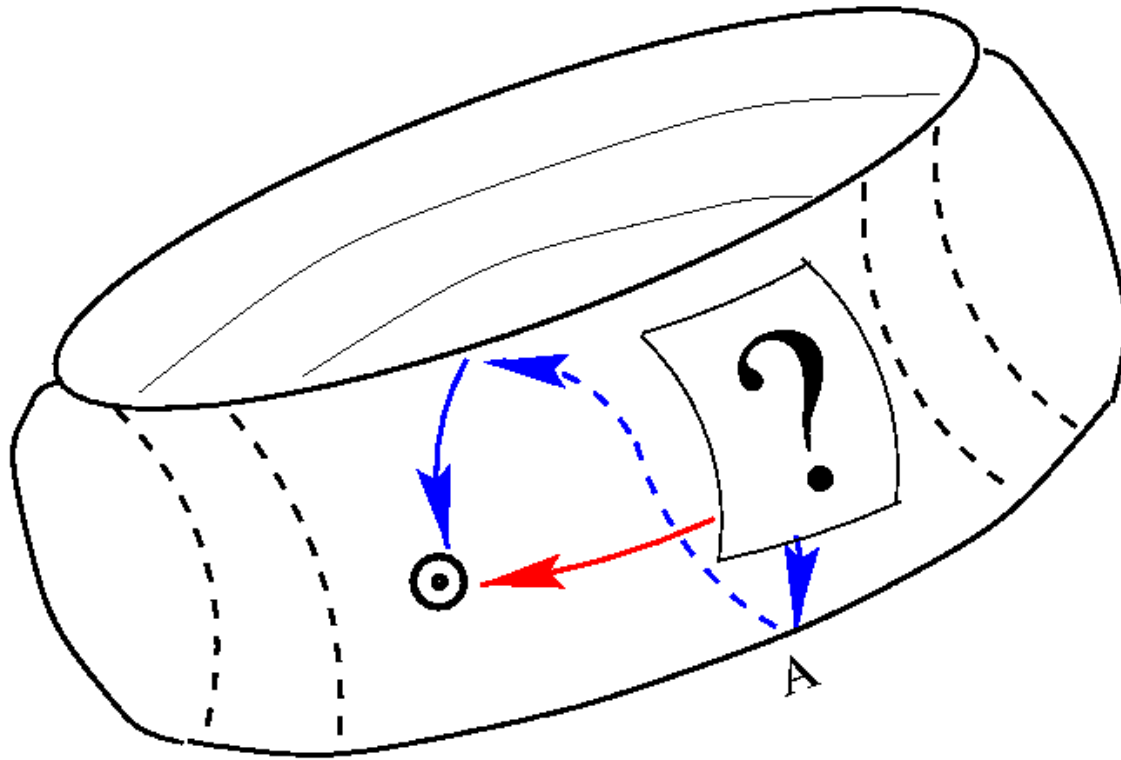
■ W:

# 2D topology intuition ( $k = 0$ )

- intuition 1: embed in higher dim. space

# 2D topology intuition ( $k = 0$ )

- intuition 1: embed in higher dim. space

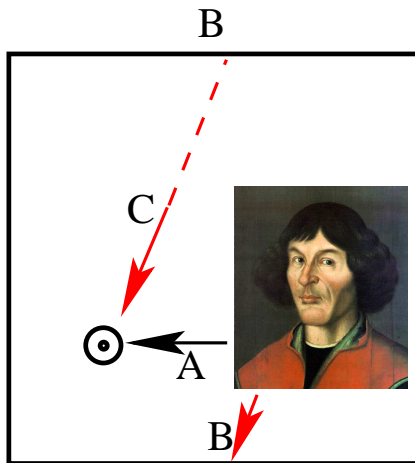


# 2D topology intuition ( $k = 0$ )

- intuition 1: embed in higher dim. space
- intuition 2: fundamental domain

# 2D topology intuition ( $k = 0$ )

- intuition 1: embed in higher dim. space
- intuition 2: fundamental domain

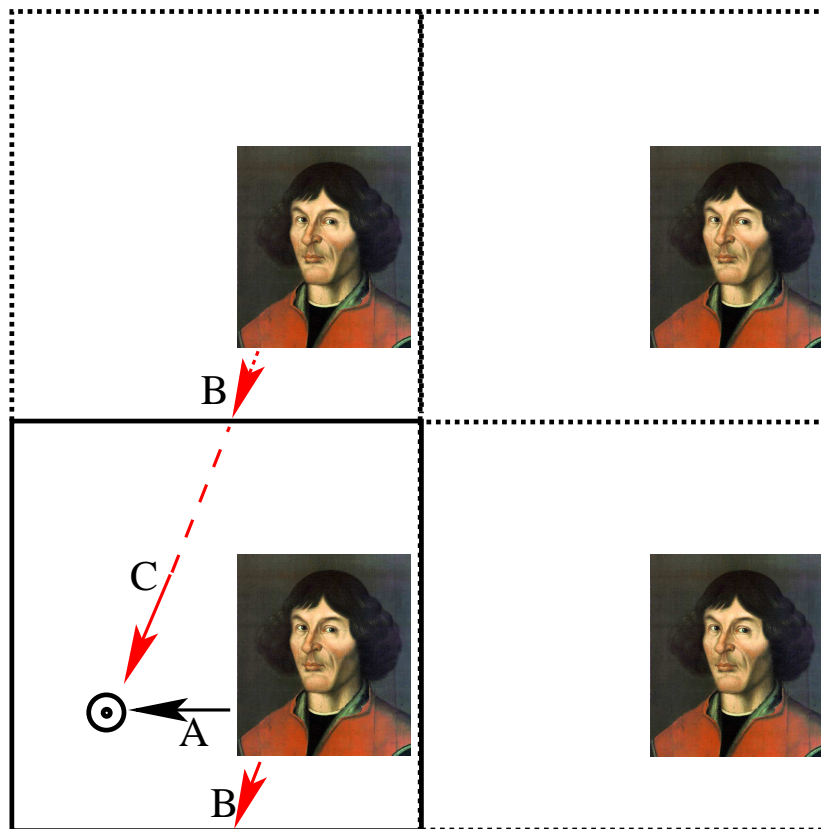


# 2D topology intuition ( $k = 0$ )

- intuition 1: embed in higher dim. space
- intuition 2: fundamental domain
- intuition 3: universal covering space

# 2D topology intuition ( $k = 0$ )

- intuition 1: embed in higher dim. space
- intuition 2: fundamental domain
- intuition 3: universal covering space



# expansion

- Hubble law (Lemaître 1927; [ADS:1927ASSB...47...49L](#)):  
 $\dot{a}(t_0)/a(t_0) \approx 600 \text{ km/s/Mpc}$



# expansion

- Hubble law (Lemaître 1927; [ADS:1927ASSB...47...49L](#)):  
 $\dot{a}(t_0)/a(t_0) \approx 600 \text{ km/s/Mpc} > 0$  (also Hubble 1929)

# expansion

- Hubble law (Lemaître 1927; [ADS:1927ASSB...47...49L](#)):  
 $\dot{a}(t_0)/a(t_0) \approx 600 \text{ km/s/Mpc} > 0$  (also Hubble 1929)
- so what is the “Big Bang”?

# expansion

- Hubble law (Lemaître 1927; [ADS:1927ASSB...47...49L](#)):  
 $\dot{a}(t_0)/a(t_0) \approx 600 \text{ km/s/Mpc} > 0$  (also Hubble 1929)
- so what is the “Big Bang”?
- it is:  $\exists t_b$  such that  $t \rightarrow t_b^+ \Rightarrow a(t) \rightarrow 0^+$

# expansion

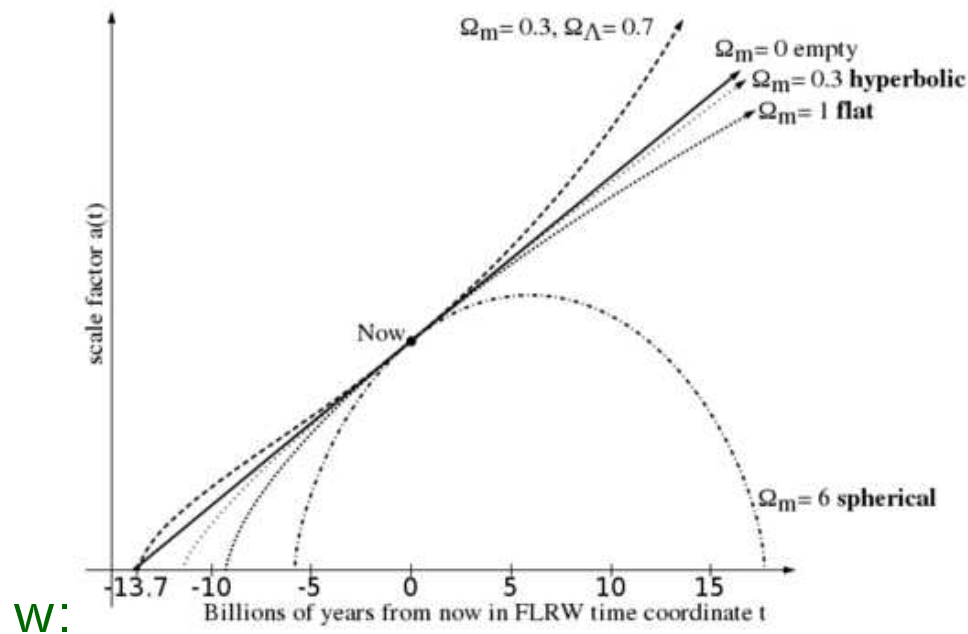
- Hubble law (Lemaître 1927; [ADS:1927ASSB...47...49L](#)):  
 $\dot{a}(t_0)/a(t_0) \approx 600 \text{ km/s/Mpc} > 0$  (also Hubble 1929)
- so what is the “Big Bang”?
- convention:  $t_b := 0$ , giving  $t \rightarrow 0^+ \Rightarrow a(t) \rightarrow 0^+$

# expansion

- Hubble law (Lemaître 1927; [ADS:1927ASSB...47...49L](#)):  
 $\dot{a}(t_0)/a(t_0) \approx 600 \text{ km/s/Mpc} > 0$  (also Hubble 1929)
- so what is the “Big Bang”?
- it is:  $\exists t_b$  such that  $t \rightarrow t_b^+ \Rightarrow a(t) \rightarrow 0^+$

# expansion

- Hubble law (Lemaître 1927; [ADS:1927ASSB...47...49L](#)):  
 $\dot{a}(t_0)/a(t_0) \approx 600 \text{ km/s/Mpc} > 0$  (also Hubble 1929)
- so what is the “Big Bang”?
- it is:  $\exists t_b$  such that  $t \rightarrow t_b^+ \Rightarrow a(t) \rightarrow 0^+$



# expansion

- matter density:  $\rho_m \propto a^{-3}$

# expansion

- matter density:  $\rho_m \propto a^{-3}$
- wavelength:  $\lambda \propto a$  [GR—transport four-velocity: Synge (1960); Narlikar (1994; [ADS:1994AmJPh..62..903N](#))]



# expansion

- matter density:  $\rho_m \propto a^{-3}$
- wavelength:  $\lambda \propto a$  [GR—transport four-velocity: Synge (1960); Narlikar (1994; [ADS:1994AmJPh..62..903N](#))]

⇔

$$\frac{\lambda_{\text{obs}}}{\lambda_{\text{em}}} = \frac{a_{\text{obs}}}{a_{\text{em}}}$$

# expansion

- matter density:  $\rho_m \propto a^{-3}$
- wavelength:  $\lambda \propto a$  [GR—transport four-velocity: Synge (1960); Narlikar (1994; [ADS:1994AmJPh..62..903N](#))]



$$\frac{\lambda_{\text{obs}}}{\lambda_{\text{em}}} = \frac{1}{a_{\text{em}}}$$

(Defn:  $a_0 := 1$ )

# expansion

- matter density:  $\rho_m \propto a^{-3}$
- wavelength:  $\lambda \propto a$  [GR—transport four-velocity: Synge (1960); Narlikar (1994; [ADS:1994AmJPh..62..903N](#))]



$$1 + z = \frac{1}{a_{\text{em}}}$$

(Defn of redshift  $z$ )

# expansion

- matter density:  $\rho_m \propto a^{-3} = (1+z)^3$
- wavelength:  $\lambda \propto a$  [GR—transport four-velocity: Synge (1960); Narlikar (1994; [ADS:1994AmJPh..62..903N](#))]



$$1 + z = a^{-1}$$

(light-cone convention:  $a$  often means  $a_{\text{em}}$ )

# expansion

- matter density:  $\rho_m \propto a^{-3} = (1+z)^3$
- wavelength:  $\lambda \propto a$  [GR—transport four-velocity: Synge (1960); Narlikar (1994; [ADS:1994AmJPh..62..903N](#))]



$$1 + z = a^{-1}$$

(light-cone convention:  $a$  often means  $a_{\text{em}}$ )

- radiation density:  $E = h\nu \Rightarrow \rho_r \propto a^{-4} = (1+z)^4$

# Black body: COBE ( $\sim 1992$ )

■ Planck's Law: 
$$I(\nu, T) = \frac{2h\nu^3}{c^2} \frac{1}{e^{\frac{h\nu}{kT}} - 1}$$

# Black body: COBE ( $\sim 1992$ )

■ Planck's Law: 
$$I(\nu, T) = \frac{2h\nu^3}{c^2} \frac{1}{e^{\frac{h\nu}{kT}} - 1}$$

■ wavelength:  $\lambda \propto a \Rightarrow$  frequency:  $\nu \propto a^{-1} = (1 + z)$

# Black body: COBE ( $\sim 1992$ )

- Planck's Law:  $I(\nu, T) = \frac{2h\nu^3}{c^2} \frac{1}{e^{\frac{h\nu}{kT}} - 1}$
- wavelength:  $\lambda \propto a \Rightarrow$  frequency:  $\nu \propto a^{-1} = (1 + z)$
- $\Rightarrow$  temperature:  $kT \propto h\nu \propto (1 + z)$



# Black body: COBE ( $\sim 1992$ )

- Planck's Law:  $I(\nu, T) = \frac{2h\nu^3}{c^2} \frac{1}{e^{\frac{h\nu}{kT}} - 1}$
- wavelength:  $\lambda \propto a \Rightarrow$  frequency:  $\nu \propto a^{-1} = (1 + z)$
- $\Rightarrow$  temperature:  $kT \propto h\nu \propto (1 + z)$
- $z \gg 1 \Rightarrow$  early Universe dominated by hot, dense plasma = protons, electrons, photons

# Black body: COBE ( $\sim 1992$ )

- Planck's Law:  $I(\nu, T) = \frac{2h\nu^3}{c^2} \frac{1}{e^{\frac{h\nu}{kT}} - 1}$
- wavelength:  $\lambda \propto a \Rightarrow$  frequency:  $\nu \propto a^{-1} = (1 + z)$
- $\Rightarrow$  temperature:  $kT \propto h\nu \propto (1 + z)$
- $z \gg 1 \Rightarrow$  early Universe dominated by hot, dense plasma = protons, electrons, photons
- $\Rightarrow$  black body

# Black body: COBE ( $\sim 1992$ )

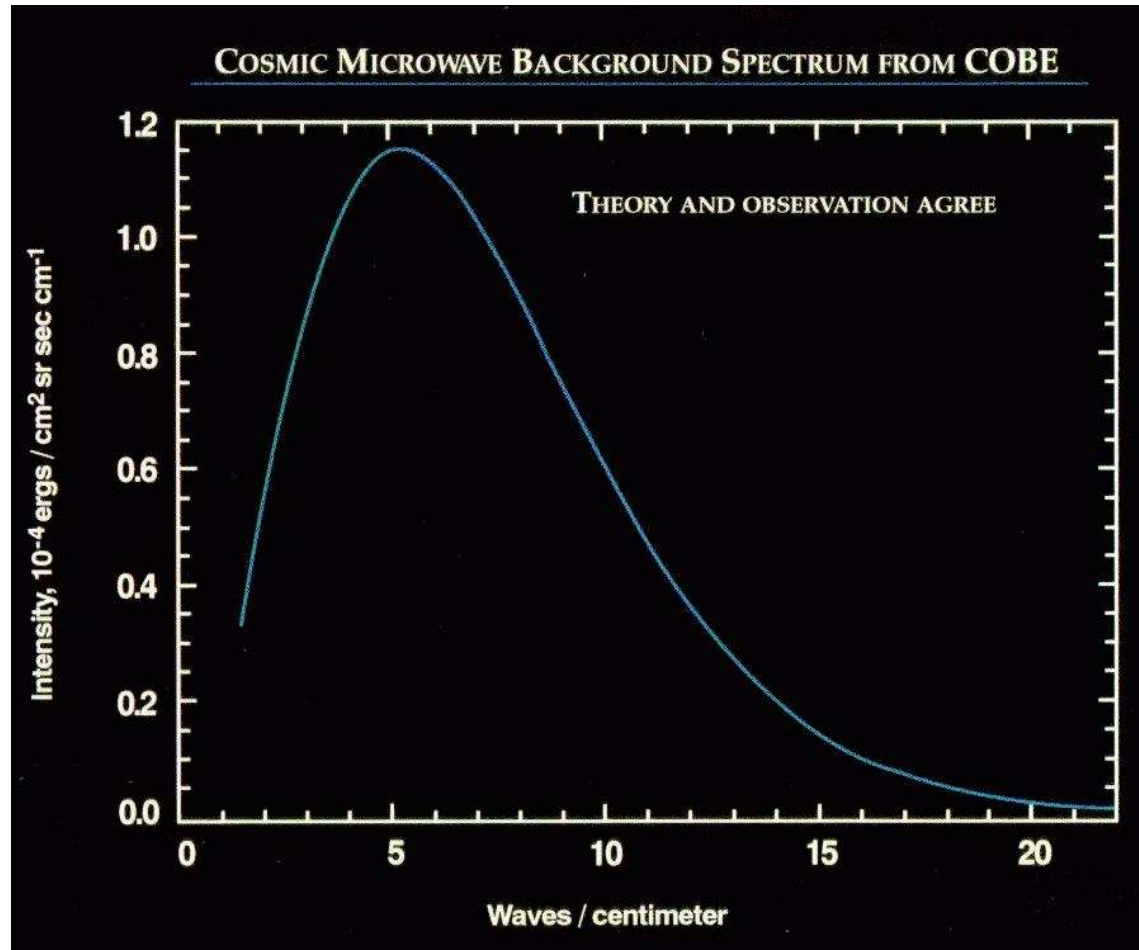
- Planck's Law:  $I(\nu, T) = \frac{2h\nu^3}{c^2} \frac{1}{e^{\frac{h\nu}{kT}} - 1}$
- wavelength:  $\lambda \propto a \Rightarrow$  frequency:  $\nu \propto a^{-1} = (1 + z)$
- $\Rightarrow$  temperature:  $kT \propto h\nu \propto (1 + z)$
- $z \gg 1 \Rightarrow$  early Universe dominated by hot, dense plasma = protons, electrons, photons
- $\Rightarrow$  black body + primordial nucleosynthesis

# Black body: COBE ( $\sim 1992$ )

- COBE /FIRAS (Far Infrared Absolute Spectrophotometer)

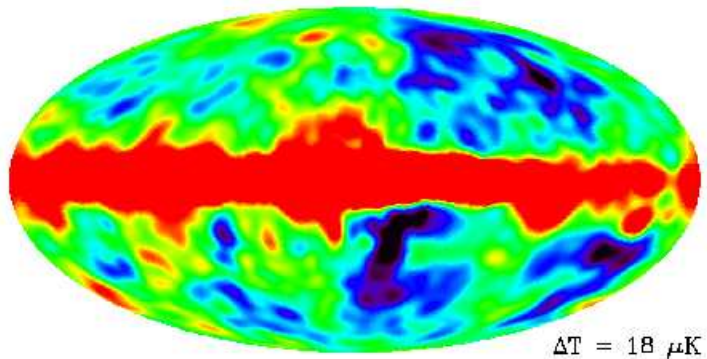
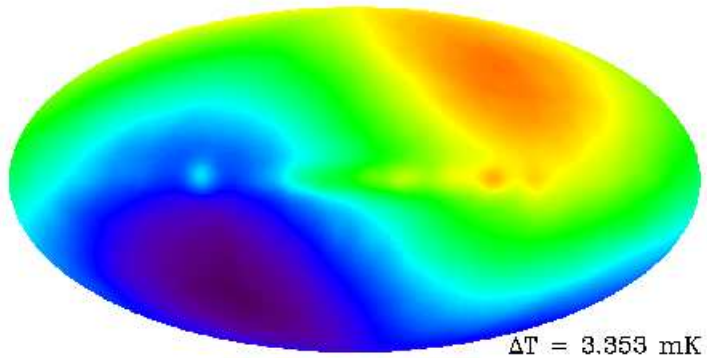
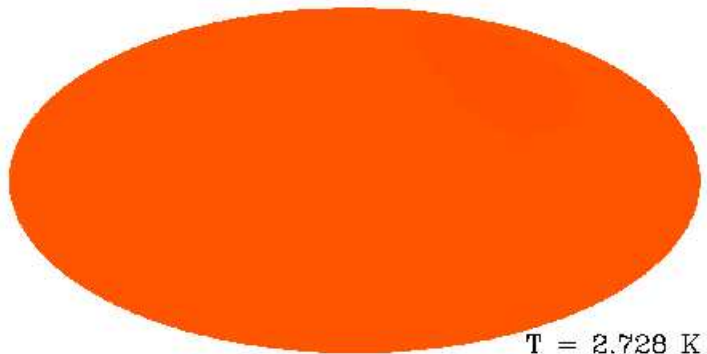
# Black body: COBE ( $\sim 1992$ )

- COBE /FIRAS (Far Infrared Absolute Spectrophotometer)



# Black body: COBE ( $\sim 1992$ )

- COBE /DMR (Differential Microwave Radiometer)

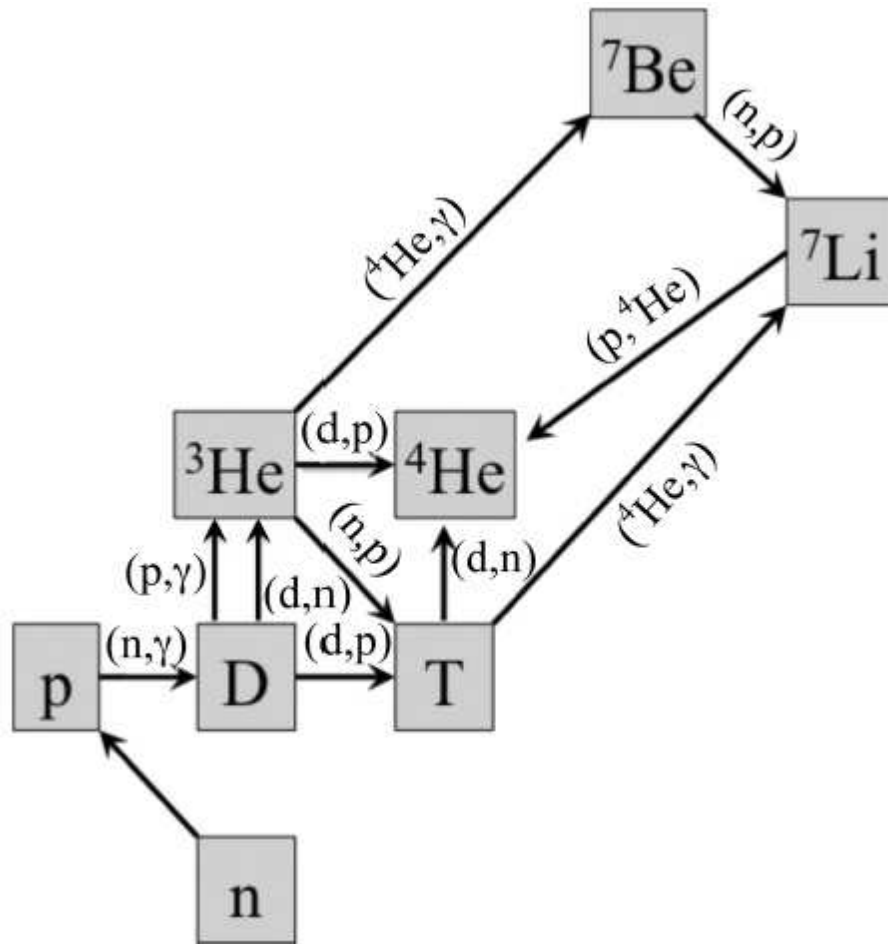


# BBN: Big bang nucleosynthesis

- Alpher, Bethe, & Gamow (1948; [ADS:1948PhRv...73..803A](#))

# BBN: Big bang nucleosynthesis

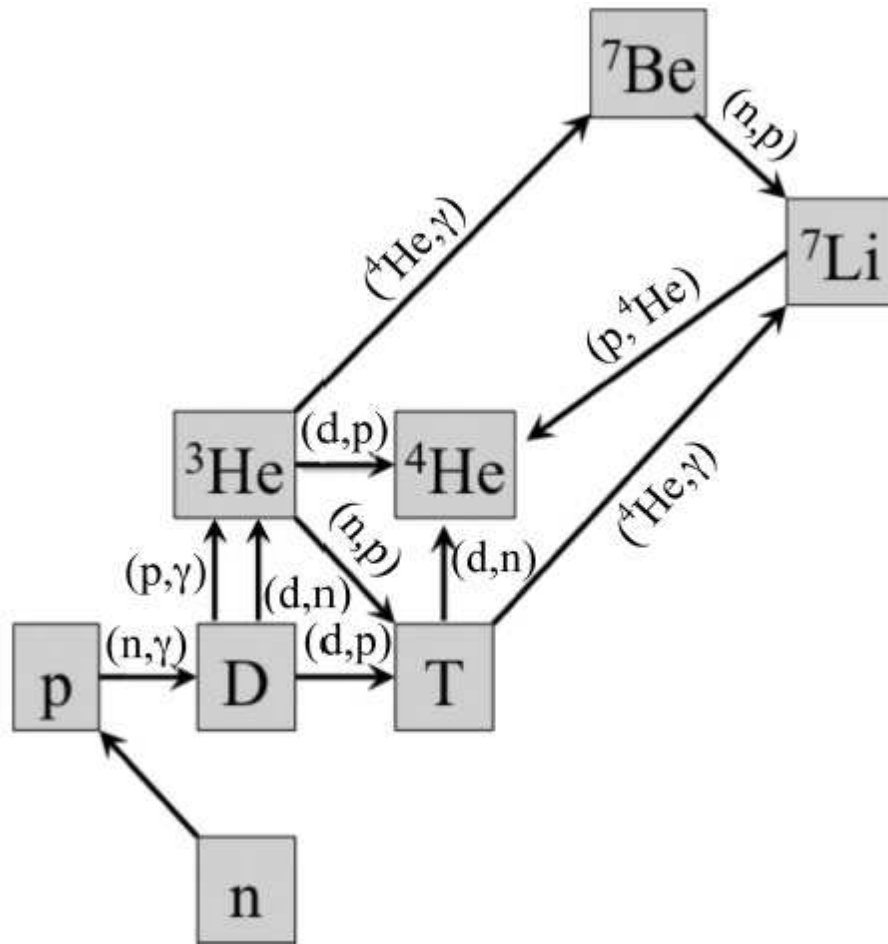
- Alpher, Bethe, & Gamow (1948; [ADS:1948PhRv...73..803A](#))





# BBN: Big bang nucleosynthesis

- Alpher, Bethe, & Gamow (1948; [ADS:1948PhRv...73..803A](#))



w:Big Bang nucleosynthesis

# FLRW: $a(t) = ?$

- second choice FLRW coord system:  $r :=$  w:orthographic projection  
of radial comoving distance (cf  $r :=$  radial comoving distance)

# FLRW: $a(t) = ?$

- second choice FLRW coord system:  $r :=$  w:orthographic projection of radial comoving distance (cf  $r :=$  radial comoving distance)
- $\Rightarrow$  coord singularity at equator ( $:= \pi/2$  from centre) if  $k > 0$

# FLRW: $a(t) = ?$

- second choice FLRW coord system:  $r :=$  w:orthographic projection of radial comoving distance (cf  $r :=$  radial comoving distance)

$\Rightarrow$  coord singularity at equator ( $:= \pi/2$  from centre) if  $k > 0$



# FLRW: $a(t) = ?$

- second choice FLRW coord system:  $r :=$  w:orthographic projection of radial comoving distance (cf  $r :=$  radial comoving distance)

$\Rightarrow$  coord singularity at equator ( $:= \pi/2$  from centre) if  $k > 0$

- $$ds^2 = -dt^2 + a^2(t) \left[ \frac{dr^2}{1 - kr^2} + r^2(d\theta^2 + \cos^2 \theta d\phi^2) \right]$$

# FLRW: $a(t) = ?$

- $$ds^2 = -dt^2 + a^2(t) \left[ \frac{dr^2}{1 - kr^2} + r^2(d\theta^2 + \cos^2 \theta d\phi^2) \right]$$

- universe content:  $\text{diag}(\mathbf{T}) = (-\rho, p, p, p)$

# FLRW: $a(t) = ?$

$$ds^2 = -dt^2 + a^2(t) \left[ \frac{dr^2}{1 - kr^2} + r^2(d\theta^2 + \cos^2 \theta d\phi^2) \right]$$

- universe content:  $\text{diag}(\mathbf{T}) = (-\rho, p, p, p)$
- MAXIMA: calculate  $\mathbf{G}$  and  $\mathbf{G} = 8\pi\mathbf{T}$  and simplify:  
<http://cosmo.torun.pl/Cosmo/FLRWEquationsGR>

# FLRW: $a(t) = ?$

- $$ds^2 = -dt^2 + a^2(t) \left[ \frac{dr^2}{1 - kr^2} + r^2(d\theta^2 + \cos^2 \theta d\phi^2) \right]$$

- universe content:  $\text{diag}(\mathbf{T}) = (-\rho, p, p, p)$

- MAXIMA: calculate  $\mathbf{G}$  and  $\mathbf{G} = 8\pi\mathbf{T}$  and simplify:  
<http://cosmo.torun.pl/Cosmo/FLRWEquationsGR>

- Friedmann Eqn: 
$$\frac{c^2 k}{a^2} + \frac{\dot{a}^2}{a^2} = \frac{8\pi G \rho}{3}$$



# FLRW: $a(t) = ?$

- $$ds^2 = -dt^2 + a^2(t) \left[ \frac{dr^2}{1 - kr^2} + r^2(d\theta^2 + \cos^2 \theta d\phi^2) \right]$$

- universe content:  $\text{diag}(\mathbf{T}) = (-\rho, p, p, p)$

- MAXIMA: calculate  $\mathbf{G}$  and  $\mathbf{G} = 8\pi\mathbf{T}$  and simplify:  
<http://cosmo.torun.pl/Cosmo/FLRWEquationsGR>

- Friedmann Eqn: 
$$\frac{c^2 k}{a^2} + \frac{\dot{a}^2}{a^2} = \frac{8\pi G \rho}{3}$$

- acceleration Eqn: 
$$\frac{\ddot{a}}{a} = -\frac{4\pi G (\rho + 3p/c^2)}{3}$$

# FLRW matter-dominated epoch

■ Friedmann Eqn:

$$\frac{\dot{a}^2}{a^2} = \frac{8\pi G\rho}{3} - \frac{c^2 k}{a^2}$$

# FLRW matter-dominated epoch

■ Friedmann Eqn: 
$$\frac{\dot{a}^2}{a^2} = \frac{8\pi G\rho}{3} - \frac{c^2 k}{a^2}$$

■ matter-dominated epoch:  $\rho = \rho_m$

# FLRW matter-dominated epoch

- Friedmann Eqn: 
$$\frac{\dot{a}^2}{a^2} = \frac{8\pi G\rho}{3} - \frac{c^2 k}{a^2}$$
- matter-dominated epoch:  $\rho = \rho_m = \rho_{m0} a^{-3}$

# FLRW matter-dominated epoch

- Friedmann Eqn:

$$\frac{\dot{a}^2}{a^2} = \frac{8\pi G \rho_{m0}}{3a^3} - \frac{c^2 k}{a^2}$$

- matter-dominated epoch:  $\rho = \rho_m = \rho_{m0} a^{-3}$

# FLRW matter-dominated epoch

- Friedmann Eqn:

$$\dot{a}^2 = \frac{8\pi G \rho_{m0}}{3a} - c^2 k$$

- matter-dominated epoch:  $\rho = \rho_m = \rho_{m0} a^{-3}$

# FLRW matter-dominated epoch

- Friedmann Eqn:

$$\dot{a}^2 = \frac{8\pi G \rho_{m0}}{3a} - c^2 k$$

- matter-dominated epoch:  $\rho = \rho_m = \rho_{m0} a^{-3}$

- $k = 0$  case:  $\dot{a}^2 \propto a^{-1}$

# FLRW matter-dominated epoch

- Friedmann Eqn:

$$\dot{a}^2 = \frac{8\pi G \rho_{m0}}{3a} - c^2 k$$

- matter-dominated epoch:  $\rho = \rho_m = \rho_{m0} a^{-3}$
- $k = 0$  case:  $\dot{a} \propto a^{-1/2}$  for  $a > 0$



# FLRW matter-dominated epoch

- Friedmann Eqn:

$$\dot{a}^2 = \frac{8\pi G \rho_{m0}}{3a} - c^2 k$$

- matter-dominated epoch:  $\rho = \rho_m = \rho_{m0} a^{-3}$

- $k = 0$  case:  $da \propto a^{-1/2} dt$  for  $a > 0$

# FLRW matter-dominated epoch

- Friedmann Eqn:

$$\dot{a}^2 = \frac{8\pi G \rho_{m0}}{3a} - c^2 k$$

- matter-dominated epoch:  $\rho = \rho_m = \rho_{m0} a^{-3}$

- $k = 0$  case:  $a^{1/2} da \propto dt$

# FLRW matter-dominated epoch

- Friedmann Eqn:

$$\dot{a}^2 = \frac{8\pi G \rho_{m0}}{3a} - c^2 k$$

- matter-dominated epoch:  $\rho = \rho_m = \rho_{m0} a^{-3}$

- $k = 0$  case:  $(2/3)a^{3/2} \propto t$

# FLRW matter-dominated epoch

- Friedmann Eqn:

$$\dot{a}^2 = \frac{8\pi G \rho_{m0}}{3a} - c^2 k$$

- matter-dominated epoch:  $\rho = \rho_m = \rho_{m0} a^{-3}$

- $k = 0$  case:  $a \propto t^{2/3}$

# FLRW matter-dominated epoch

- Friedmann Eqn:

$$\dot{a}^2 = \frac{8\pi G \rho_{m0}}{3a} - c^2 k$$

- matter-dominated epoch:  $\rho = \rho_m = \rho_{m0} a^{-3}$

- $k = 0$  case:

$$a = \left(\frac{t}{t_0}\right)^{2/3}$$

Einstein–de Sitter model (EdS)

# FLRW matter-dominated epoch

- Friedmann Eqn:

$$\dot{a}^2 = \frac{8\pi G \rho_{m0}}{3a} - c^2 k$$

- matter-dominated epoch:  $\rho = \rho_m = \rho_{m0} a^{-3}$

- $k = 0$  case:

$$a = \left(\frac{t}{t_0}\right)^{2/3}$$

Einstein–de Sitter model (EdS)

Defn: Hubble parameter  $H := \dot{a}/a$

# FLRW matter-dominated epoch

■ Friedmann Eqn:  $\dot{a}^2 = \frac{8\pi G \rho_{m0}}{3a} - c^2 k$

■ matter-dominated epoch:  $\rho = \rho_m = \rho_{m0} a^{-3}$

■  $k = 0$  case:  $a = \left(\frac{t}{t_0}\right)^{2/3}$  Einstein–de Sitter model (EdS)

Defn:  $H := \dot{a}/a$

$\Rightarrow$  Hubble constant  $H_0 := H(z = 0)$

# FLRW matter-dominated epoch

- Friedmann Eqn:

$$\dot{a}^2 = \frac{8\pi G \rho_{m0}}{3a} - c^2 k$$

- matter-dominated epoch:  $\rho = \rho_m = \rho_{m0} a^{-3}$

- $k = 0$  case:

$$a = \left(\frac{t}{t_0}\right)^{2/3}$$

Einstein–de Sitter model (EdS)

Defn: Hubble parameter  $H := \dot{a}/a$

$\Rightarrow$  Hubble constant  $H_0 := H(z = 0)$

$\Rightarrow$   $k = 0$  case:  $\frac{\dot{a}}{a} = \frac{2}{3t}$



# FLRW matter-dominated epoch

- Friedmann Eqn:

$$\dot{a}^2 = \frac{8\pi G \rho_{m0}}{3a} - c^2 k$$

- matter-dominated epoch:  $\rho = \rho_m = \rho_{m0} a^{-3}$

- $k = 0$  case:

$$a = \left(\frac{t}{t_0}\right)^{2/3}$$

Einstein–de Sitter model (EdS)

Defn: Hubble parameter  $H := \dot{a}/a$

$\Rightarrow$  Hubble constant  $H_0 := H(z = 0)$

$\Rightarrow$   $k = 0$  case:  $\frac{\dot{a}}{a} = \frac{2}{3t}$

$\Rightarrow H(t) = \frac{2}{3t}$

# FLRW matter-dominated epoch

- Friedmann Eqn:

$$\dot{a}^2 = \frac{8\pi G \rho_{m0}}{3a} - c^2 k$$

- matter-dominated epoch:  $\rho = \rho_m = \rho_{m0} a^{-3}$

- $k = 0$  case:

$$a = \left(\frac{t}{t_0}\right)^{2/3}$$

Einstein–de Sitter model (EdS)

Defn: Hubble parameter  $H := \dot{a}/a$

$\Rightarrow$  Hubble constant  $H_0 := H(z = 0)$

$\Rightarrow$   $k = 0$  case:  $\frac{\dot{a}}{a} = \frac{2}{3t}$

$\Rightarrow H(t) = \frac{2}{3t}; H_0 = H(t_0) = \frac{2}{3t_0}$

# FLRW matter-dominated epoch

- Friedmann Eqn:

$$\dot{a}^2 = \frac{8\pi G \rho_{m0}}{3a} - c^2 k$$

- matter-dominated epoch:  $\rho = \rho_m = \rho_{m0} a^{-3}$

- $k = 0$  case:  $a = \left(\frac{t}{t_0}\right)^{2/3}$  Einstein–de Sitter model (EdS)

$$\Rightarrow H(t) = \frac{2}{3t}; H_0 = H(t_0) = \frac{2}{3t_0}$$

- convenient conversion:  $1 \text{ km/s} \approx 1.04 \text{ kpc/Gyr} \approx 1 \text{ kpc/Gyr}$

# FLRW matter-dominated epoch

- Friedmann Eqn:

$$\dot{a}^2 = \frac{8\pi G \rho_{m0}}{3a} - c^2 k$$

- matter-dominated epoch:  $\rho = \rho_m = \rho_{m0} a^{-3}$

- $k = 0$  case:  $a = \left(\frac{t}{t_0}\right)^{2/3}$  Einstein–de Sitter model (EdS)

$$\Rightarrow H(t) = \frac{2}{3t}; H_0 = H(t_0) = \frac{2}{3t_0}$$

- Lemaître (1927) [ADS:1927ASSB...47...49L](#):  $H_0 \approx 600$  km/s/Mpc

# FLRW matter-dominated epoch

■ Friedmann Eqn: 
$$\dot{a}^2 = \frac{8\pi G \rho_{m0}}{3a} - c^2 k$$

■ matter-dominated epoch:  $\rho = \rho_m = \rho_{m0} a^{-3}$

■  $k = 0$  case: 
$$a = \left(\frac{t}{t_0}\right)^{2/3}$$
 Einstein–de Sitter model (EdS)

$\Rightarrow H(t) = \frac{2}{3t}; H_0 = H(t_0) = \frac{2}{3t_0}$

■ Lemaître (1927) [ADS:1927ASSB...47...49L](#):  $H_0 \approx 600$   
kpc/Gyr/Mpc

# FLRW matter-dominated epoch

- Friedmann Eqn:

$$\dot{a}^2 = \frac{8\pi G \rho_{m0}}{3a} - c^2 k$$

- matter-dominated epoch:  $\rho = \rho_m = \rho_{m0} a^{-3}$

- $k = 0$  case:  $a = \left(\frac{t}{t_0}\right)^{2/3}$  Einstein–de Sitter model (EdS)

$$\Rightarrow H(t) = \frac{2}{3t}; H_0 = H(t_0) = \frac{2}{3t_0}$$

- Lemaître (1927) [ADS:1927ASSB...47...49L](#):  $H_0 \approx 0.6 \text{ Gyr}^{-1}$

# FLRW matter-dominated epoch

- Friedmann Eqn:

$$\dot{a}^2 = \frac{8\pi G \rho_{m0}}{3a} - c^2 k$$

- matter-dominated epoch:  $\rho = \rho_m = \rho_{m0} a^{-3}$

- $k = 0$  case:  $a = \left(\frac{t}{t_0}\right)^{2/3}$  Einstein–de Sitter model (EdS)

$$\Rightarrow H(t) = \frac{2}{3t}; H_0 = H(t_0) = \frac{2}{3t_0}$$

- Lemaître (1927) [ADS:1927ASSB...47...49L](#):  $H_0 \approx 0.6 \text{ Gyr}^{-1}$

- Hubble (1929) [ADS:1929PNAS...15..168H](#):  $H_0 \approx 500 \text{ km/s/Mpc}$

# FLRW matter-dominated epoch

- Friedmann Eqn:

$$\dot{a}^2 = \frac{8\pi G \rho_{m0}}{3a} - c^2 k$$

- matter-dominated epoch:  $\rho = \rho_m = \rho_{m0} a^{-3}$

- $k = 0$  case:  $a = \left(\frac{t}{t_0}\right)^{2/3}$  Einstein–de Sitter model (EdS)

$$\Rightarrow H(t) = \frac{2}{3t}; H_0 = H(t_0) = \frac{2}{3t_0}$$

- Lemaître (1927) [ADS:1927ASSB...47...49L](#):  $H_0 \approx 0.6 \text{ Gyr}^{-1}$

- Hubble (1929) [ADS:1929PNAS...15..168H](#):  $H_0 \approx 500 \text{ kpc/Gyr/Mpc}$



# FLRW matter-dominated epoch

- Friedmann Eqn:

$$\dot{a}^2 = \frac{8\pi G \rho_{m0}}{3a} - c^2 k$$

- matter-dominated epoch:  $\rho = \rho_m = \rho_{m0} a^{-3}$

- $k = 0$  case:  $a = \left(\frac{t}{t_0}\right)^{2/3}$  Einstein–de Sitter model (EdS)

$$\Rightarrow H(t) = \frac{2}{3t}; H_0 = H(t_0) = \frac{2}{3t_0}$$

- Lemaître (1927) [ADS:1927ASSB...47...49L](#):  $H_0 \approx 0.6 \text{ Gyr}^{-1}$

- Hubble (1929) [ADS:1929PNAS...15..168H](#):  $H_0 \approx 0.5 \text{ Gyr}^{-1}$

# FLRW matter-dominated epoch

- Friedmann Eqn:

$$\dot{a}^2 = \frac{8\pi G \rho_{m0}}{3a} - c^2 k$$

- matter-dominated epoch:  $\rho = \rho_m = \rho_{m0} a^{-3}$

- $k = 0$  case:  $a = \left(\frac{t}{t_0}\right)^{2/3}$  Einstein–de Sitter model (EdS)

$$\Rightarrow H(t) = \frac{2}{3t}; H_0 = H(t_0) = \frac{2}{3t_0}$$

- Lemaître (1927) [ADS:1927ASSB...47...49L](#):  $H_0 \approx 0.6 \text{ Gyr}^{-1}$

- Hubble (1929) [ADS:1929PNAS...15..168H](#):  $H_0 \approx 0.5 \text{ Gyr}^{-1}$

$$\Rightarrow \text{EdS would give } t_0 = \frac{2}{3H_0}$$

# FLRW matter-dominated epoch

- Friedmann Eqn:

$$\dot{a}^2 = \frac{8\pi G \rho_{m0}}{3a} - c^2 k$$

- matter-dominated epoch:  $\rho = \rho_m = \rho_{m0} a^{-3}$

- $k = 0$  case:  $a = \left(\frac{t}{t_0}\right)^{2/3}$  Einstein–de Sitter model (EdS)

$$\Rightarrow H(t) = \frac{2}{3t}; H_0 = H(t_0) = \frac{2}{3t_0}$$

- Lemaître (1927) [ADS:1927ASSB...47...49L](#):  $H_0 \approx 0.6 \text{ Gyr}^{-1}$

- Hubble (1929) [ADS:1929PNAS...15..168H](#):  $H_0 \approx 0.5 \text{ Gyr}^{-1}$

$$\Rightarrow \text{EdS would give } t_0 = \frac{2}{3H_0} \approx 1.3 \text{ Gyr}$$

# FLRW matter-dominated epoch

- Friedmann Eqn:

$$\dot{a}^2 = \frac{8\pi G \rho_{m0}}{3a} - c^2 k$$

- matter-dominated epoch:  $\rho = \rho_m = \rho_{m0} a^{-3}$

- $k = 0$  case:  $a = \left(\frac{t}{t_0}\right)^{2/3}$  Einstein–de Sitter model (EdS)

$$\Rightarrow H(t) = \frac{2}{3t}; H_0 = H(t_0) = \frac{2}{3t_0}$$

- Lemaître (1927) [ADS:1927ASSB...47...49L](#):  $H_0 \approx 0.6 \text{ Gyr}^{-1}$

- Hubble (1929) [ADS:1929PNAS...15..168H](#):  $H_0 \approx 0.5 \text{ Gyr}^{-1}$

$$\Rightarrow \text{EdS would give } t_0 = \frac{2}{3H_0} \approx 1.3 \text{ Gyr} < t_{\text{Earth}} \approx 4.5 \text{ Gyr}$$

# FLRW matter-dominated epoch

- Friedmann Eqn:

$$\dot{a}^2 = \frac{8\pi G \rho_{m0}}{3a} - c^2 k$$

- matter-dominated epoch:  $\rho = \rho_m = \rho_{m0} a^{-3}$

- $k = 0$  case:  $a = \left(\frac{t}{t_0}\right)^{2/3}$  Einstein–de Sitter model (EdS)

$$\Rightarrow H(t) = \frac{2}{3t}; H_0 = H(t_0) = \frac{2}{3t_0}$$

- Lemaître (1927) [ADS:1927ASSB...47...49L](#):  $H_0 \approx 0.6 \text{ Gyr}^{-1}$

- Hubble (1929) [ADS:1929PNAS...15..168H](#):  $H_0 \approx 0.5 \text{ Gyr}^{-1}$

$$\Rightarrow \text{EdS would give } t_0 = \frac{2}{3H_0} \approx 1.3 \text{ Gyr} < t_{\text{Earth}} \approx 4.5 \text{ Gyr}$$

- 1980's:  $H_0 \approx 50$  or  $100 \text{ km/s/Mpc}$

# FLRW matter-dominated epoch

- Friedmann Eqn:

$$\dot{a}^2 = \frac{8\pi G \rho_{m0}}{3a} - c^2 k$$

- matter-dominated epoch:  $\rho = \rho_m = \rho_{m0} a^{-3}$

- $k = 0$  case:  $a = \left(\frac{t}{t_0}\right)^{2/3}$  Einstein–de Sitter model (EdS)

$$\Rightarrow H(t) = \frac{2}{3t}; H_0 = H(t_0) = \frac{2}{3t_0}$$

- Lemaître (1927) [ADS:1927ASSB...47...49L](#):  $H_0 \approx 0.6 \text{ Gyr}^{-1}$

- Hubble (1929) [ADS:1929PNAS...15..168H](#):  $H_0 \approx 0.5 \text{ Gyr}^{-1}$

$$\Rightarrow \text{EdS would give } t_0 = \frac{2}{3H_0} \approx 1.3 \text{ Gyr} < t_{\text{Earth}} \approx 4.5 \text{ Gyr}$$

- 1980's:  $H_0 \approx 50$  or  $100 \text{ kpc/Gyr/Mpc}$

# FLRW matter-dominated epoch

- Friedmann Eqn:

$$\dot{a}^2 = \frac{8\pi G \rho_{m0}}{3a} - c^2 k$$

- matter-dominated epoch:  $\rho = \rho_m = \rho_{m0} a^{-3}$

- $k = 0$  case:  $a = \left(\frac{t}{t_0}\right)^{2/3}$  Einstein–de Sitter model (EdS)

$$\Rightarrow H(t) = \frac{2}{3t}; H_0 = H(t_0) = \frac{2}{3t_0}$$

- Lemaître (1927) [ADS:1927ASSB...47...49L](#):  $H_0 \approx 0.6 \text{ Gyr}^{-1}$

- Hubble (1929) [ADS:1929PNAS...15..168H](#):  $H_0 \approx 0.5 \text{ Gyr}^{-1}$

$$\Rightarrow \text{EdS would give } t_0 = \frac{2}{3H_0} \approx 1.3 \text{ Gyr} < t_{\text{Earth}} \approx 4.5 \text{ Gyr}$$

- 1980's:  $H_0 \approx 0.05$  or  $0.1 \text{ Gyr}^{-1}$

# FLRW matter-dominated epoch

- Friedmann Eqn:

$$\dot{a}^2 = \frac{8\pi G \rho_{m0}}{3a} - c^2 k$$

- matter-dominated epoch:  $\rho = \rho_m = \rho_{m0} a^{-3}$

- $k = 0$  case:  $a = \left(\frac{t}{t_0}\right)^{2/3}$  Einstein–de Sitter model (EdS)

$$\Rightarrow H(t) = \frac{2}{3t}; H_0 = H(t_0) = \frac{2}{3t_0}$$

- Lemaître (1927) [ADS:1927ASSB...47...49L](#):  $H_0 \approx 0.6 \text{ Gyr}^{-1}$

- Hubble (1929) [ADS:1929PNAS...15..168H](#):  $H_0 \approx 0.5 \text{ Gyr}^{-1}$

$$\Rightarrow \text{EdS would give } t_0 = \frac{2}{3H_0} \approx 1.3 \text{ Gyr} < t_{\text{Earth}} \approx 4.5 \text{ Gyr}$$

- 1980's:  $H_0 \approx 0.05$  or  $0.1 \text{ Gyr}^{-1} \Rightarrow t_0(\text{EdS}) \approx 13.0$  or  $6.5 \text{ Gyr}$ , resp.



# FLRW: $\rho_{\text{crit}}$

■ Friedmann Eqn:

$$\frac{\dot{a}^2}{a^2} = \frac{8\pi G\rho}{3} - \frac{c^2 k}{a^2}$$

# FLRW: $\rho_{\text{crit}}$

■ Friedmann Eqn:

$$H^2 = \frac{8\pi G\rho}{3} - \frac{c^2 k}{a^2}$$

# FLRW: $\rho_{\text{crit}}$

- Friedmann Eqn:

$$H^2 = \frac{8\pi G\rho}{3} - \frac{c^2 k}{a^2}$$

- consider a fixed observation, e.g.  $H_0 = 100 \text{ km/s/Mpc}$

# FLRW: $\rho_{\text{crit}}$

- Friedmann Eqn:

$$H^2 = \frac{8\pi G\rho}{3} - \frac{c^2 k}{a^2}$$

- consider a fixed observation, e.g.  $H_0 = 100 \text{ km/s/Mpc}$

- ◆  $\rho_{\text{m}0} = \frac{3H_0^2}{8\pi G} \Leftrightarrow k = 0$

# FLRW: $\rho_{\text{crit}}$

- Friedmann Eqn:

$$H^2 = \frac{8\pi G\rho}{3} - \frac{c^2 k}{a^2}$$

- consider a fixed observation, e.g.  $H_0 = 100 \text{ km/s/Mpc}$

- ◆  $\rho_{\text{m}0} = \frac{3H_0^2}{8\pi G} \Leftrightarrow k = 0 \text{ flat}$

# FLRW: $\rho_{\text{crit}}$

- Friedmann Eqn:

$$H^2 = \frac{8\pi G\rho}{3} - \frac{c^2 k}{a^2}$$

- consider a fixed observation, e.g.  $H_0 = 100 \text{ km/s/Mpc}$

- ◆  $\rho_{\text{m}0} = \frac{3H_0^2}{8\pi G} \Leftrightarrow k = 0 \text{ flat}$

- ◆  $\rho_{\text{m}0} > \frac{3H_0^2}{8\pi G} \Leftrightarrow k > 0$

# FLRW: $\rho_{\text{crit}}$

- Friedmann Eqn:

$$H^2 = \frac{8\pi G\rho}{3} - \frac{c^2 k}{a^2}$$

- consider a fixed observation, e.g.  $H_0 = 100 \text{ km/s/Mpc}$

- ◆  $\rho_{\text{m}0} = \frac{3H_0^2}{8\pi G} \Leftrightarrow k = 0 \text{ flat}$

- ◆  $\rho_{\text{m}0} > \frac{3H_0^2}{8\pi G} \Leftrightarrow k > 0 \text{ spherical}$

# FLRW: $\rho_{\text{crit}}$

- Friedmann Eqn:

$$H^2 = \frac{8\pi G\rho}{3} - \frac{c^2 k}{a^2}$$

- consider a fixed observation, e.g.  $H_0 = 100 \text{ km/s/Mpc}$

- ◆  $\rho_{\text{m}0} = \frac{3H_0^2}{8\pi G} \Leftrightarrow k = 0 \text{ flat}$

- ◆  $\rho_{\text{m}0} > \frac{3H_0^2}{8\pi G} \Leftrightarrow k > 0 \text{ spherical}$

- ◆  $\rho_{\text{m}0} < \frac{3H_0^2}{8\pi G} \Leftrightarrow k < 0$



# FLRW: $\rho_{\text{crit}}$

- Friedmann Eqn:

$$H^2 = \frac{8\pi G \rho}{3} - \frac{c^2 k}{a^2}$$

- consider a fixed observation, e.g.  $H_0 = 100 \text{ km/s/Mpc}$

- ◆  $\rho_{\text{m}0} = \frac{3H_0^2}{8\pi G} \Leftrightarrow k = 0 \text{ flat}$
- ◆  $\rho_{\text{m}0} > \frac{3H_0^2}{8\pi G} \Leftrightarrow k > 0 \text{ spherical}$
- ◆  $\rho_{\text{m}0} < \frac{3H_0^2}{8\pi G} \Leftrightarrow k < 0 \text{ hyperbolic}$

# FLRW: $\rho_{\text{crit}}$

■ Friedmann Eqn:

$$H^2 = \frac{8\pi G\rho}{3} - \frac{c^2 k}{a^2}$$

# FLRW: $\rho_{\text{crit}}$

■ Friedmann Eqn:

$$H^2 = \frac{8\pi G\rho}{3} - \frac{c^2 k}{a^2}$$

Defn:  $\rho_{\text{crit}} := \frac{3H^2}{8\pi G}$  critical density

# FLRW: $\rho_{\text{crit}}$

■ Friedmann Eqn: 
$$H^2 = \frac{\rho H^2}{\rho_{\text{crit}}} - \frac{c^2 k}{a^2}$$

Defn:  $\rho_{\text{crit}} := \frac{3H^2}{8\pi G}$  critical density

# FLRW: $\rho_{\text{crit}}$

■ Friedmann Eqn: 
$$H^2 = \frac{\rho H^2}{\rho_{\text{crit}}} - \frac{c^2 k}{a^2}$$

Defn:  $\rho_{\text{crit}} := \frac{3H^2}{8\pi G}$  critical density

Defn:  $\Omega_{\text{m}} := \frac{\rho}{\rho_{\text{crit}}}$  matter density parameter

# FLRW: $\rho_{\text{crit}}$

■ Friedmann Eqn: 
$$H^2 = \Omega_m H^2 - \frac{c^2 k}{a^2}$$

Defn:  $\rho_{\text{crit}} := \frac{3H^2}{8\pi G}$  critical density

Defn:  $\Omega_m := \frac{\rho}{\rho_{\text{crit}}}$  matter density parameter

# FLRW: $\rho_{\text{crit}}$

■ Friedmann Eqn: 
$$H^2 = \Omega_m H^2 - \frac{c^2 k}{a^2}$$

Defn:  $\rho_{\text{crit}} := \frac{3H^2}{8\pi G}$  critical density

Defn:  $\Omega_m := \frac{\rho}{\rho_{\text{crit}}}$  matter density parameter

Defn:  $\Omega_k := -\frac{c^2 k}{a^2 H^2}$  curvature density parameter

# FLRW: $\rho_{\text{crit}}$

■ Friedmann Eqn: 
$$H^2 = \Omega_m H^2 - \frac{c^2 k}{a^2}$$

Defn:  $\rho_{\text{crit}} := \frac{3H^2}{8\pi G}$  critical density

Defn:  $\Omega_m := \frac{\rho}{\rho_{\text{crit}}}$  matter density parameter

Defn:  $\Omega_k := -\frac{c^2 k}{a^2 H^2}$  curvature density parameter (sign reversal!)



# FLRW: $\rho_{\text{crit}}$

■ Friedmann Eqn:  $H^2 = \Omega_m H^2 + \Omega_k H^2$

Defn:  $\rho_{\text{crit}} := \frac{3H^2}{8\pi G}$  critical density

Defn:  $\Omega_m := \frac{\rho}{\rho_{\text{crit}}}$  matter density parameter

Defn:  $\Omega_k := -\frac{c^2 k}{a^2 H^2}$  curvature density parameter (sign reversal!)

# FLRW: $\rho_{\text{crit}}$

■ Friedmann Eqn:

$$1 = \Omega_m + \Omega_k$$

Defn:  $\rho_{\text{crit}} := \frac{3H^2}{8\pi G}$  critical density

Defn:  $\Omega_m := \frac{\rho}{\rho_{\text{crit}}}$  matter density parameter

Defn:  $\Omega_k := -\frac{c^2 k}{a^2 H^2}$  curvature density parameter (sign reversal!)

# FLRW: $\rho_{\text{crit}}$

■ Friedmann Eqn:

$$1 = \Omega_m + \Omega_k$$

Defn:  $\rho_{\text{crit}} := \frac{3H^2}{8\pi G}$  critical density

Defn:  $\Omega_m := \frac{\rho}{\rho_{\text{crit}}}$  matter density parameter

Defn:  $\Omega_k := -\frac{c^2 k}{a^2 H^2}$  curvature density parameter (sign reversal!)

■ consider a fixed observation, e.g.  $H_0 = 100 \text{ km/s/Mpc}$

# FLRW: $\rho_{\text{crit}}$

■ Friedmann Eqn:

$$1 = \Omega_m + \Omega_k$$

Defn:  $\rho_{\text{crit}} := \frac{3H^2}{8\pi G}$  critical density

Defn:  $\Omega_m := \frac{\rho}{\rho_{\text{crit}}}$  matter density parameter

Defn:  $\Omega_k := -\frac{c^2 k}{a^2 H^2}$  curvature density parameter (sign reversal!)

■ consider a fixed observation, e.g.  $H_0 = 100 \text{ km/s/Mpc}$

◆  $\Omega_{m0} = 1 \Leftrightarrow k = 0$

# FLRW: $\rho_{\text{crit}}$

■ Friedmann Eqn:

$$1 = \Omega_m + \Omega_k$$

Defn:  $\rho_{\text{crit}} := \frac{3H^2}{8\pi G}$  critical density

Defn:  $\Omega_m := \frac{\rho}{\rho_{\text{crit}}}$  matter density parameter

Defn:  $\Omega_k := -\frac{c^2 k}{a^2 H^2}$  curvature density parameter (sign reversal!)

■ consider a fixed observation, e.g.  $H_0 = 100 \text{ km/s/Mpc}$

◆  $\Omega_{m0} = 1 \Leftrightarrow k = 0$  flat

# FLRW: $\rho_{\text{crit}}$

■ Friedmann Eqn:

$$1 = \Omega_m + \Omega_k$$

Defn:  $\rho_{\text{crit}} := \frac{3H^2}{8\pi G}$  critical density

Defn:  $\Omega_m := \frac{\rho}{\rho_{\text{crit}}}$  matter density parameter

Defn:  $\Omega_k := -\frac{c^2 k}{a^2 H^2}$  curvature density parameter (sign reversal!)

■ consider a fixed observation, e.g.  $H_0 = 100 \text{ km/s/Mpc}$

◆  $\Omega_{m0} = 1 \Leftrightarrow k = 0$  flat

◆  $\Omega_{m0} > 1 \Leftrightarrow k > 0$

# FLRW: $\rho_{\text{crit}}$

■ Friedmann Eqn:

$$1 = \Omega_m + \Omega_k$$

Defn:  $\rho_{\text{crit}} := \frac{3H^2}{8\pi G}$  critical density

Defn:  $\Omega_m := \frac{\rho}{\rho_{\text{crit}}}$  matter density parameter

Defn:  $\Omega_k := -\frac{c^2 k}{a^2 H^2}$  curvature density parameter (sign reversal!)

■ consider a fixed observation, e.g.  $H_0 = 100 \text{ km/s/Mpc}$

◆  $\Omega_{m0} = 1 \Leftrightarrow k = 0$  flat

◆  $\Omega_{m0} > 1 \Leftrightarrow k > 0$  spherical

# FLRW: $\rho_{\text{crit}}$

■ Friedmann Eqn:

$$1 = \Omega_m + \Omega_k$$

Defn:  $\rho_{\text{crit}} := \frac{3H^2}{8\pi G}$  critical density

Defn:  $\Omega_m := \frac{\rho}{\rho_{\text{crit}}}$  matter density parameter

Defn:  $\Omega_k := -\frac{c^2 k}{a^2 H^2}$  curvature density parameter (sign reversal!)

■ consider a fixed observation, e.g.  $H_0 = 100 \text{ km/s/Mpc}$

- ◆  $\Omega_{m0} = 1 \Leftrightarrow k = 0$  flat
- ◆  $\Omega_{m0} > 1 \Leftrightarrow k > 0$  spherical
- ◆  $\Omega_{m0} < 1 \Leftrightarrow k < 0$



# FLRW: $\rho_{\text{crit}}$

■ Friedmann Eqn:

$$1 = \Omega_m + \Omega_k$$

Defn:  $\rho_{\text{crit}} := \frac{3H^2}{8\pi G}$  critical density

Defn:  $\Omega_m := \frac{\rho}{\rho_{\text{crit}}}$  matter density parameter

Defn:  $\Omega_k := -\frac{c^2 k}{a^2 H^2}$  curvature density parameter (sign reversal!)

■ consider a fixed observation, e.g.  $H_0 = 100 \text{ km/s/Mpc}$

- ◆  $\Omega_{m0} = 1 \Leftrightarrow k = 0$  flat
- ◆  $\Omega_{m0} > 1 \Leftrightarrow k > 0$  spherical
- ◆  $\Omega_{m0} < 1 \Leftrightarrow k < 0$  hyperbolic

# FLRW: $\rho_{\text{crit}}$

■ Friedmann Eqn:

$$1 = \Omega_{\text{tot}} + \Omega_k$$

Defn:  $\rho_{\text{crit}} := \frac{3H^2}{8\pi G}$  critical density

Defn:  $\Omega_m := \frac{\rho}{\rho_{\text{crit}}}$  matter density parameter

Defn:  $\Omega_k := -\frac{c^2 k}{a^2 H^2}$  curvature density parameter (sign reversal!)

■ consider a fixed observation, e.g.  $H_0 = 100 \text{ km/s/Mpc}$

- ◆  $\Omega_{m0} = 1 \Leftrightarrow k = 0$  flat
- ◆  $\Omega_{m0} > 1 \Leftrightarrow k > 0$  spherical
- ◆  $\Omega_{m0} < 1 \Leftrightarrow k < 0$  hyperbolic

■  $\Omega_{\text{tot}} :=$

# FLRW: $\rho_{\text{crit}}$

■ Friedmann Eqn:

$$1 = \Omega_{\text{tot}} + \Omega_k$$

Defn:  $\rho_{\text{crit}} := \frac{3H^2}{8\pi G}$  critical density

Defn:  $\Omega_m := \frac{\rho}{\rho_{\text{crit}}}$  matter density parameter

Defn:  $\Omega_k := -\frac{c^2 k}{a^2 H^2}$  curvature density parameter (sign reversal!)

■ consider a fixed observation, e.g.  $H_0 = 100 \text{ km/s/Mpc}$

- ◆  $\Omega_{m0} = 1 \Leftrightarrow k = 0$  flat
- ◆  $\Omega_{m0} > 1 \Leftrightarrow k > 0$  spherical
- ◆  $\Omega_{m0} < 1 \Leftrightarrow k < 0$  hyperbolic

■  $\Omega_{\text{tot}} := \Omega_m +$

# FLRW: $\rho_{\text{crit}}$

■ Friedmann Eqn:

$$1 = \Omega_{\text{tot}} + \Omega_k$$

Defn:  $\rho_{\text{crit}} := \frac{3H^2}{8\pi G}$  critical density

Defn:  $\Omega_{\text{m}} := \frac{\rho}{\rho_{\text{crit}}}$  matter density parameter

Defn:  $\Omega_k := -\frac{c^2 k}{a^2 H^2}$  curvature density parameter (sign reversal!)

■ consider a fixed observation, e.g.  $H_0 = 100 \text{ km/s/Mpc}$

- ◆  $\Omega_{\text{m}0} = 1 \Leftrightarrow k = 0$  flat
- ◆  $\Omega_{\text{m}0} > 1 \Leftrightarrow k > 0$  spherical
- ◆  $\Omega_{\text{m}0} < 1 \Leftrightarrow k < 0$  hyperbolic

■  $\Omega_{\text{tot}} := \Omega_{\text{m}} + \Omega_{\text{r}} +$

# FLRW: $\rho_{\text{crit}}$

■ Friedmann Eqn:

$$1 = \Omega_{\text{tot}} + \Omega_k$$

Defn:  $\rho_{\text{crit}} := \frac{3H^2}{8\pi G}$  critical density

Defn:  $\Omega_m := \frac{\rho}{\rho_{\text{crit}}}$  matter density parameter

Defn:  $\Omega_k := -\frac{c^2 k}{a^2 H^2}$  curvature density parameter (sign reversal!)

■ consider a fixed observation, e.g.  $H_0 = 100 \text{ km/s/Mpc}$

- ◆  $\Omega_{m0} = 1 \Leftrightarrow k = 0$  flat
- ◆  $\Omega_{m0} > 1 \Leftrightarrow k > 0$  spherical
- ◆  $\Omega_{m0} < 1 \Leftrightarrow k < 0$  hyperbolic

■  $\Omega_{\text{tot}} := \Omega_m + \Omega_r + \Omega_\Lambda$

# FLRW: $\rho_{\text{crit}}$

■ Friedmann Eqn:

$$1 = \Omega_{\text{tot}} + \Omega_k$$

Defn:  $\rho_{\text{crit}} := \frac{3H^2}{8\pi G}$  critical density

Defn:  $\Omega_m := \frac{\rho}{\rho_{\text{crit}}}$  matter density parameter

Defn:  $\Omega_k := -\frac{c^2 k}{a^2 H^2}$  curvature density parameter (sign reversal!)

■ consider a fixed observation, e.g.  $H_0 = 100 \text{ km/s/Mpc}$

- ◆  $\Omega_{m0} = 1 \Leftrightarrow k = 0$  flat
- ◆  $\Omega_{m0} > 1 \Leftrightarrow k > 0$  spherical
- ◆  $\Omega_{m0} < 1 \Leftrightarrow k < 0$  hyperbolic

■  $\Omega_{\text{tot}} := \Omega_b + \Omega_{\text{nbDM}} + \Omega_r + \Omega_\Lambda$

# FLRW curvature constant

■ metric in

◆ azimuthal equidistant coords:  $R_C$

# FLRW curvature constant

■ metric in

- ◆ azimuthal equidistant coords:  $R_C$
- ◆ orthographic coords:  $k$



# FLRW curvature constant

■ metric in

◆ azimuthal equidistant coords:  $R_C$

◆ orthographic coords:  $k$

■ orthographic:  $1 - kr^2 = 0$  coord singularity at equator

# FLRW curvature constant

■ metric in

◆ azimuthal equidistant coords:  $R_C$

◆ orthographic coords:  $k$

■ orthographic:  $1 - kr^2 = 0$  coord singularity at equator

■  $\Rightarrow kR_C^2 = 1$

# FLRW curvature constant

■ metric in

◆ azimuthal equidistant coords:  $R_C$

◆ orthographic coords:  $k$

■ orthographic:  $1 - kr^2 = 0$  coord singularity at equator

■  $\Rightarrow kR_C^2 = 1 \Rightarrow k = 1/R_C^2$

# FLRW curvature constant

■ metric in

◆ azimuthal equidistant coords:  $R_C$

◆ orthographic coords:  $k$

■ orthographic:  $1 - kr^2 = 0$  coord singularity at equator

■  $\Rightarrow kR_C^2 = 1 \Rightarrow k = 1/R_C^2$

■ Defn:  $\Omega_k := -\frac{c^2 k}{a^2 H^2}$

# FLRW curvature constant

■ metric in

◆ azimuthal equidistant coords:  $R_C$

◆ orthographic coords:  $k$

■ orthographic:  $1 - kr^2 = 0$  coord singularity at equator

■  $\Rightarrow kR_C^2 = 1 \Rightarrow k = 1/R_C^2$

■ Defn:  $\Omega_k := -\frac{c^2 k}{a^2 H^2}$   $\Rightarrow \Omega_{k0} = -\frac{c^2 k}{H_0^2}$

# FLRW curvature constant

■ metric in

◆ azimuthal equidistant coords:  $R_C$

◆ orthographic coords:  $k$

■ orthographic:  $1 - kr^2 = 0$  coord singularity at equator

■  $\Rightarrow kR_C^2 = 1 \Rightarrow k = 1/R_C^2$

■ Defn:  $\Omega_k := -\frac{c^2 k}{a^2 H^2}$   $\Rightarrow k = -\frac{\Omega_{k0} H_0^2}{c^2}$

# FLRW curvature constant

■ metric in

◆ azimuthal equidistant coords:  $R_C$

◆ orthographic coords:  $k$

■ orthographic:  $1 - kr^2 = 0$  coord singularity at equator

■  $\Rightarrow kR_C^2 = 1 \Rightarrow k = 1/R_C^2$

■ Defn:  $\Omega_k := -\frac{c^2 k}{a^2 H^2}$   $\Rightarrow R_C^2 = -\frac{c^2}{\Omega_{k0} H_0^2}$

# FLRW curvature constant

■ metric in

◆ azimuthal equidistant coords:  $R_C$

◆ orthographic coords:  $k$

■ orthographic:  $1 - kr^2 = 0$  coord singularity at equator

■  $\Rightarrow kR_C^2 = 1 \Rightarrow k = 1/R_C^2$

■ Defn:  $\Omega_k := -\frac{c^2 k}{a^2 H^2} \Rightarrow R_C^2 = -\frac{c^2}{H_0^2} \frac{1}{\Omega_{k0}}$



# FLRW curvature constant

■ metric in

◆ azimuthal equidistant coords:  $R_C$

◆ orthographic coords:  $k$

■ orthographic:  $1 - kr^2 = 0$  coord singularity at equator

■  $\Rightarrow kR_C^2 = 1 \Rightarrow k = 1/R_C^2$

■ Defn:  $\Omega_k := -\frac{c^2 k}{a^2 H^2}$   $\Rightarrow R_C^2 = -\frac{c^2}{H_0^2} \frac{1}{1 - \Omega_{\text{tot}0}}$

# FLRW curvature constant

■ metric in

◆ azimuthal equidistant coords:  $R_C$

◆ orthographic coords:  $k$

■ orthographic:  $1 - kr^2 = 0$  coord singularity at equator

■  $\Rightarrow kR_C^2 = 1 \Rightarrow k = 1/R_C^2$

■ Defn:  $\Omega_k := -\frac{c^2 k}{a^2 H^2}$   $\Rightarrow$   $R_C = \frac{c}{H_0} \frac{1}{\sqrt{\Omega_{\text{tot}0} - 1}}$

# FLRW curvature constant

■ metric in

◆ azimuthal equidistant coords:  $R_C$

◆ orthographic coords:  $k$

■ orthographic:  $1 - kr^2 = 0$  coord singularity at equator

■  $\Rightarrow kR_C^2 = 1 \Rightarrow k = 1/R_C^2$

■ Defn:  $\Omega_k := -\frac{c^2 k}{a^2 H^2}$   $\Rightarrow$   $R_C = \frac{c}{H_0} \frac{1}{\sqrt{\Omega_{\text{tot}0} - 1}}$

■  $\Omega_{\text{tot}0} > 1$

# FLRW curvature constant

■ metric in

◆ azimuthal equidistant coords:  $R_C$

◆ orthographic coords:  $k$

■ orthographic:  $1 - kr^2 = 0$  coord singularity at equator

■  $\Rightarrow kR_C^2 = 1 \Rightarrow k = 1/R_C^2$

■ Defn:  $\Omega_k := -\frac{c^2 k}{a^2 H^2}$   $\Rightarrow$   $R_C = \frac{c}{H_0} \frac{1}{\sqrt{\Omega_{\text{tot}0} - 1}}$

■  $\Omega_{\text{tot}0} > 1$  *spherical*

# FLRW curvature constant

■ metric in

◆ azimuthal equidistant coords:  $R_C$

◆ orthographic coords:  $k$

■ orthographic:  $1 - kr^2 = 0$  coord singularity at equator

■  $\Rightarrow kR_C^2 = 1 \Rightarrow k = 1/R_C^2$

■ Defn:  $\Omega_k := -\frac{c^2 k}{a^2 H^2}$   $\Rightarrow$   $R_C = \frac{c}{H_0} \frac{1}{\sqrt{\Omega_{\text{tot}0} - 1}}$

■  $\Omega_{\text{tot}0} > 1$  *spherical*  $R_C$  real

# FLRW curvature constant

■ metric in

◆ azimuthal equidistant coords:  $R_C$

◆ orthographic coords:  $k$

■ orthographic:  $1 - kr^2 = 0$  coord singularity at equator

■  $\Rightarrow kR_C^2 = 1 \Rightarrow k = 1/R_C^2$

■ Defn:  $\Omega_k := -\frac{c^2 k}{a^2 H^2}$   $\Rightarrow$   $R_C = \frac{c}{H_0} \frac{1}{\sqrt{\Omega_{\text{tot}0} - 1}}$

■  $\Omega_{\text{tot}0} > 1$  *spherical*  $R_C$  real

■  $\Omega_{\text{tot}0} = 1$

# FLRW curvature constant

■ metric in

◆ azimuthal equidistant coords:  $R_C$

◆ orthographic coords:  $k$

■ orthographic:  $1 - kr^2 = 0$  coord singularity at equator

■  $\Rightarrow kR_C^2 = 1 \Rightarrow k = 1/R_C^2$

■ Defn:  $\Omega_k := -\frac{c^2 k}{a^2 H^2}$   $\Rightarrow$   $R_C = \frac{c}{H_0} \frac{1}{\sqrt{\Omega_{\text{tot}0} - 1}}$

■  $\Omega_{\text{tot}0} > 1$  *spherical*  $R_C$  real

■  $\Omega_{\text{tot}0} = 1$  *flat*

# FLRW curvature constant

■ metric in

◆ azimuthal equidistant coords:  $R_C$

◆ orthographic coords:  $k$

■ orthographic:  $1 - kr^2 = 0$  coord singularity at equator

■  $\Rightarrow kR_C^2 = 1 \Rightarrow k = 1/R_C^2$

■ Defn:  $\Omega_k := -\frac{c^2 k}{a^2 H^2}$   $\Rightarrow$   $R_C = \frac{c}{H_0} \frac{1}{\sqrt{\Omega_{\text{tot}0} - 1}}$

■  $\Omega_{\text{tot}0} > 1$  *spherical*  $R_C$  real

■  $\Omega_{\text{tot}0} = 1$  *flat*  $R_C$  undefined



# FLRW curvature constant

■ metric in

◆ azimuthal equidistant coords:  $R_C$

◆ orthographic coords:  $k$

■ orthographic:  $1 - kr^2 = 0$  coord singularity at equator

■  $\Rightarrow kR_C^2 = 1 \Rightarrow k = 1/R_C^2$

■ Defn:  $\Omega_k := -\frac{c^2 k}{a^2 H^2}$   $\Rightarrow$   $R_C = \frac{c}{H_0} \frac{1}{\sqrt{\Omega_{\text{tot}0} - 1}}$

■  $\Omega_{\text{tot}0} > 1$  *spherical*  $R_C$  real

■  $\Omega_{\text{tot}0} = 1$  *flat*  $R_C$  undefined

■  $\Omega_{\text{tot}0} < 1$

# FLRW curvature constant

■ metric in

◆ azimuthal equidistant coords:  $R_C$

◆ orthographic coords:  $k$

■ orthographic:  $1 - kr^2 = 0$  coord singularity at equator

■  $\Rightarrow kR_C^2 = 1 \Rightarrow k = 1/R_C^2$

■ Defn:  $\Omega_k := -\frac{c^2 k}{a^2 H^2}$   $\Rightarrow$   $R_C = \frac{c}{H_0} \frac{1}{\sqrt{\Omega_{\text{tot}0} - 1}}$

■  $\Omega_{\text{tot}0} > 1$  *spherical*  $R_C$  real

■  $\Omega_{\text{tot}0} = 1$  *flat*  $R_C$  undefined

■  $\Omega_{\text{tot}0} < 1$  *hyperbolic*

# FLRW curvature constant

■ metric in

◆ azimuthal equidistant coords:  $R_C$

◆ orthographic coords:  $k$

■ orthographic:  $1 - kr^2 = 0$  coord singularity at equator

■  $\Rightarrow kR_C^2 = 1 \Rightarrow k = 1/R_C^2$

■ Defn:  $\Omega_k := -\frac{c^2 k}{a^2 H^2}$   $\Rightarrow$   $R_C = \frac{c}{H_0} \frac{1}{\sqrt{\Omega_{\text{tot}0} - 1}}$

■  $\Omega_{\text{tot}0} > 1$  *spherical*  $R_C$  real

■  $\Omega_{\text{tot}0} = 1$  *flat*  $R_C$  undefined

■  $\Omega_{\text{tot}0} < 1$  *hyperbolic*  $R_C$  imaginary (or use  $|R_C|$ )

# Einstein's free parameter: $\Lambda$

- Einstein: prevent expansion/contraction via  $\Lambda$   
ADS:1917SPAW.....142E

# Einstein's free parameter: $\Lambda$

- Einstein: prevent expansion/contraction via  $\Lambda$   
[ADS:1917SPAW.....142E](#)
- MAXIMA: calculate  $\mathbf{G}$  and  $\mathbf{G} - \mathbf{g}\Lambda = 8\pi\mathbf{T}$  and simplify:  
<http://cosmo.torun.pl/Cosmo/FLRWEquationsGR>

# Einstein's free parameter: $\Lambda$

- Einstein: prevent expansion/contraction via  $\Lambda$   
[ADS:1917SPAW.....142E](#)
- MAXIMA: calculate  $\mathbf{G}$  and  $\mathbf{G} - \mathbf{g}\Lambda = 8\pi\mathbf{T}$  and simplify:  
<http://cosmo.torun.pl/Cosmo/FLRWEquationsGR>
- *hint*: mixed index form of  $\mathbf{g}$  is easy

# Einstein's free parameter: $\Lambda$

- Einstein: prevent expansion/contraction via  $\Lambda$   
ADS:1917SPAW.....142E

Friedmann Eqn ( $\Lambda \neq 0$ ):

$$\frac{c^2 k}{a^2} + \frac{\dot{a}^2}{a^2} = \frac{8 \pi G \rho}{3} + \frac{c^2 \Lambda}{3}$$

# Einstein's free parameter: $\Lambda$

- Einstein: prevent expansion/contraction via  $\Lambda$   
ADS:1917SPAW.....142E

Friedmann Eqn ( $\Lambda \neq 0$ ):

$$\frac{c^2 k}{a^2} + \frac{\dot{a}^2}{a^2} = \frac{8 \pi G \rho}{3} + \frac{c^2 \Lambda}{3}$$

acceleration Eqn ( $\Lambda \neq 0$ ):

$$\frac{\ddot{a}}{a} = -\frac{4 \pi G (\rho + 3 p/c^2)}{3} + \frac{c^2 \Lambda}{3}$$



# Einstein's free parameter: $\Lambda$

- Einstein: prevent expansion/contraction via  $\Lambda$   
ADS:1917SPAW.....142E

Friedmann Eqn ( $\Lambda \neq 0$ ):

$$\frac{c^2 k}{a^2} + \frac{\dot{a}^2}{a^2} = \frac{8 \pi G \rho}{3} + \frac{c^2 \Lambda}{3}$$

acceleration Eqn ( $\Lambda \neq 0$ ):

$$\frac{\ddot{a}}{a} = -\frac{4 \pi G (\rho + 3 p/c^2)}{3} + \frac{c^2 \Lambda}{3}$$

Defn: “dust solution”:  $p(t) = 0$

# Einstein's free parameter: $\Lambda$

- Einstein: prevent expansion/contraction via  $\Lambda$   
ADS:1917SPAW.....142E

Friedmann Eqn ( $\Lambda \neq 0$ ):

$$\frac{c^2 k}{a^2} + \frac{\dot{a}^2}{a^2} = \frac{8 \pi G \rho}{3} + \frac{c^2 \Lambda}{3}$$

acceleration Eqn ( $\Lambda \neq 0$ ):

$$\frac{\ddot{a}}{a} = -\frac{4 \pi G (\rho + 3 p/c^2)}{3} + \frac{c^2 \Lambda}{3}$$

Defn: “dust solution”:  $p(t) = 0$

Defn:  $\Omega_\Lambda := \frac{c^2 \Lambda}{3 H^2}$

# Einstein's free parameter: $\Lambda$

- Einstein: prevent expansion/contraction via  $\Lambda$   
ADS:1917SPAW.....142E

Friedmann Eqn ( $\Lambda \neq 0$ ):

$$\frac{c^2 k}{a^2} + \frac{\dot{a}^2}{a^2} = \frac{8 \pi G \rho}{3} + \frac{c^2 \Lambda}{3}$$

Defn: "dust solution":  $p(t) = 0$

Defn:  $\Omega_\Lambda := \frac{c^2 \Lambda}{3 H^2}$

acceleration Eqn ( $\Lambda \neq 0$ ):

$$\frac{\ddot{a}}{a} = -\frac{4 \pi G (\rho + 3 p/c^2)}{3} + \Omega_\Lambda H^2$$

# Einstein's free parameter: $\Lambda$

- Einstein: prevent expansion/contraction via  $\Lambda$   
ADS:1917SPAW.....142E

Friedmann Eqn ( $\Lambda \neq 0$ ):

$$\frac{c^2 k}{a^2} + \frac{\dot{a}^2}{a^2} = \frac{8 \pi G \rho}{3} + \frac{c^2 \Lambda}{3}$$

Defn: "dust solution":  $p(t) = 0$

Defn:  $\Omega_\Lambda := \frac{c^2 \Lambda}{3 H^2}$

acceleration Eqn ( $\Lambda \neq 0$ ):

$$\frac{\ddot{a}}{a} = -\frac{H^2}{2} \frac{\rho}{\rho_{\text{crit}}} + \Omega_\Lambda H^2$$

# Einstein's free parameter: $\Lambda$

- Einstein: prevent expansion/contraction via  $\Lambda$   
ADS:1917SPAW.....142E

Friedmann Eqn ( $\Lambda \neq 0$ ):

$$\frac{c^2 k}{a^2} + \frac{\dot{a}^2}{a^2} = \frac{8 \pi G \rho}{3} + \frac{c^2 \Lambda}{3}$$

Defn: "dust solution":  $p(t) = 0$

Defn:  $\Omega_\Lambda := \frac{c^2 \Lambda}{3 H^2}$

acceleration Eqn ( $\Lambda \neq 0$ ):

$$\frac{\ddot{a}}{a} = -\frac{H^2 \Omega_m}{2} + \Omega_\Lambda H^2$$

# Einstein's free parameter: $\Lambda$

- Einstein: prevent expansion/contraction via  $\Lambda$   
ADS:1917SPAW.....142E

$$\text{Friedmann Eqn } (\Lambda \neq 0): \quad \frac{c^2 k}{a^2} + \frac{\dot{a}^2}{a^2} = \frac{8 \pi G \rho}{3} + \frac{c^2 \Lambda}{3}$$

Defn: "dust solution":  $p(t) = 0$

$$\text{Defn: } \Omega_\Lambda := \frac{c^2 \Lambda}{3 H^2}$$

$$\text{acceleration Eqn } (\Lambda \neq 0): \quad \frac{\ddot{a} a^2}{a \dot{a}^2} = -\frac{\Omega_m}{2} + \Omega_\Lambda$$

# Einstein's free parameter: $\Lambda$

- Einstein: prevent expansion/contraction via  $\Lambda$   
ADS:1917SPAW.....142E

Friedmann Eqn ( $\Lambda \neq 0$ ):

$$\frac{c^2 k}{a^2} + \frac{\dot{a}^2}{a^2} = \frac{8 \pi G \rho}{3} + \frac{c^2 \Lambda}{3}$$

Defn: "dust solution":  $p(t) = 0$

Defn:  $\Omega_\Lambda := \frac{c^2 \Lambda}{3 H^2}$

acceleration Eqn ( $\Lambda \neq 0$ ):

$$\frac{\ddot{a} a^2}{a \dot{a}^2} = -\frac{\Omega_m}{2} + \Omega_\Lambda$$

Defn:  $q := -\frac{\ddot{a} a}{\dot{a}^2}$

# Einstein's free parameter: $\Lambda$

- Einstein: prevent expansion/contraction via  $\Lambda$   
ADS:1917SPAW.....142E

$$\text{Friedmann Eqn } (\Lambda \neq 0): \quad \frac{c^2 k}{a^2} + \frac{\dot{a}^2}{a^2} = \frac{8 \pi G \rho}{3} + \frac{c^2 \Lambda}{3}$$

Defn: “dust solution”:  $p(t) = 0$

$$\text{Defn: } \Omega_\Lambda := \frac{c^2 \Lambda}{3 H^2}$$

Defn:  $q := -\frac{\ddot{a}a}{\dot{a}^2}$  “deceleration parameter”

$$\text{acceleration Eqn } (\Lambda \neq 0): \quad q = \frac{\Omega_m}{2} - \Omega_\Lambda$$



# Einstein's free parameter: $\Lambda$

- Einstein: prevent expansion/contraction via  $\Lambda$   
ADS:1917SPAW.....142E

Friedmann Eqn ( $\Lambda \neq 0$ ):

$$\frac{c^2 k}{a^2} + \frac{\dot{a}^2}{a^2} = \frac{8 \pi G \rho}{3} + \frac{c^2 \Lambda}{3}$$

Defn: “dust solution”:  $p(t) = 0$

Defn:  $\Omega_\Lambda := \frac{c^2 \Lambda}{3 H^2}$

Defn:  $q := -\frac{\ddot{a}a}{\dot{a}^2}$  “deceleration parameter”

- $q = \frac{\Omega_m}{2} - \Omega_\Lambda$

# Einstein's free parameter: $\Lambda$

- Einstein: prevent expansion/contraction via  $\Lambda$   
ADS:1917SPAW.....142E

Friedmann Eqn ( $\Lambda \neq 0$ ):

$$\frac{c^2 k}{a^2} + \frac{\dot{a}^2}{a^2} = \frac{8 \pi G \rho}{3} + \frac{c^2 \Lambda}{3}$$

Defn: “dust solution”:  $p(t) = 0$

Defn:  $\Omega_\Lambda := \frac{c^2 \Lambda}{3 H^2}$

Defn:  $q := -\frac{\ddot{a}a}{\dot{a}^2}$  “deceleration parameter”

- $\boxed{q = \frac{\Omega_m}{2} - \Omega_\Lambda}$  acceleration equation

- if  $\Lambda = 0$  and  $\Omega_m > 0$  then  $\frac{\ddot{a}}{a} < 0$ , i.e.  $q > 0$

# Einstein's free parameter: $\Lambda$

- Einstein: prevent expansion/contraction via  $\Lambda$   
ADS:1917SPAW.....142E

Friedmann Eqn ( $\Lambda \neq 0$ ):

$$\frac{c^2 k}{a^2} + \frac{\dot{a}^2}{a^2} = \frac{8 \pi G \rho}{3} + \frac{c^2 \Lambda}{3}$$

Defn: “dust solution”:  $p(t) = 0$

Defn:  $\Omega_\Lambda := \frac{c^2 \Lambda}{3 H^2}$

Defn:  $q := -\frac{\ddot{a}a}{\dot{a}^2}$  “deceleration parameter”

- $\boxed{q = \frac{\Omega_m}{2} - \Omega_\Lambda}$  acceleration equation

- if  $\Lambda = 0$  and  $\Omega_m > 0$  then  $\frac{\ddot{a}}{a} < 0$ , i.e.  $q > 0$  *deceleration*

# Einstein's free parameter: $\Lambda$

- Einstein: prevent expansion/contraction via  $\Lambda$   
ADS:1917SPAW.....142E

$$\text{Friedmann Eqn } (\Lambda \neq 0): \quad \frac{c^2 k}{a^2} + \frac{\dot{a}^2}{a^2} = \frac{8 \pi G \rho}{3} + \frac{c^2 \Lambda}{3}$$

Defn: “dust solution”:  $p(t) = 0$

$$\text{Defn: } \Omega_\Lambda := \frac{c^2 \Lambda}{3 H^2}$$

Defn:  $q := -\frac{\ddot{a}a}{\dot{a}^2}$  “deceleration parameter”

- $\boxed{q = \frac{\Omega_m}{2} - \Omega_\Lambda}$  acceleration equation

- if  $\Omega_\Lambda > \Omega_m/2$  then  $\frac{\ddot{a}}{a} > 0$ , i.e.  $q < 0$  *acceleration*

# distances in FLRW cosmology

- azimuthal equidistant  $r$ :

# distances in FLRW cosmology

- azimuthal equidistant  $r$ : proper distance at  $t_0 \equiv$  comoving radial distance  $r = \int_t^{t_0} \frac{c dt'}{a(t')}$

# distances in FLRW cosmology

- azimuthal equidistant  $r$ : proper distance at  $t_0 \equiv$  comoving radial distance  $r = \int_t^{t_0} \frac{c dt'}{a(t')}$
- Friedmann Eq:  $1 = \Omega_m + \Omega_k + \Omega_\Lambda$

# distances in FLRW cosmology

- azimuthal equidistant  $r$ : proper distance at  $t_0 \equiv$  comoving radial distance  $r = \int_t^{t_0} \frac{c dt'}{a(t')}$
- Friedmann Eq:  $1 = \Omega_m + \Omega_k + \Omega_\Lambda$
- $\Omega_m$  =  $\frac{\rho}{\rho_{\text{crit}}}$



# distances in FLRW cosmology

- azimuthal equidistant  $r$ : proper distance at  $t_0 \equiv$  comoving radial distance  $r = \int_t^{t_0} \frac{c dt'}{a(t')}$
- Friedmann Eq:  $1 = \Omega_m + \Omega_k + \Omega_\Lambda$
- $\Omega_m = \frac{\rho}{\rho_{\text{crit}}} = \frac{\rho_0 a^{-3}}{\rho_{\text{crit}0} (H^2/H_0^2)}$

# distances in FLRW cosmology

- azimuthal equidistant  $r$ : proper distance at  $t_0 \equiv$  comoving radial distance  $r = \int_t^{t_0} \frac{c dt'}{a(t')}$
- Friedmann Eq:  $1 = \Omega_m + \Omega_k + \Omega_\Lambda$
- $\Omega_m = \frac{\rho}{\rho_{\text{crit}}} = \frac{\rho_0 a^{-3}}{\rho_{\text{crit}0} (H^2/H_0^2)} = \Omega_{m0} H_0^2 a^{-1} \dot{a}^{-2}$

# distances in FLRW cosmology

- azimuthal equidistant  $r$ : proper distance at  $t_0 \equiv$  comoving radial

$$\text{distance } r = \int_t^{t_0} \frac{c dt'}{a(t')}$$

- Friedmann Eq:  $1 = \Omega_m + \Omega_k + \Omega_\Lambda$

- $\Omega_m = \frac{\rho}{\rho_{\text{crit}}} = \frac{\rho_0 a^{-3}}{\rho_{\text{crit}0} (H^2/H_0^2)} = \Omega_{m0} H_0^2 a^{-1} \dot{a}^{-2}$

- $\Omega_k = \Omega_{k0} \dot{a}^{-2} H_0^2$

# distances in FLRW cosmology

- azimuthal equidistant  $r$ : proper distance at  $t_0 \equiv$  comoving radial

$$\text{distance } r = \int_t^{t_0} \frac{c dt'}{a(t')}$$

- Friedmann Eq:  $1 = \Omega_m + \Omega_k + \Omega_\Lambda$

- $\underline{\Omega_m} = \frac{\rho}{\rho_{\text{crit}}} = \frac{\rho_0 a^{-3}}{\rho_{\text{crit}0} (H^2/H_0^2)} = \Omega_{m0} H_0^2 a^{-1} \dot{a}^{-2}$

- $\Omega_k = \Omega_{k0} \dot{a}^{-2} H_0^2$

- $\underline{\Omega_\Lambda} = \Omega_{\Lambda0} a^2 \dot{a}^{-2} H_0^2$

# distances in FLRW cosmology

- azimuthal equidistant  $r$ : proper distance at  $t_0 \equiv$  comoving radial

$$\text{distance } r = \int_t^{t_0} \frac{c dt'}{a(t')}$$

- Friedmann Eq:  $1 = \Omega_m + \Omega_k + \Omega_\Lambda$

- $\underline{\Omega_m} = \frac{\rho}{\rho_{\text{crit}}} = \frac{\rho_0 a^{-3}}{\rho_{\text{crit}0} (H^2/H_0^2)} = \Omega_{m0} H_0^2 a^{-1} \dot{a}^{-2}$

- $\Omega_k = \Omega_{k0} \dot{a}^{-2} H_0^2$

- $\underline{\Omega_\Lambda} = \Omega_{\Lambda0} a^2 \dot{a}^{-2} H_0^2$

- $1 = \Omega_{m0} H_0^2 a^{-1} \dot{a}^{-2} + \Omega_{k0} \dot{a}^{-2} H_0^2 + \Omega_{\Lambda0} a^2 \dot{a}^{-2} H_0^2$

# distances in FLRW cosmology

- azimuthal equidistant  $r$ : proper distance at  $t_0 \equiv$  comoving radial

$$\text{distance } r = \int_t^{t_0} \frac{c dt'}{a(t')}$$

- Friedmann Eq:  $1 = \Omega_m + \Omega_k + \Omega_\Lambda$

- $\underline{\Omega_m} = \frac{\rho}{\rho_{\text{crit}}} = \frac{\rho_0 a^{-3}}{\rho_{\text{crit}0} (H^2/H_0^2)} = \Omega_{m0} H_0^2 a^{-1} \dot{a}^{-2}$

- $\Omega_k = \Omega_{k0} \dot{a}^{-2} H_0^2$

- $\underline{\Omega_\Lambda} = \Omega_{\Lambda0} a^2 \dot{a}^{-2} H_0^2$

- $1 = \Omega_{m0} H_0^2 a^{-1} \dot{a}^{-2} + \Omega_{k0} \dot{a}^{-2} H_0^2 + \Omega_{\Lambda0} a^2 \dot{a}^{-2} H_0^2$

$$\Rightarrow \dot{a}^2 = H_0^2 (\Omega_{m0} a^{-1} + \Omega_{k0} + \Omega_{\Lambda0} a^2)$$

# distances in FLRW cosmology

- azimuthal equidistant  $r$ : proper distance at  $t_0 \equiv$  comoving radial

$$\text{distance } r = \int_t^{t_0} \frac{c dt'}{a(t')}$$

- Friedmann Eq:  $1 = \Omega_m + \Omega_k + \Omega_\Lambda$

- $\underline{\Omega_m} = \frac{\rho}{\rho_{\text{crit}}} = \frac{\rho_0 a^{-3}}{\rho_{\text{crit}0} (H^2/H_0^2)} = \Omega_{m0} H_0^2 a^{-1} \dot{a}^{-2}$

- $\Omega_k = \Omega_{k0} \dot{a}^{-2} H_0^2$

- $\underline{\Omega_\Lambda} = \Omega_{\Lambda0} a^2 \dot{a}^{-2} H_0^2$

- $1 = \Omega_{m0} H_0^2 a^{-1} \dot{a}^{-2} + \Omega_{k0} \dot{a}^{-2} H_0^2 + \Omega_{\Lambda0} a^2 \dot{a}^{-2} H_0^2$

$$\Rightarrow da = H_0 \sqrt{\Omega_{m0} a^{-1} + \Omega_{k0} + \Omega_{\Lambda0} a^2} dt \quad \text{if } \dot{a}, a, H_0 > 0$$

# distances in FLRW cosmology

- azimuthal equidistant  $r$ : proper distance at  $t_0 \equiv$  comoving radial

$$\text{distance } r = \int_t^{t_0} \frac{c dt'}{a(t')}$$

- Friedmann Eq:  $1 = \Omega_m + \Omega_k + \Omega_\Lambda$

- $\underline{\Omega_m} = \frac{\rho}{\rho_{\text{crit}}} = \frac{\rho_0 a^{-3}}{\rho_{\text{crit}0} (H^2/H_0^2)} = \Omega_{m0} H_0^2 a^{-1} \dot{a}^{-2}$

- $\Omega_k = \Omega_{k0} \dot{a}^{-2} H_0^2$

- $\underline{\Omega_\Lambda} = \Omega_{\Lambda0} a^2 \dot{a}^{-2} H_0^2$

- $1 = \Omega_{m0} H_0^2 a^{-1} \dot{a}^{-2} + \Omega_{k0} \dot{a}^{-2} H_0^2 + \Omega_{\Lambda0} a^2 \dot{a}^{-2} H_0^2$

$$\Rightarrow da = H_0 \sqrt{\Omega_{m0} a^{-1} + \Omega_{k0} + \Omega_{\Lambda0} a^2} dt \quad \text{if } \dot{a}, a, H_0 > 0$$

$$\Rightarrow dt = \frac{da}{H_0 \sqrt{\Omega_{m0} a^{-1} + \Omega_{k0} + \Omega_{\Lambda0} a^2}} \quad \text{if } \dot{a}, a, H_0 > 0$$



# distances in FLRW cosmology

- azimuthal equidistant  $r$ : proper distance at  $t_0 \equiv$  comoving radial

$$\text{distance } r = \int_t^{t_0} \frac{c dt'}{a(t')}$$

- Friedmann Eq:  $1 = \Omega_m + \Omega_k + \Omega_\Lambda$

- $\underline{\Omega_m} = \frac{\rho}{\rho_{\text{crit}}} = \frac{\rho_0 a^{-3}}{\rho_{\text{crit}0} (H^2/H_0^2)} = \Omega_{m0} H_0^2 a^{-1} \dot{a}^{-2}$

- $\Omega_k = \Omega_{k0} \dot{a}^{-2} H_0^2$

- $\underline{\Omega_\Lambda} = \Omega_{\Lambda0} a^2 \dot{a}^{-2} H_0^2$

- $1 = \Omega_{m0} H_0^2 a^{-1} \dot{a}^{-2} + \Omega_{k0} \dot{a}^{-2} H_0^2 + \Omega_{\Lambda0} a^2 \dot{a}^{-2} H_0^2$

$$\Rightarrow da = H_0 \sqrt{\Omega_{m0} a^{-1} + \Omega_{k0} + \Omega_{\Lambda0} a^2} dt \quad \text{if } \dot{a}, a, H_0 > 0$$

$$\Rightarrow \boxed{r = \int_{1/(1+z)}^1 \frac{c}{H_0} \frac{da}{a \sqrt{\Omega_{m0} a^{-1} + \Omega_{k0} + \Omega_{\Lambda0} a^2}}}$$

# Low $z$ limit Hubble law

- use  $r =$  comoving radial distance

# Low $z$ limit Hubble law

- use  $r =$  comoving radial distance
- consider  $0 < z \ll 1$  limit  $\Leftrightarrow 1 > a \gg 0.5$

# Low $z$ limit Hubble law

- use  $r =$  comoving radial distance
- consider  $0 < z \ll 1$  limit  $\Leftrightarrow 1 > a \gg 0.5$
- $a = \frac{1}{1+z}$

# Low $z$ limit Hubble law

- use  $r =$  comoving radial distance
- consider  $0 < z \ll 1$  limit  $\Leftrightarrow 1 > a \gg 0.5$
- $a = \frac{1}{1+z} \Rightarrow da = -(1+z)^{-2}dz$

# Low $z$ limit Hubble law

- use  $r =$  comoving radial distance
- consider  $0 < z \ll 1$  limit  $\Leftrightarrow 1 > a \gg 0.5$
- $a = \frac{1}{1+z} \Rightarrow da = -(1+z)^{-2}dz = -a^2dz$

# Low $z$ limit Hubble law

- use  $r =$  comoving radial distance
- consider  $0 < z \ll 1$  limit  $\Leftrightarrow 1 > a \gg 0.5$
- $a = \frac{1}{1+z} \Rightarrow da = -(1+z)^{-2}dz = -a^2dz$

$$r = \int_{1/(1+z)}^1 \frac{c}{H_0 a \sqrt{\Omega_{m0}a^{-1} + \Omega_{k0} + \Omega_{\Lambda0}a^2}} da$$

# Low $z$ limit Hubble law

- use  $r =$  comoving radial distance
- consider  $0 < z \ll 1$  limit  $\Leftrightarrow 1 > a \gg 0.5$
- $a = \frac{1}{1+z} \Rightarrow da = -(1+z)^{-2} dz = -a^2 dz$

$$\begin{aligned} r &= \int_{1/(1+z)}^1 \frac{c}{H_0 a \sqrt{\Omega_{m0} a^{-1} + \Omega_{k0} + \Omega_{\Lambda0} a^2}} da \\ &= - \int_z^0 \frac{c}{H_0 a \sqrt{\Omega_{m0} a^{-1} + \Omega_{k0} + \Omega_{\Lambda0} a^2}} a^2 dz \end{aligned}$$



# Low $z$ limit Hubble law

- use  $r =$  comoving radial distance
- consider  $0 < z \ll 1$  limit  $\Leftrightarrow 1 > a \gg 0.5$
- $a = \frac{1}{1+z} \Rightarrow da = -(1+z)^{-2} dz = -a^2 dz$

$$\begin{aligned} r &= \int_{1/(1+z)}^1 \frac{c}{H_0 a \sqrt{\Omega_{m0} a^{-1} + \Omega_{k0} + \Omega_{\Lambda0} a^2}} da \\ &= - \int_z^0 \frac{c}{H_0 \sqrt{\Omega_{m0} a^{-1} + \Omega_{k0} + \Omega_{\Lambda0} a^2}} a dz \end{aligned}$$

# Low $z$ limit Hubble law

- use  $r =$  comoving radial distance
- consider  $0 < z \ll 1$  limit  $\Leftrightarrow 1 > a \gg 0.5$
- $a = \frac{1}{1+z} \Rightarrow da = -(1+z)^{-2} dz = -a^2 dz$

$$r = \int_{1/(1+z)}^1 \frac{c}{H_0 a \sqrt{\Omega_{m0} a^{-1} + \Omega_{k0} + \Omega_{\Lambda 0} a^2}} da$$
$$\approx - \int_z^0 \frac{c}{H_0 \sqrt{\Omega_{m0} + \Omega_{k0} + \Omega_{\Lambda 0}}} dz$$

# Low $z$ limit Hubble law

- use  $r =$  comoving radial distance
- consider  $0 < z \ll 1$  limit  $\Leftrightarrow 1 > a \gg 0.5$
- $a = \frac{1}{1+z} \Rightarrow da = -(1+z)^{-2} dz = -a^2 dz$

$$r = \int_{1/(1+z)}^1 \frac{c}{H_0 a \sqrt{\Omega_{m0} a^{-1} + \Omega_{k0} + \Omega_{\Lambda0} a^2}} da$$
$$\approx - \int_z^0 \frac{c}{H_0} \frac{1}{\sqrt{1}} dz$$

# Low $z$ limit Hubble law

- use  $r =$  comoving radial distance
- consider  $0 < z \ll 1$  limit  $\Leftrightarrow 1 > a \gg 0.5$
- $a = \frac{1}{1+z} \Rightarrow da = -(1+z)^{-2} dz = -a^2 dz$

$$\begin{aligned} r &= \int_{1/(1+z)}^1 \frac{c}{H_0 a \sqrt{\Omega_{m0} a^{-1} + \Omega_{k0} + \Omega_{\Lambda0} a^2}} da \\ &\approx - \int_z^0 \frac{c}{H_0} \frac{1}{\sqrt{1}} dz \\ &= - \int_z^0 \frac{c}{H_0} dz \end{aligned}$$

# Low $z$ limit Hubble law

- use  $r =$  comoving radial distance
- consider  $0 < z \ll 1$  limit  $\Leftrightarrow 1 > a \gg 0.5$
- $a = \frac{1}{1+z} \Rightarrow da = -(1+z)^{-2} dz = -a^2 dz$

$$\begin{aligned} r &= \int_{1/(1+z)}^1 \frac{c}{H_0 a \sqrt{\Omega_{m0} a^{-1} + \Omega_{k0} + \Omega_{\Lambda 0} a^2}} da \\ &\approx - \int_z^0 \frac{c}{H_0} \frac{1}{\sqrt{1}} dz \\ &= - \int_z^0 \frac{c}{H_0} dz \end{aligned}$$

$$\Rightarrow \frac{dr}{dz} \approx \frac{c}{H_0}$$

# Low $z$ limit Hubble law

- use  $r =$  comoving radial distance
- consider  $0 < z \ll 1$  limit  $\Leftrightarrow 1 > a \gg 0.5$
- $a = \frac{1}{1+z} \Rightarrow da = -(1+z)^{-2} dz = -a^2 dz$

$$\begin{aligned} r &= \int_{1/(1+z)}^1 \frac{c}{H_0 a \sqrt{\Omega_{m0} a^{-1} + \Omega_{k0} + \Omega_{\Lambda 0} a^2}} da \\ &\approx - \int_z^0 \frac{c}{H_0} \frac{1}{\sqrt{1}} dz \\ &= - \int_z^0 \frac{c}{H_0} dz \end{aligned}$$

$$\Rightarrow \frac{dr}{dz} \approx \frac{c}{H_0}$$

$$\Rightarrow \text{when } 0 < z \ll 1, \beta := v/c, r \approx \frac{z c}{H_0}$$

# Low $z$ limit Hubble law

- use  $r =$  comoving radial distance
- consider  $0 < z \ll 1$  limit  $\Leftrightarrow 1 > a \gg 0.5$
- $a = \frac{1}{1+z} \Rightarrow da = -(1+z)^{-2} dz = -a^2 dz$

$$\begin{aligned} r &= \int_{1/(1+z)}^1 \frac{c}{H_0 a \sqrt{\Omega_{m0} a^{-1} + \Omega_{k0} + \Omega_{\Lambda0} a^2}} da \\ &\approx - \int_z^0 \frac{c}{H_0} \frac{1 dz}{\sqrt{1}} \\ &= - \int_z^0 \frac{c}{H_0} dz \end{aligned}$$

$$\Rightarrow \frac{dr}{dz} \approx \frac{c}{H_0}$$

$$\Rightarrow \text{when } 0 < z \ll 1, \beta := v/c, r \approx \frac{z c}{H_0} \approx \frac{\beta c}{H_0}$$

# Low $z$ limit Hubble law

- use  $r =$  comoving radial distance
- consider  $0 < z \ll 1$  limit  $\Leftrightarrow 1 > a \gg 0.5$
- $a = \frac{1}{1+z} \Rightarrow da = -(1+z)^{-2} dz = -a^2 dz$

$$\begin{aligned} r &= \int_{1/(1+z)}^1 \frac{c}{H_0 a \sqrt{\Omega_{m0} a^{-1} + \Omega_{k0} + \Omega_{\Lambda 0} a^2}} da \\ &\approx - \int_z^0 \frac{c}{H_0} \frac{1}{\sqrt{1}} dz \\ &= - \int_z^0 \frac{c}{H_0} dz \end{aligned}$$

$$\Rightarrow \frac{dr}{dz} \approx \frac{c}{H_0}$$

$$\Rightarrow \text{when } 0 < z \ll 1, \beta := v/c, \quad \boxed{r \approx \frac{z c}{H_0} \approx \frac{\beta c}{H_0} = \frac{v}{H_0}}$$



# EdS: radial comoving distance

$$\blacksquare r = \int_{1/(1+z)}^1 \frac{c}{H_0} \frac{da}{a \sqrt{\Omega_{m0} a^{-1} + \Omega_{k0} + \Omega_{\Lambda 0} a^2}}$$

# EdS: radial comoving distance

$$\blacksquare r = \int_{1/(1+z)}^1 \frac{c}{H_0} \frac{da}{a \sqrt{\Omega_{m0} a^{-1} + \Omega_{k0} + \Omega_{\Lambda 0} a^2}}$$

$$\text{EdS: } \Omega_{m0} = 1, \Omega_{k0} = 0, \Omega_{\Lambda 0} = 0$$

# EdS: radial comoving distance

$$\blacksquare r = \int_{1/(1+z)}^1 \frac{c}{H_0} \frac{da}{a \sqrt{\Omega_{m0} a^{-1} + \Omega_{k0} + \Omega_{\Lambda 0} a^2}}$$

$$\text{EdS: } \Omega_{m0} = 1, \Omega_{k0} = 0, \Omega_{\Lambda 0} = 0$$

$$\Rightarrow r_{\text{EdS}} = \int_{1/(1+z)}^1 \frac{c}{H_0} \frac{da}{a \sqrt{a^{-1}}}$$

# EdS: radial comoving distance

$$\blacksquare r = \int_{1/(1+z)}^1 \frac{c}{H_0} \frac{da}{a \sqrt{\Omega_{m0} a^{-1} + \Omega_{k0} + \Omega_{\Lambda0} a^2}}$$

$$\text{EdS: } \Omega_{m0} = 1, \Omega_{k0} = 0, \Omega_{\Lambda0} = 0$$

$$\begin{aligned} \Rightarrow r_{\text{EdS}} &= \int_{1/(1+z)}^1 \frac{c}{H_0} \frac{da}{a \sqrt{a^{-1}}} \\ &= \frac{c}{H_0} \int_{1/(1+z)}^1 a^{-1/2} da \end{aligned}$$

# EdS: radial comoving distance

$$\blacksquare r = \int_{1/(1+z)}^1 \frac{c}{H_0} \frac{da}{a \sqrt{\Omega_{m0} a^{-1} + \Omega_{k0} + \Omega_{\Lambda0} a^2}}$$

$$\text{EdS: } \Omega_{m0} = 1, \Omega_{k0} = 0, \Omega_{\Lambda0} = 0$$

$$\begin{aligned} \Rightarrow r_{\text{EdS}} &= \int_{1/(1+z)}^1 \frac{c}{H_0} \frac{da}{a \sqrt{a^{-1}}} \\ &= \frac{c}{H_0} \int_{1/(1+z)}^1 a^{-1/2} da \\ &= \frac{c}{H_0} \left[ 2a^{1/2} \right]_{1/(1+z)}^1 \end{aligned}$$

# EdS: radial comoving distance

$$\blacksquare r = \int_{1/(1+z)}^1 \frac{c}{H_0} \frac{da}{a \sqrt{\Omega_{m0} a^{-1} + \Omega_{k0} + \Omega_{\Lambda0} a^2}}$$

$$\text{EdS: } \Omega_{m0} = 1, \Omega_{k0} = 0, \Omega_{\Lambda0} = 0$$

$$\begin{aligned} \Rightarrow r_{\text{EdS}} &= \int_{1/(1+z)}^1 \frac{c}{H_0} \frac{da}{a \sqrt{a^{-1}}} \\ &= \frac{c}{H_0} \int_{1/(1+z)}^1 a^{-1/2} da \\ &= \frac{c}{H_0} \left[ 2a^{1/2} \right]_{1/(1+z)}^1 \\ &= \boxed{\frac{2c}{H_0} \left( 1 - \frac{1}{\sqrt{1+z}} \right)} \end{aligned}$$

# EdS: radial comoving distance

$$\blacksquare r = \int_{1/(1+z)}^1 \frac{c}{H_0} \frac{da}{a \sqrt{\Omega_{m0} a^{-1} + \Omega_{k0} + \Omega_{\Lambda0} a^2}}$$

$$\text{EdS: } \Omega_{m0} = 1, \Omega_{k0} = 0, \Omega_{\Lambda0} = 0$$

$\Rightarrow$

$$\begin{aligned} r_{\text{EdS}} &= \int_{1/(1+z)}^1 \frac{c}{H_0} \frac{da}{a \sqrt{a^{-1}}} \\ &= \frac{c}{H_0} \int_{1/(1+z)}^1 a^{-1/2} da \\ &= \frac{c}{H_0} \left[ 2a^{1/2} \right]_{1/(1+z)}^1 \\ &= \boxed{\frac{2c}{H_0} \left( 1 - \frac{1}{\sqrt{1+z}} \right)} \\ &\approx 6h^{-1} \text{Gpc} \left( 1 - \frac{1}{\sqrt{1+z}} \right) \end{aligned}$$

# distances in FLRW cosmology

$$\blacksquare r = \int_{1/(1+z)}^1 \frac{c}{H_0} \frac{da}{a \sqrt{\Omega_{m0} a^{-1} + \Omega_{k0} + \Omega_{\Lambda 0} a^2}}$$



# distances in FLRW cosmology

- $$r = \int_{1/(1+z)}^1 \frac{c}{H_0} \frac{da}{a \sqrt{\Omega_{m0} a^{-1} + \Omega_{k0} + \Omega_{\Lambda 0} a^2}}$$

- GPL numerical package: `cosmdist`

<http://cosmo.torun.pl/GPLdownload/cosmdist/>

# distances in FLRW cosmology

- $r = \int_{1/(1+z)}^1 \frac{c}{H_0} \frac{da}{a \sqrt{\Omega_{m0} a^{-1} + \Omega_{k0} + \Omega_{\Lambda 0} a^2}}$
- GPL numerical package: cosmdist  
<http://cosmo.torun.pl/GPLdownload/cosmdist/>
- standard install to /usr/local:  
`./configure && make && make install`

# distances in FLRW cosmology

- $r = \int_{1/(1+z)}^1 \frac{c}{H_0} \frac{da}{a \sqrt{\Omega_{m0} a^{-1} + \Omega_{k0} + \Omega_{\Lambda0} a^2}}$
- GPL numerical package: `cosmdist`  
<http://cosmo.torun.pl/GPLdownload/cosmdist/>
- standard install to `/usr/local`:  
`./configure && make && make install`
- command line pipe:  
`cat myredshifts.lis | cosmdist`

# distances in FLRW cosmology

- $r = \int_{1/(1+z)}^1 \frac{c}{H_0} \frac{da}{a \sqrt{\Omega_{m0} a^{-1} + \Omega_{k0} + \Omega_{\Lambda 0} a^2}}$
- GPL numerical package: `cosmdist`  
<http://cosmo.torun.pl/GPLdownload/cosmdist/>
- standard install to `/usr/local`:  
`./configure && make && make install`
- command line pipe:  
`cat myredshifts.lis | cosmdist`
- command line options: `cosmdist --help`

# distances in FLRW cosmology

- $r = \int_{1/(1+z)}^1 \frac{c}{H_0} \frac{da}{a \sqrt{\Omega_{m0} a^{-1} + \Omega_{k0} + \Omega_{\Lambda 0} a^2}}$
- GPL numerical package: `cosmdist`  
<http://cosmo.torun.pl/GPLdownload/cosmdist/>
- standard install to `/usr/local`:  
`./configure && make && make install`
- command line pipe:  
`cat myredshifts.lis | cosmdist`
- command line options: `cosmdist --help`
- static fortran or C library: link to `libcosmdist.a`

# distances in FLRW cosmology

- $r = \int_{1/(1+z)}^1 \frac{c}{H_0} \frac{da}{a \sqrt{\Omega_{m0} a^{-1} + \Omega_{k0} + \Omega_{\Lambda 0} a^2}}$
- GPL numerical package: `cosmdist`  
<http://cosmo.torun.pl/GPLdownload/cosmdist/>
- standard install to `/usr/local`:  
`./configure && make && make install`
- command line pipe:  
`cat myredshifts.lis | cosmdist`
- command line options: `cosmdist --help`
- static fortran or C library: link to `libcosmdist.a`
- high-level frontends (e.g. python) should be easy to write

# tangential arc-lengths: $r_{\perp}$ vs $d_A$

- $r_{\perp}$  =: “proper motion” distance (Weinberg 1972)

# tangential arc-lengths: $r_{\perp}$ vs $d_A$

- $r_{\perp}$  =: “proper motion” distance (Weinberg 1972)
- $\equiv$  **comoving** tangential arc-length for 1 radian



# tangential arc-lengths: $r_{\perp}$ vs $d_A$

■  $r_{\perp}$  =: “proper motion” distance (Weinberg 1972)

■  $\equiv$  **comoving** tangential arc-length for 1 radian

■ azimuthal equidistant coords:  $r_{\perp} = \begin{cases} R_C \sinh \frac{r}{R_C} & k < 0 \\ r & k = 0 \\ R_C \sin \frac{r}{R_C} & k > 0 \end{cases}$

# tangential arc-lengths: $r_{\perp}$ vs $d_A$

■  $r_{\perp}$  =: “proper motion” distance (Weinberg 1972)

■  $\equiv$  **comoving** tangential arc-length for 1 radian

■ azimuthal equidistant coords:  $r_{\perp} = \begin{cases} R_C \sinh \frac{r}{R_C} & k < 0 \\ r & k = 0 \\ R_C \sin \frac{r}{R_C} & k > 0 \end{cases}$  where

$r$  = comoving radial distance

■ angular diameter distance  $d_A$

# tangential arc-lengths: $r_{\perp}$ vs $d_A$

■  $r_{\perp}$  =: “proper motion” distance (Weinberg 1972)

■  $\equiv$  **comoving** tangential arc-length for 1 radian

■ azimuthal equidistant coords:  $r_{\perp} = \begin{cases} R_C \sinh \frac{r}{R_C} & k < 0 \\ r & k = 0 \\ R_C \sin \frac{r}{R_C} & k > 0 \end{cases}$  where

$r$  = comoving radial distance

■ angular diameter distance  $d_A$

■  $\equiv$  **physical/**

# tangential arc-lengths: $r_{\perp}$ vs $d_A$

■  $r_{\perp}$  =: “proper motion” distance (Weinberg 1972)

■  $\equiv$  **comoving** tangential arc-length for 1 radian

■ azimuthal equidistant coords:  $r_{\perp} = \begin{cases} R_C \sinh \frac{r}{R_C} & k < 0 \\ r & k = 0 \\ R_C \sin \frac{r}{R_C} & k > 0 \end{cases}$  where

$r$  = comoving radial distance

■ angular diameter distance  $d_A$

■  $\equiv$  **physical/ metric /**

# tangential arc-lengths: $r_{\perp}$ vs $d_A$

■  $r_{\perp}$  =: “proper motion” distance (Weinberg 1972)

■  $\equiv$  **comoving** tangential arc-length for 1 radian

■ azimuthal equidistant coords:  $r_{\perp} = \begin{cases} R_C \sinh \frac{r}{R_C} & k < 0 \\ r & k = 0 \\ R_C \sin \frac{r}{R_C} & k > 0 \end{cases}$  where

$r$  = comoving radial distance

■ angular diameter distance  $d_A$

■  $\equiv$  **physical / metric / local**

# tangential arc-lengths: $r_{\perp}$ vs $d_A$

■  $r_{\perp}$  =: “proper motion” distance (Weinberg 1972)

■  $\equiv$  **comoving** tangential arc-length for 1 radian

■ azimuthal equidistant coords:  $r_{\perp} = \begin{cases} R_C \sinh \frac{r}{R_C} & k < 0 \\ r & k = 0 \\ R_C \sin \frac{r}{R_C} & k > 0 \end{cases}$  where

$r$  = comoving radial distance

■ angular diameter distance  $d_A$

■  $\equiv$  **physical / metric / local** arc-length for 1 radian

# tangential arc-lengths: $r_{\perp}$ vs $d_A$

- $r_{\perp}$  =: “proper motion” distance (Weinberg 1972)

- $\equiv$  **comoving** tangential arc-length for 1 radian

- azimuthal equidistant coords:  $r_{\perp} = \begin{cases} R_C \sinh \frac{r}{R_C} & k < 0 \\ r & k = 0 \\ R_C \sin \frac{r}{R_C} & k > 0 \end{cases}$  where

$r$  = comoving radial distance

- angular diameter distance  $d_A$

- $\equiv$  **physical/ metric / local** arc-length for 1 radian

- $d_A = r_{\perp} a$  (scaled by the scale factor)

# tangential arc-lengths: $r_{\perp}$ vs $d_A$

- $r_{\perp}$  =: “proper motion” distance (Weinberg 1972)

- $\equiv$  **comoving** tangential arc-length for 1 radian

- azimuthal equidistant coords:  $r_{\perp} = \begin{cases} R_C \sinh \frac{r}{R_C} & k < 0 \\ r & k = 0 \\ R_C \sin \frac{r}{R_C} & k > 0 \end{cases}$  where

$r$  = comoving radial distance

- angular diameter distance  $d_A$

- $\equiv$  **physical / metric / local** arc-length for 1 radian

- $$d_A = \frac{r_{\perp}}{1+z}$$



# luminosity distance

want to define  $d_L$  such that flux is  $F = \frac{L}{4\pi d_L^2}$

# luminosity distance

want to define  $d_L$  such that flux is  $F = \frac{L}{4\pi d_L^2}$

- surface area of 2-sphere =  $4\pi r_{\perp}^2$  (comoving)

# luminosity distance

want to define  $d_L$  such that flux is  $F = \frac{L}{4\pi d_L^2}$

■  $E = h\nu$  of photon:

# luminosity distance

want to define  $d_L$  such that flux is  $F = \frac{L}{4\pi d_L^2}$

- $E = h\nu$  of photon: emitter frame to observer frame  
 $F \propto 1/\lambda \propto a_{\text{em}}$

# luminosity distance

want to define  $d_L$  such that flux is  $F = \frac{L}{4\pi d_L^2}$

- $E = h\nu$  of photon: emitter frame to observer frame  
 $F \propto 1/\lambda \propto a_{\text{em}}$
- erg per s: time dilation:

# luminosity distance

want to define  $d_L$  such that flux is  $F = \frac{L}{4\pi d_L^2}$

■  $E = h\nu$  of photon: emitter frame to observer frame

$$F \propto 1/\lambda \propto a_{\text{em}}$$

■ **erg per s**: time dilation: emitter frame to observer frame  $F \propto a_{\text{em}}$

# luminosity distance

want to define  $d_L$  such that flux is  $F = \frac{L}{4\pi d_L^2}$

■  $E = h\nu$  of photon: emitter frame to observer frame

$$F \propto 1/\lambda \propto a_{\text{em}}$$

■ **erg per s**: time dilation: emitter frame to observer frame  $F \propto a_{\text{em}}$

■  $\Rightarrow F = \frac{L}{4\pi r_{\perp}^2} a^2$

# luminosity distance

want to define  $d_L$  such that flux is  $F = \frac{L}{4\pi d_L^2}$

■  $E = h\nu$  of photon: emitter frame to observer frame

$$F \propto 1/\lambda \propto a_{\text{em}}$$

■ **erg per s**: time dilation: emitter frame to observer frame  $F \propto a_{\text{em}}$

$$\Rightarrow F = \frac{L}{4\pi r_{\perp}^2} a^2 = \frac{L}{4\pi r_{\perp}^2 (1+z)^2}$$



# luminosity distance

want to define  $d_L$  such that flux is  $F = \frac{L}{4\pi d_L^2}$

■  $E = h\nu$  of photon: emitter frame to observer frame

$$F \propto 1/\lambda \propto a_{\text{em}}$$

■ erg per s: time dilation: emitter frame to observer frame  $F \propto a_{\text{em}}$

$$\Rightarrow F = \frac{L}{4\pi r_{\perp}^2} a^2 = \frac{L}{4\pi r_{\perp}^2 (1+z)^2}$$

■  $\Rightarrow$  luminosity distance  $d_L := r_{\perp}(1+z) = d_A(1+z)^2$

# luminosity distance

want to define  $d_L$  such that flux is  $F = \frac{L}{4\pi d_L^2}$

- $E = h\nu$  of photon: emitter frame to observer frame

$$F \propto 1/\lambda \propto a_{\text{em}}$$

- erg per s: time dilation: emitter frame to observer frame  $F \propto a_{\text{em}}$

- $\Rightarrow F = \frac{L}{4\pi r_{\perp}^2} a^2 = \frac{L}{4\pi r_{\perp}^2 (1+z)^2}$

- $\Rightarrow$  luminosity distance  $d_L := r_{\perp}(1+z) = d_A(1+z)^2$

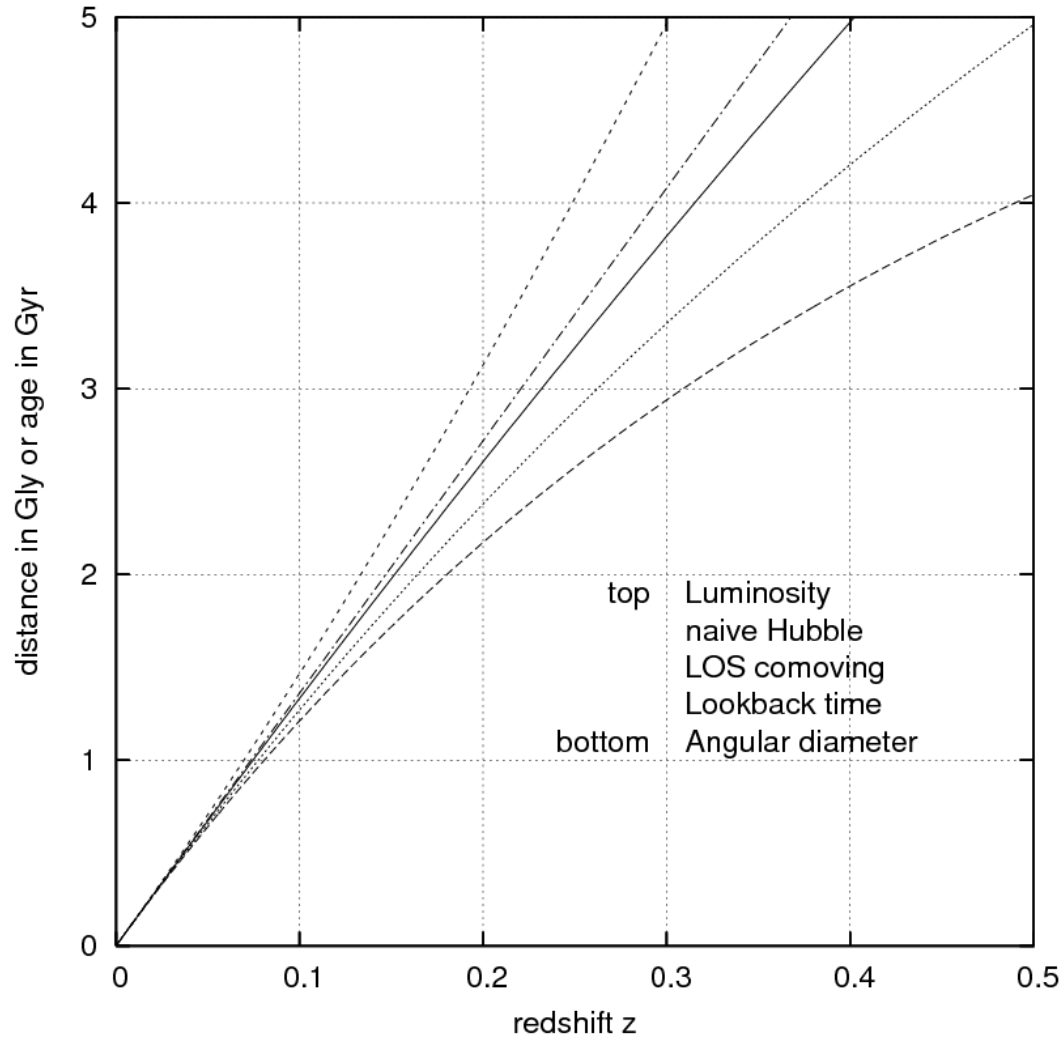
- w:Distance measures (cosmology)



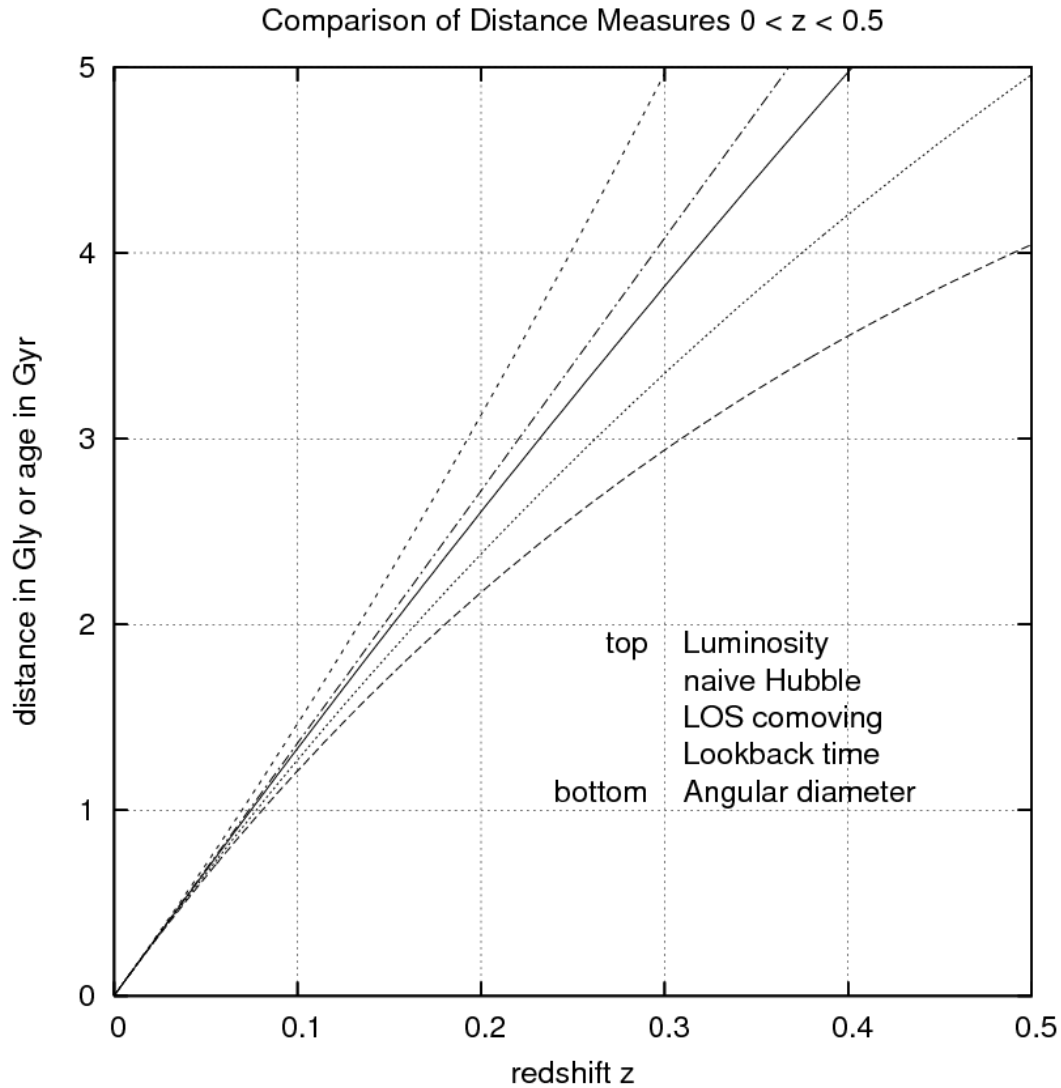
# FLRW distances e.g. $\Lambda$ CDM



Comparison of Distance Measures  $0 < z < 0.5$

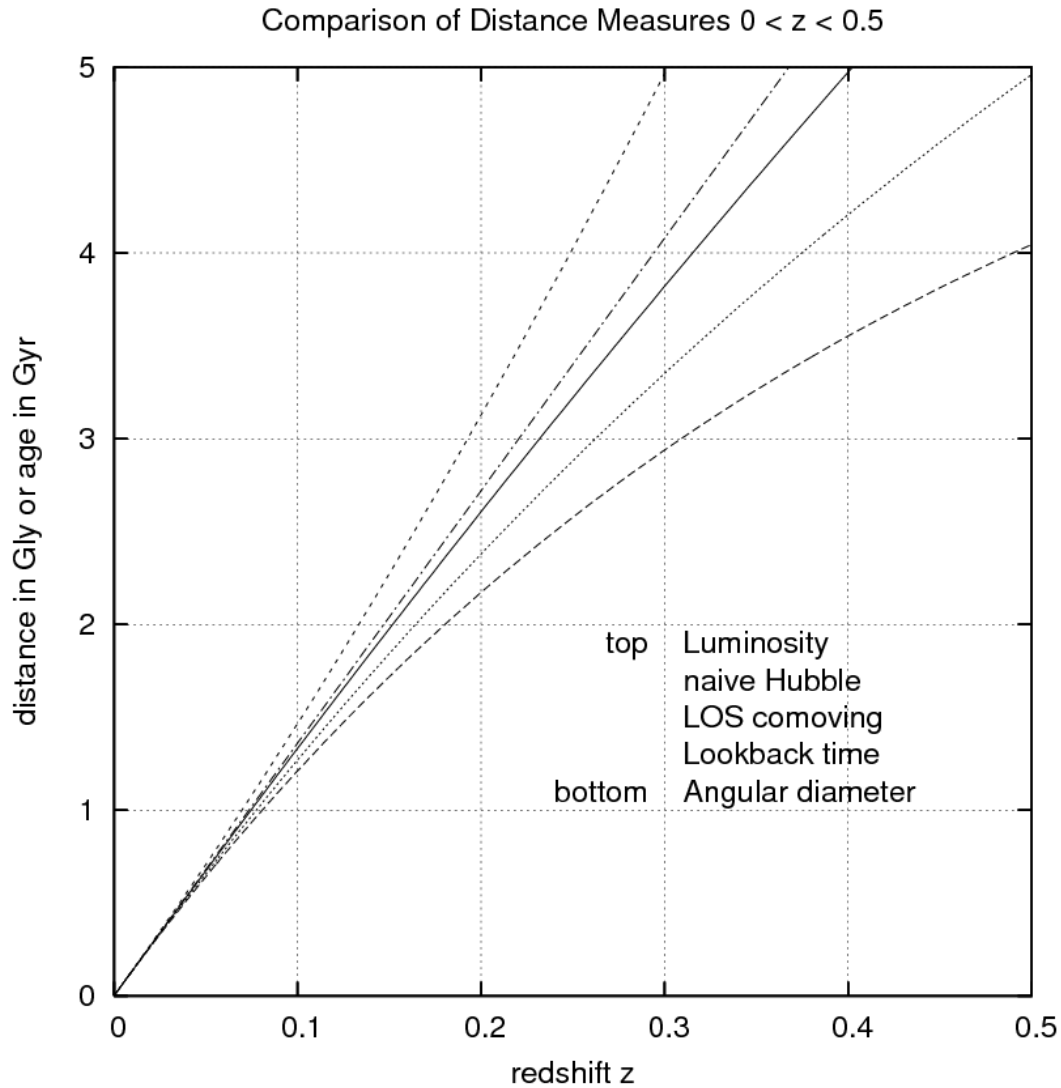


# FLRW distances e.g. $\Lambda$ CDM



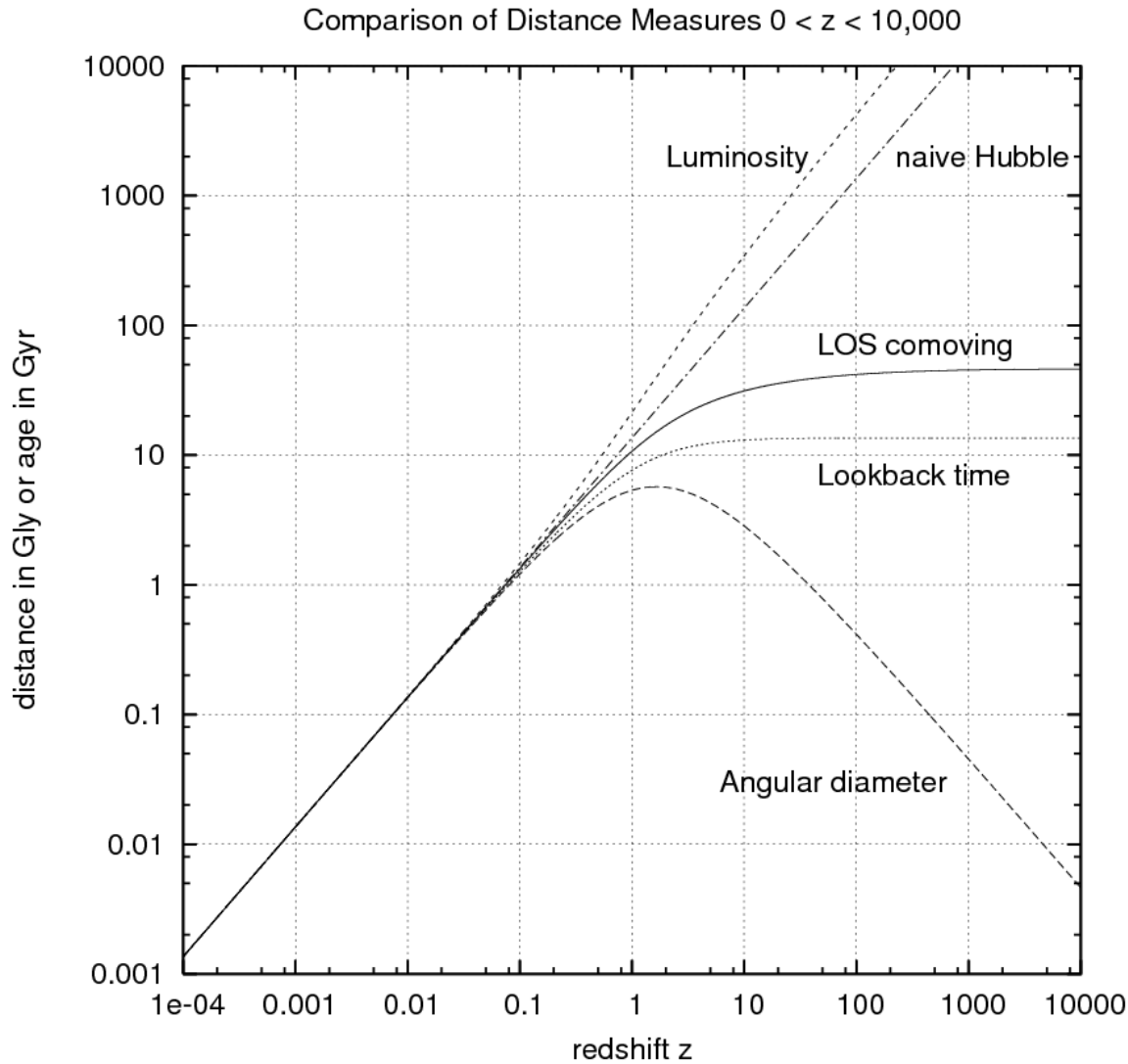
Defn:  $h := H_0/100 \text{ km/s/Mpc}$  (without a "0" subscript on  $h$ )

# FLRW distances e.g. $\Lambda$ CDM



$$\Omega_{m0} = 0.3, \Omega_{r0} = 10^{-4}, \Omega_{\Lambda 0} = 1.0 - (\Omega_{m0} + \Omega_{r0}), h = 0.7, \Omega_{k0} = 0$$

# FLRW distances e.g. $\Lambda$ CDM



$$\Omega_{m0} = 0.3, \Omega_{r0} = 10^{-4}, \Omega_{\Lambda0} = 1.0 - (\Omega_{m0} + \Omega_{r0}), h = 0.7, \Omega_{k0} = 0$$

# “Velocity” $v > c$ in FLRW?

- low  $z$  Hubble law  $r \approx \frac{\beta c}{H_0}$ :  $r > 3h^{-1}$  Gpc  $\Rightarrow \beta > 1$ ?

# “Velocity” $v > c$ in FLRW?

- low  $z$  Hubble law  $r \approx \frac{\beta c}{H_0}$ :  $r > 3h^{-1}$  Gpc  $\Rightarrow \beta > 1$ ?
- possible definitions of “velocity”:
  1. galaxy as fundamental observer,



# “Velocity” $v > c$ in FLRW?

■ low  $z$  Hubble law  $r \approx \frac{\beta c}{H_0}$ :  $r > 3h^{-1}$  Gpc  $\Rightarrow \beta > 1$ ?

■ possible definitions of “velocity”:

1. galaxy as fundamental observer, local velocity  $v_{\text{gal}} = 0$  by defn of comoving coords,

# “Velocity” $v > c$ in FLRW?

■ low  $z$  Hubble law  $r \approx \frac{\beta c}{H_0}$ :  $r > 3h^{-1}$  Gpc  $\Rightarrow \beta > 1$ ?

■ possible definitions of “velocity”:

1. galaxy as fundamental observer, local velocity  $v_{\text{gal}} = 0$  by defn of comoving coords, or
2. Hubble law assumes  $\parallel$  transport of 4-velocity  $\vec{u}_{\text{gal}}$ ; then compare  $\vec{u}_{\text{obs}}, \vec{u}_{\text{gal}}$  *locally* in Minkowski spacetime

# “Velocity” $v > c$ in FLRW?

■ low  $z$  Hubble law  $r \approx \frac{\beta c}{H_0}$ :  $r > 3h^{-1}$  Gpc  $\Rightarrow \beta > 1$ ?

■ possible definitions of “velocity”:

1. galaxy as fundamental observer, local velocity  $v_{\text{gal}} = 0$  by defn of comoving coords, or

2. Hubble law assumes  $\parallel$  transport of 4-velocity  $\vec{u}_{\text{gal}}$ ; then compare  $\vec{u}_{\text{obs}}, \vec{u}_{\text{gal}}$  *locally* in Minkowski spacetime i.e.

$1 + z = \sqrt{\frac{1+\beta}{1-\beta}}$  is *assumed* in derivation

# “Velocity” $v > c$ in FLRW?

■ low  $z$  Hubble law  $r \approx \frac{\beta c}{H_0}$ :  $r > 3h^{-1}$  Gpc  $\Rightarrow \beta > 1$ ?

■ possible definitions of “velocity”:

1. galaxy as fundamental observer, local velocity  $v_{\text{gal}} = 0$  by defn of comoving coords, or

2. Hubble law assumes  $\parallel$  transport of 4-velocity  $\vec{u}_{\text{gal}}$ ; then compare  $\vec{u}_{\text{obs}}, \vec{u}_{\text{gal}}$  *locally* in Minkowski spacetime i.e.

$1 + z = \sqrt{\frac{1+\beta}{1-\beta}}$  is *assumed* in derivation  $v_{\text{gal}} < c$  by assumption ( $\parallel$  transport + Mink.)

# “Velocity” $v > c$ in FLRW?

■ low  $z$  Hubble law  $r \approx \frac{\beta c}{H_0}$ :  $r > 3h^{-1}$  Gpc  $\Rightarrow \beta > 1$ ?

■ possible definitions of “velocity”:

1. galaxy as fundamental observer, local velocity  $v_{\text{gal}} = 0$  by defn of comoving coords, or

2.  $v_{\text{gal}} < c$  by assumption ( $\parallel$  transport + Mink.) or

3. define differential “global” comoving velocity

$$v_{\text{gal}}^{\text{glob}} := \frac{d\left(\int_0^r ds|_{t,\theta,\phi}\right)}{dt}$$

# “Velocity” $v > c$ in FLRW?

■ low  $z$  Hubble law  $r \approx \frac{\beta c}{H_0}$ :  $r > 3h^{-1}$  Gpc  $\Rightarrow \beta > 1$ ?

■ possible definitions of “velocity”:

1. galaxy as fundamental observer, local velocity  $v_{\text{gal}} = 0$  by defn of comoving coords, or

2.  $v_{\text{gal}} < c$  by assumption ( $\parallel$  transport + Mink.) or

3. define differential “global” comoving velocity

$$v_{\text{gal}}^{\text{glob}} := \frac{d\left(\int_0^r ds|_{t,\theta,\phi}\right)}{dt} = \frac{d(ar)}{dt}$$

# “Velocity” $v > c$ in FLRW?

■ low  $z$  Hubble law  $r \approx \frac{\beta c}{H_0}$ :  $r > 3h^{-1}$  Gpc  $\Rightarrow \beta > 1$ ?

■ possible definitions of “velocity”:

1. galaxy as fundamental observer, local velocity  $v_{\text{gal}} = 0$  by defn of comoving coords, or

2.  $v_{\text{gal}} < c$  by assumption ( $\parallel$  transport + Mink.) or

3. define differential “global” comoving velocity

$$v_{\text{gal}}^{\text{glob}} := \frac{d\left(\int_0^r ds|_{t,\theta,\phi}\right)}{dt} = \frac{d(ar)}{dt} = r \dot{a}$$

# “Velocity” $v > c$ in FLRW?

■ low  $z$  Hubble law  $r \approx \frac{\beta c}{H_0}$ :  $r > 3h^{-1}$  Gpc  $\Rightarrow \beta > 1$ ?

■ possible definitions of “velocity”:

1. galaxy as fundamental observer, local velocity  $v_{\text{gal}} = 0$  by defn of comoving coords, or

2.  $v_{\text{gal}} < c$  by assumption ( $\parallel$  transport + Mink.) or

3. define differential “global” comoving velocity

$$v_{\text{gal}}^{\text{glob}} := \frac{d\left(\int_0^r ds|_{t,\theta,\phi}\right)}{dt} = \frac{d(ar)}{dt} = r \dot{a} = r \frac{\dot{a}}{a} a$$



# “Velocity” $v > c$ in FLRW?

■ low  $z$  Hubble law  $r \approx \frac{\beta c}{H_0}$ :  $r > 3h^{-1}$  Gpc  $\Rightarrow \beta > 1$ ?

■ possible definitions of “velocity”:

1. galaxy as fundamental observer, local velocity  $v_{\text{gal}} = 0$  by defn of comoving coords, or

2.  $v_{\text{gal}} < c$  by assumption ( $\parallel$  transport + Mink.) or

3. define differential “global” comoving velocity

$$v_{\text{gal}}^{\text{glob}} := \frac{d\left(\int_0^r ds|_{t,\theta,\phi}\right)}{dt} = \frac{d(ar)}{dt} = r \dot{a} = rH_0$$

# “Velocity” $v > c$ in FLRW?

■ low  $z$  Hubble law  $r \approx \frac{\beta c}{H_0}$ :  $r > 3h^{-1}$  Gpc  $\Rightarrow \beta > 1$ ?

■ possible definitions of “velocity”:

1. galaxy as fundamental observer, local velocity  $v_{\text{gal}} = 0$  by defn of comoving coords, or

2.  $v_{\text{gal}} < c$  by assumption ( $\parallel$  transport + Mink.) or

3.  $v_{\text{gal}}^{\text{glob}} := \frac{d(\int_0^r ds|_{t,\theta,\phi})}{dt} = \frac{d(ar)}{dt} = r \dot{a} = rH_0$  or

4. define mean “global” comoving velocity

$$\langle v \rangle_{\text{gal}}^{\text{glob}} := \frac{\int_0^r ds|_{t,\theta,\phi}}{t_0}$$

# “Velocity” $v > c$ in FLRW?

■ low  $z$  Hubble law  $r \approx \frac{\beta c}{H_0}$ :  $r > 3h^{-1}$  Gpc  $\Rightarrow \beta > 1$ ?

■ possible definitions of “velocity”:

1. galaxy as fundamental observer, local velocity  $v_{\text{gal}} = 0$  by defn of comoving coords, or

2.  $v_{\text{gal}} < c$  by assumption ( $\parallel$  transport + Mink.) or

3.  $v_{\text{gal}}^{\text{glob}} := \frac{d(\int_0^r ds|_{t,\theta,\phi})}{dt} = \frac{d(ar)}{dt} = r \dot{a} = rH_0$  or

4. define mean “global” comoving velocity

$$\langle v \rangle_{\text{gal}}^{\text{glob}} := \frac{\int_0^r ds|_{t,\theta,\phi}}{t_0} = \frac{ar}{t_0}$$

# “Velocity” $v > c$ in FLRW?

■ low  $z$  Hubble law  $r \approx \frac{\beta c}{H_0}$ :  $r > 3h^{-1}$  Gpc  $\Rightarrow \beta > 1$ ?

■ possible definitions of “velocity”:

1. galaxy as fundamental observer, local velocity  $v_{\text{gal}} = 0$  by defn of comoving coords, or

2.  $v_{\text{gal}} < c$  by assumption ( $\parallel$  transport + Mink.) or

3.  $v_{\text{gal}}^{\text{glob}} := \frac{d(\int_0^r ds|_{t,\theta,\phi})}{dt} = \frac{d(ar)}{dt} = r \dot{a} = rH_0$  or

4. define mean “global” comoving velocity

$$\langle v \rangle_{\text{gal}}^{\text{glob}} := \frac{\int_0^r ds|_{t,\theta,\phi}}{t_0} = \frac{ar}{t_0} = \frac{r}{t_0}$$

# “Velocity” $v > c$ in FLRW?

■ low  $z$  Hubble law  $r \approx \frac{\beta c}{H_0}$ :  $r > 3h^{-1}$  Gpc  $\Rightarrow \beta > 1$ ?

■ possible definitions of “velocity”:

1. galaxy as fundamental observer, local velocity  $v_{\text{gal}} = 0$  by defn of comoving coords, or

2.  $v_{\text{gal}} < c$  by assumption ( $\parallel$  transport + Mink.) or

3.  $v_{\text{gal}}^{\text{glob}} := \frac{d(\int_0^r ds|_{t,\theta,\phi})}{dt} = \frac{d(ar)}{dt} = r \dot{a} = rH_0$  or

4.  $\langle v \rangle_{\text{gal}}^{\text{glob}} := \frac{\int_0^r ds|_{t,\theta,\phi}}{t_0} = \frac{ar}{t_0} = \frac{r}{t_0}$

■  $v_{\text{gal}}^{\text{glob}} \geq c$  when

# “Velocity” $v > c$ in FLRW?

■ low  $z$  Hubble law  $r \approx \frac{\beta c}{H_0}$ :  $r > 3h^{-1}$  Gpc  $\Rightarrow \beta > 1$ ?

■ possible definitions of “velocity”:

1. galaxy as fundamental observer, local velocity  $v_{\text{gal}} = 0$  by defn of comoving coords, or

2.  $v_{\text{gal}} < c$  by assumption ( $\parallel$  transport + Mink.) or

3.  $v_{\text{gal}}^{\text{glob}} := \frac{d(\int_0^r ds|_{t,\theta,\phi})}{dt} = \frac{d(ar)}{dt} = r \dot{a} = rH_0$  or

4.  $\langle v \rangle_{\text{gal}}^{\text{glob}} := \frac{\int_0^r ds|_{t,\theta,\phi}}{t_0} = \frac{ar}{t_0} = \frac{r}{t_0}$

■  $v_{\text{gal}}^{\text{glob}} \geq c$  when  $r \geq c/H_0 \approx 3h^{-1}$  Gpc

# “Velocity” $v > c$ in FLRW?

■ low  $z$  Hubble law  $r \approx \frac{\beta c}{H_0}$ :  $r > 3h^{-1}$  Gpc  $\Rightarrow \beta > 1$ ?

■ possible definitions of “velocity”:

1. galaxy as fundamental observer, local velocity  $v_{\text{gal}} = 0$  by defn of comoving coords, or

2.  $v_{\text{gal}} < c$  by assumption ( $\parallel$  transport + Mink.) or

3.  $v_{\text{gal}}^{\text{glob}} := \frac{d(\int_0^r ds|_{t,\theta,\phi})}{dt} = \frac{d(ar)}{dt} = r \dot{a} = rH_0$  or

4.  $\langle v \rangle_{\text{gal}}^{\text{glob}} := \frac{\int_0^r ds|_{t,\theta,\phi}}{t_0} = \frac{ar}{t_0} = \frac{r}{t_0}$

■  $v_{\text{gal}}^{\text{glob}} \geq c$  when  $r \geq c/H_0 \approx 3h^{-1}$  Gpc

■  $\langle v \rangle_{\text{gal}}^{\text{glob}} \geq c$  when  $r/t_0 \geq c$

# “Velocity” $v > c$ in FLRW?

■ low  $z$  Hubble law  $r \approx \frac{\beta c}{H_0}$ :  $r > 3h^{-1}$  Gpc  $\Rightarrow \beta > 1$ ?

■ possible definitions of “velocity”:

1. galaxy as fundamental observer, local velocity  $v_{\text{gal}} = 0$  by defn of comoving coords, or

2.  $v_{\text{gal}} < c$  by assumption ( $\parallel$  transport + Mink.) or

3.  $v_{\text{gal}}^{\text{glob}} := \frac{d(\int_0^r ds|_{t,\theta,\phi})}{dt} = \frac{d(ar)}{dt} = r \dot{a} = rH_0$  or

4.  $\langle v \rangle_{\text{gal}}^{\text{glob}} := \frac{\int_0^r ds|_{t,\theta,\phi}}{t_0} = \frac{ar}{t_0} = \frac{r}{t_0}$

■  $v_{\text{gal}}^{\text{glob}} \geq c$  when  $r \geq c/H_0 \approx 3h^{-1}$  Gpc

■  $\langle v \rangle_{\text{gal}}^{\text{glob}} \geq c$  when  $r \geq ct_0 \approx 4.5h^{-1}$  Gpc



# “Velocity” $v > c$ in FLRW?

- low  $z$  Hubble law  $r \approx \frac{\beta c}{H_0}$ :  $r > 3h^{-1}$  Gpc  $\Rightarrow \beta > 1$ ?
- possible definitions of “velocity”:
  1. galaxy as fundamental observer, local velocity  $v_{\text{gal}} = 0$  by defn of comoving coords, or
  2.  $v_{\text{gal}} < c$  by assumption ( $\parallel$  transport + Mink.) or
  3.  $v_{\text{gal}}^{\text{glob}} := \frac{d(\int_0^r ds|_{t,\theta,\phi})}{dt} = \frac{d(ar)}{dt} = r \dot{a} = rH_0$  or
  4.  $\langle v \rangle_{\text{gal}}^{\text{glob}} := \frac{\int_0^r ds|_{t,\theta,\phi}}{t_0} = \frac{ar}{t_0} = \frac{r}{t_0}$
- $v_{\text{gal}}^{\text{glob}} \geq c$  when  $r \geq c/H_0 \approx 3h^{-1}$  Gpc
- $\langle v \rangle_{\text{gal}}^{\text{glob}} \geq c$  when  $r \geq ct_0 \approx 4.5h^{-1}$  Gpc
- $v_{\text{gal}}^{\text{glob}}, \langle v \rangle_{\text{gal}}^{\text{glob}}$  = theoretical, global  $\neq$  velocity in tangent spacetime

# “Velocity” $v > c$ in FLRW?

■ low  $z$  Hubble law  $r \approx \frac{\beta c}{H_0}$ :  $r > 3h^{-1}$  Gpc  $\Rightarrow \beta > 1$ ?

■ possible definitions of “velocity”:

1. galaxy as fundamental observer, local velocity  $v_{\text{gal}} = 0$  by defn of comoving coords, or

2.  $v_{\text{gal}} < c$  by assumption ( $\parallel$  transport + Mink.) or

3.  $v_{\text{gal}}^{\text{glob}} := \frac{d(\int_0^r ds|_{t,\theta,\phi})}{dt} = \frac{d(ar)}{dt} = r \dot{a} = rH_0$  or

4.  $\langle v \rangle_{\text{gal}}^{\text{glob}} := \frac{\int_0^r ds|_{t,\theta,\phi}}{t_0} = \frac{ar}{t_0} = \frac{r}{t_0}$

■  $v_{\text{gal}}^{\text{glob}} \geq c$  when  $r \geq c/H_0 \approx 3h^{-1}$  Gpc

■  $\langle v \rangle_{\text{gal}}^{\text{glob}} \geq c$  when  $r \geq ct_0 \approx 4.5h^{-1}$  Gpc

■  $v_{\text{gal}}^{\text{glob}}, \langle v \rangle_{\text{gal}}^{\text{glob}}$  = theoretical, global  $\neq$  velocity in tangent spacetime

$\Rightarrow$  no conflict with locally Lorentzian (SR) spacetime

# Non-radial spatial geodesics

- What is the comoving distance between two objects at different celestial positions and different redshifts, for an arbitrary curvature (+, 0, -)?

# Non-radial spatial geodesics

- distances on the 2-sphere, embedded in  $\mathbb{R}^3$

# Non-radial spatial geodesics

- distances on the 2-sphere, embedded in  $\mathbb{R}^3$

$$x_i = R_C \cos \delta_i \cos \alpha_i$$

$$y_i = R_C \cos \delta_i \sin \alpha_i$$

$$w_i = R_C \sin \delta_i$$

# Non-radial spatial geodesics

- distances on the 2-sphere, embedded in  $\mathbb{R}^3$

$$x_i = R_C \cos \delta_i \cos \alpha_i$$

$$y_i = R_C \cos \delta_i \sin \alpha_i$$

$$w_i = R_C \sin \delta_i$$

$$\langle \mathbf{a}_1, \mathbf{a}_1 \rangle = x_1 x_2 + y_1 y_2 + w_1 w_2$$

# Non-radial spatial geodesics

- distances on the 2-sphere, embedded in  $\mathbb{R}^3$

$$x_i = R_C \cos \delta_i \cos \alpha_i$$

$$y_i = R_C \cos \delta_i \sin \alpha_i$$

$$w_i = R_C \sin \delta_i$$

$$\langle \mathbf{a}_1, \mathbf{a}_1 \rangle = x_1 x_2 + y_1 y_2 + w_1 w_2$$

- but also:

$$\langle \mathbf{a}_1, \mathbf{a}_1 \rangle = R^2 \cos \theta_{12}.$$

# Non-radial spatial geodesics

- distances on the 2-sphere, embedded in  $\mathbb{R}^3$

$$x_i = R_C \cos \delta_i \cos \alpha_i$$

$$y_i = R_C \cos \delta_i \sin \alpha_i$$

$$w_i = R_C \sin \delta_i$$

$$\langle \mathbf{a}_1, \mathbf{a}_1 \rangle = x_1 x_2 + y_1 y_2 + w_1 w_2$$

- but also:

$$\langle \mathbf{a}_1, \mathbf{a}_1 \rangle = R^2 \cos \theta_{12}.$$

- a distance in  $S^2 = \text{arc-length in } \mathbb{R}^3:$



# Non-radial spatial geodesics

- distances on the 2-sphere, embedded in  $\mathbb{R}^3$

$$x_i = R_C \cos \delta_i \cos \alpha_i$$

$$y_i = R_C \cos \delta_i \sin \alpha_i$$

$$w_i = R_C \sin \delta_i$$

$$\langle \mathbf{a}_1, \mathbf{a}_2 \rangle = x_1 x_2 + y_1 y_2 + w_1 w_2$$

- but also:

$$\langle \mathbf{a}_1, \mathbf{a}_2 \rangle = R^2 \cos \theta_{12}.$$

- a distance in  $S^2 =$  arc-length in  $\mathbb{R}^3$ :

$$\chi_{12} = R_C \theta_{12} = R_C \cos^{-1} \left[ \langle \mathbf{a}_1, \mathbf{a}_2 \rangle / R_C^2 \right]$$



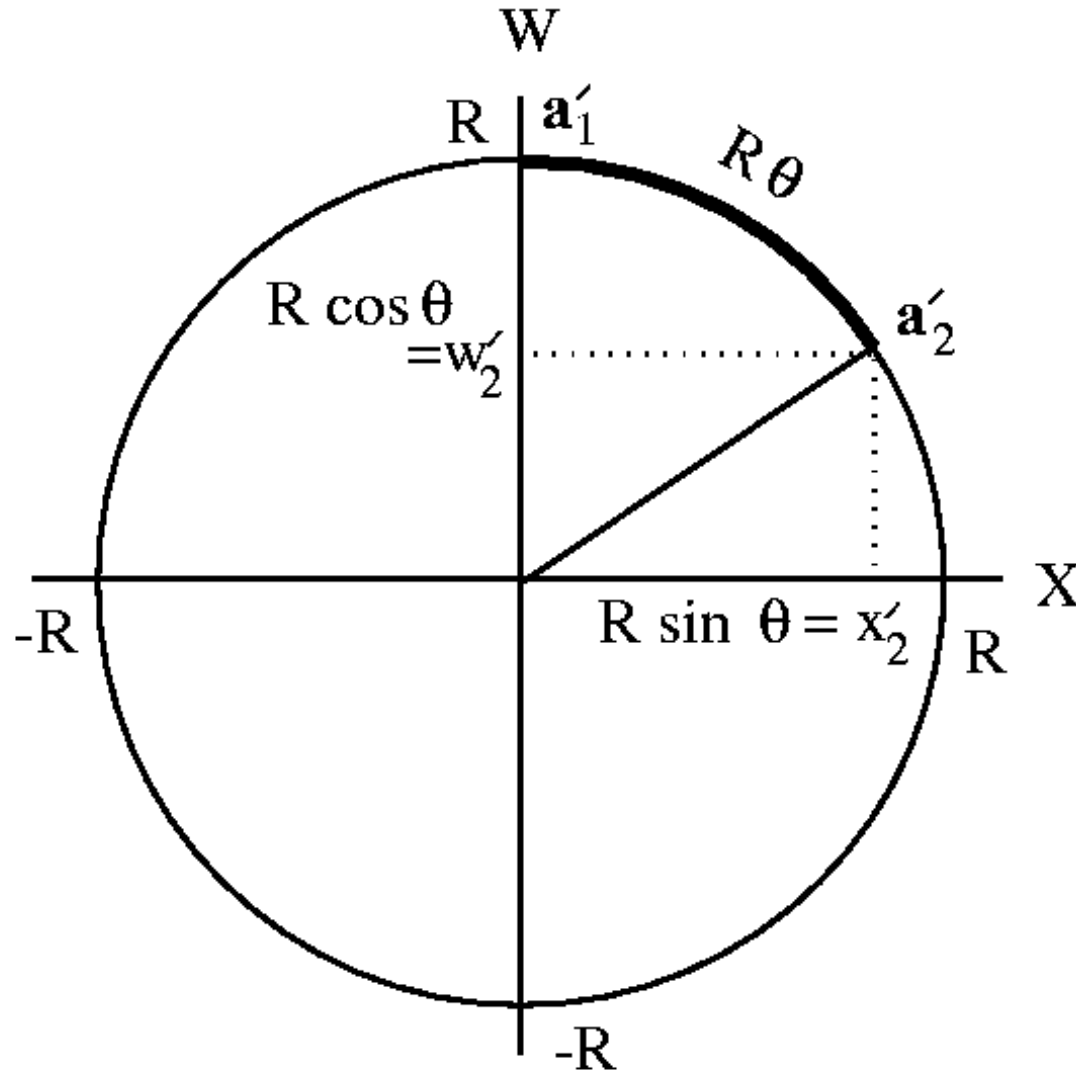
# Non-radial spatial geodesics



- positive curvature

# Non-radial spatial geodesics

- positive curvature

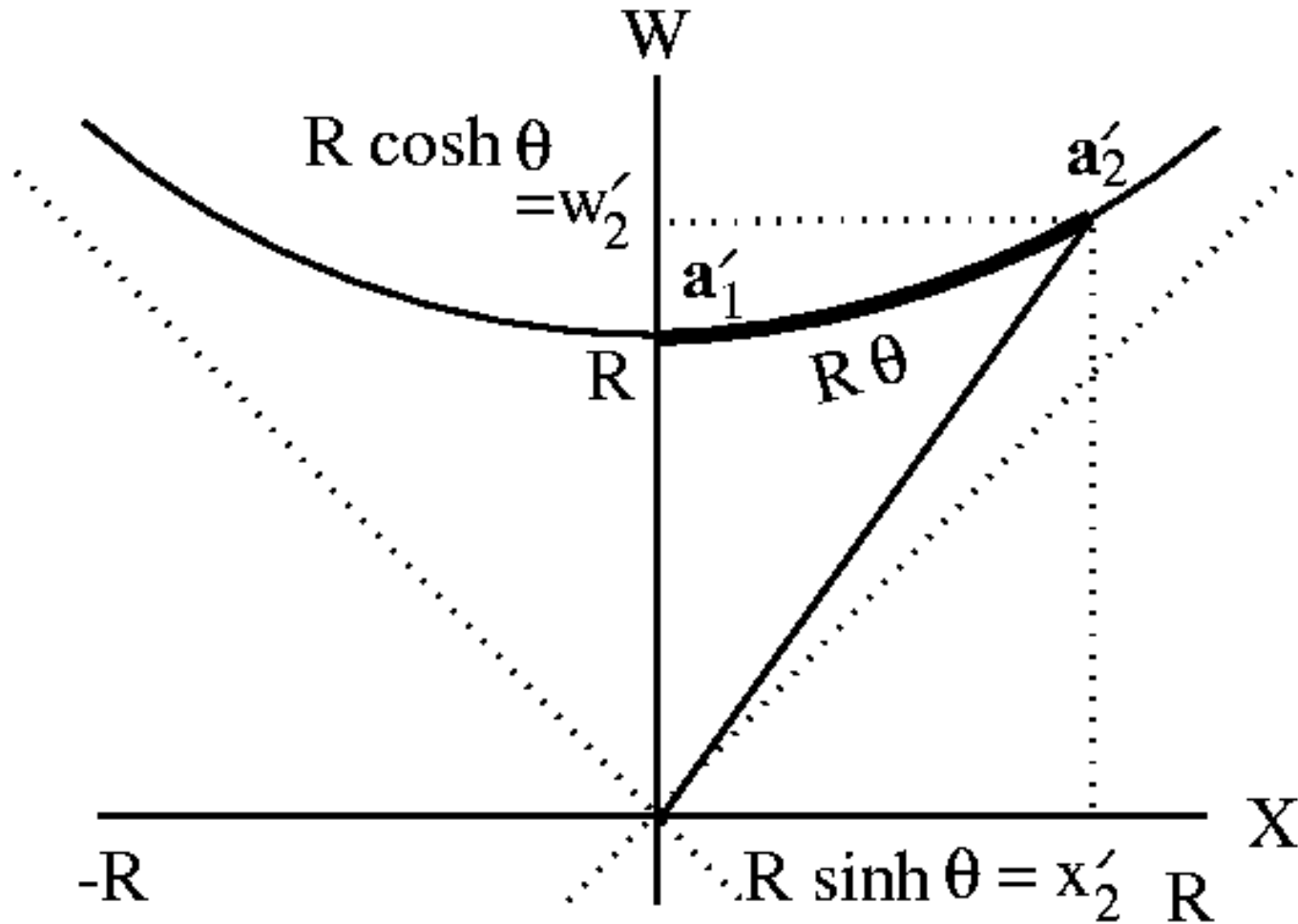


# Non-radial spatial geodesics

- negative curvature

# Non-radial spatial geodesics

- negative curvature



distances on  $S^3 \subset \mathbb{R}^4$  or  $H^3 \subset M^4$

$$\Sigma(\chi_i) := \begin{cases} R_C \sinh(\chi_i/R_C) & k < 0 \\ \chi_i & k = 0 \\ R_C \sin(\chi_i/R_C) & k > 0 \end{cases}$$

# distances on $S^3 \subset \mathbb{R}^4$ or $H^3 \subset M^4$

$$\Sigma(\chi_i) := \begin{cases} R_C \sinh(\chi_i/R_C) & k < 0 \\ \chi_i & k = 0 \\ R_C \sin(\chi_i/R_C) & k > 0 \end{cases}$$

$$x_i = \Sigma(\chi_i) \cos \delta_i \cos \alpha_i$$

$$y_i = \Sigma(\chi_i) \cos \delta_i \sin \alpha_i$$

$$z_i = \Sigma(\chi_i) \sin \delta_i$$

$$w_i = \begin{cases} R_C \cosh(\chi_i/R_C) & k < 0 \\ 0 & k = 0 \\ R_C \cos(\chi_i/R_C) & k > 0 \end{cases}$$

# distances on $S^3 \subset \mathbb{R}^4$ or $H^3 \subset M^4$

$$\Sigma(\chi_i) := \begin{cases} R_C \sinh(\chi_i/R_C) & k < 0 \\ \chi_i & k = 0 \\ R_C \sin(\chi_i/R_C) & k > 0 \end{cases}$$

$$x_i = \Sigma(\chi_i) \cos \delta_i \cos \alpha_i$$

$$y_i = \Sigma(\chi_i) \cos \delta_i \sin \alpha_i$$

$$z_i = \Sigma(\chi_i) \sin \delta_i$$

$$w_i = \begin{cases} R_C \cosh(\chi_i/R_C) & k < 0 \\ 0 & k = 0 \\ R_C \cos(\chi_i/R_C) & k > 0 \end{cases}$$

■ metric on  $S^3$  (or  $\mathbb{R}^3$  or  $H^3$ ):

$$ds^2 = \begin{cases} (k/|k|) (dx^2 + dy^2 + dz^2) + dw^2 & k \neq 0 \\ dx^2 + dy^2 + dz^2 & k = 0 \end{cases}$$



# distances on $S^3 \subset \mathbb{R}^4$ or $H^3 \subset M^4$

$$\Sigma(\chi_i) := \begin{cases} R_C \sinh(\chi_i/R_C) & k < 0 \\ \chi_i & k = 0 \\ R_C \sin(\chi_i/R_C) & k > 0 \end{cases}$$

$$x_i = \Sigma(\chi_i) \cos \delta_i \cos \alpha_i$$

$$y_i = \Sigma(\chi_i) \cos \delta_i \sin \alpha_i$$

$$z_i = \Sigma(\chi_i) \sin \delta_i$$

$$w_i = \begin{cases} R_C \cosh(\chi_i/R_C) & k < 0 \\ 0 & k = 0 \\ R_C \cos(\chi_i/R_C) & k > 0 \end{cases}$$

inner product:

$$\langle \mathbf{a}_1, \mathbf{a}_2 \rangle \equiv \begin{cases} (k/|k|) (x_1x_2 + y_1y_2 + z_1z_2) + w_1w_2 & k \neq 0 \\ x_1x_2 + y_1y_2 + z_1z_2 & k = 0 \end{cases}$$

# distances on $S^3 \subset \mathbb{R}^4$ or $H^3 \subset M^4$

$$\Sigma(\chi_i) := \begin{cases} R_C \sinh(\chi_i/R_C) & k < 0 \\ \chi_i & k = 0 \\ R_C \sin(\chi_i/R_C) & k > 0 \end{cases}$$

$$x_i = \Sigma(\chi_i) \cos \delta_i \cos \alpha_i$$

$$y_i = \Sigma(\chi_i) \cos \delta_i \sin \alpha_i$$

$$z_i = \Sigma(\chi_i) \sin \delta_i$$

$$w_i = \begin{cases} R_C \cosh(\chi_i/R_C) & k < 0 \\ 0 & k = 0 \\ R_C \cos(\chi_i/R_C) & k > 0 \end{cases}$$

$\Rightarrow$

$$\chi_{12} = \begin{cases} R_C \cosh^{-1} [\langle \mathbf{a}_1, \mathbf{a}_2 \rangle / R_C^2] & k < 0 \\ \sqrt{\langle \mathbf{a}_1 - \mathbf{a}_2, \mathbf{a}_1 - \mathbf{a}_2 \rangle} & k = 0 \\ R_C \cos^{-1} [\langle \mathbf{a}_1, \mathbf{a}_2 \rangle / R_C^2] & k > 0 \end{cases}$$

# distances on $S^3 \subset \mathbb{R}^4$ or $H^3 \subset M^4$

- $$x_i = \Sigma(\chi_i) \cos \delta_i \cos \alpha_i$$

$$y_i = \Sigma(\chi_i) \cos \delta_i \sin \alpha_i$$

$$z_i = \Sigma(\chi_i) \sin \delta_i$$

$$w_i = \begin{cases} R_C \cosh(\chi_i/R_C) & k < 0 \\ 0 & k = 0 \\ R_C \cos(\chi_i/R_C) & k > 0 \end{cases}$$

$$\Rightarrow \chi_{12} = \begin{cases} R_C \cosh^{-1} [\langle \mathbf{a}_1, \mathbf{a}_2 \rangle / R_C^2] & k < 0 \\ \sqrt{\langle \mathbf{a}_1 - \mathbf{a}_2, \mathbf{a}_1 - \mathbf{a}_2 \rangle} & k = 0 \\ R_C \cos^{-1} [\langle \mathbf{a}_1, \mathbf{a}_2 \rangle / R_C^2] & k > 0 \end{cases}$$

- a distance in  $S^3$  is an arc-length in  $\mathbb{R}^4$

# distances on $S^3 \subset \mathbb{R}^4$ or $H^3 \subset M^4$

$$\begin{aligned}
 x_i &= \Sigma(\chi_i) \cos \delta_i \cos \alpha_i \\
 y_i &= \Sigma(\chi_i) \cos \delta_i \sin \alpha_i \\
 z_i &= \Sigma(\chi_i) \sin \delta_i \\
 w_i &= \begin{cases} R_C \cosh(\chi_i/R_C) & k < 0 \\ 0 & k = 0 \\ R_C \cos(\chi_i/R_C) & k > 0 \end{cases}
 \end{aligned}$$

$$\Rightarrow \chi_{12} = \begin{cases} R_C \cosh^{-1} [\langle \mathbf{a}_1, \mathbf{a}_2 \rangle / R_C^2] & k < 0 \\ \sqrt{\langle \mathbf{a}_1 - \mathbf{a}_2, \mathbf{a}_1 - \mathbf{a}_2 \rangle} & k = 0 \\ R_C \cos^{-1} [\langle \mathbf{a}_1, \mathbf{a}_2 \rangle / R_C^2] & k > 0 \end{cases}$$

■ a distance in  $S^3$  is an arc-length in  $\mathbb{R}^4$

■ a distance in  $H^3$  is an arc-length in  $M^4$

# distances on $S^3 \subset \mathbb{R}^4$ or $H^3 \subset M^4$

$$\begin{aligned}
 x_i &= \Sigma(\chi_i) \cos \delta_i \cos \alpha_i \\
 y_i &= \Sigma(\chi_i) \cos \delta_i \sin \alpha_i \\
 z_i &= \Sigma(\chi_i) \sin \delta_i \\
 w_i &= \begin{cases} R_C \cosh(\chi_i/R_C) & k < 0 \\ 0 & k = 0 \\ R_C \cos(\chi_i/R_C) & k > 0 \end{cases}
 \end{aligned}$$

$$\Rightarrow \chi_{12} = \begin{cases} R_C \cosh^{-1} [\langle \mathbf{a}_1, \mathbf{a}_2 \rangle / R_C^2] & k < 0 \\ \sqrt{\langle \mathbf{a}_1 - \mathbf{a}_2, \mathbf{a}_1 - \mathbf{a}_2 \rangle} & k = 0 \\ R_C \cos^{-1} [\langle \mathbf{a}_1, \mathbf{a}_2 \rangle / R_C^2] & k > 0 \end{cases}$$

■ a distance in  $S^3$  is an arc-length in  $\mathbb{R}^4$  [arXiv:astro-ph/0102099](https://arxiv.org/abs/astro-ph/0102099)

■ a distance in  $H^3$  is an arc-length in  $M^4$

# Cosmic topology: definitions

- 3 intuitive methods

# Cosmic topology: definitions

- 3 intuitive methods
- $M$  = universal covering space

# Cosmic topology: definitions

- 3 intuitive methods
- $M$  = universal covering space = “apparent space”



# Cosmic topology: definitions

- 3 intuitive methods

- $M$  = universal covering space = “apparent space”

- $M = \begin{cases} \mathbb{H}^3 & k < 0 \end{cases}$

# Cosmic topology: definitions

- 3 intuitive methods

- $M$  = universal covering space = “apparent space”

- $$M = \begin{cases} \mathbb{H}^3 & k < 0 \\ \mathbb{R}^3 & k = 0 \end{cases}$$

# Cosmic topology: definitions

- 3 intuitive methods

- $M$  = universal covering space = “apparent space”

- $$M = \begin{cases} \mathbb{H}^3 & k < 0 \\ \mathbb{R}^3 & k = 0 \\ \mathbb{S}^3 & k > 0 \end{cases}$$

# Cosmic topology: definitions

- 3 intuitive methods
- $M$  = universal covering space = “apparent space”
- $M = \begin{cases} \mathbb{H}^3 & k < 0 \\ \mathbb{R}^3 & k = 0 \\ \mathbb{S}^3 & k > 0 \end{cases}$
- $\Gamma$  = a group of holonomies (isometries) of  $M$

# Cosmic topology: definitions

- 3 intuitive methods
- $M$  = universal covering space = “apparent space”
- $M = \begin{cases} \mathbb{H}^3 & k < 0 \\ \mathbb{R}^3 & k = 0 \\ \mathbb{S}^3 & k > 0 \end{cases}$
- $\Gamma$  = a group of holonomies (isometries) of  $M$
- the 3-manifold

# Cosmic topology: definitions

- 3 intuitive methods
- $M$  = universal covering space = “apparent space”
- $M = \begin{cases} \mathbb{H}^3 & k < 0 \\ \mathbb{R}^3 & k = 0 \\ \mathbb{S}^3 & k > 0 \end{cases}$
- $\Gamma$  = a group of holonomies (isometries) of  $M$
- the 3-space for FLRW cosmology is  $M/\Gamma$

# Cosmic topology: definitions

- 3 intuitive methods
- $M$  = universal covering space = “apparent space”
- $M = \begin{cases} \mathbb{H}^3 & k < 0 \\ \mathbb{R}^3 & k = 0 \\ \mathbb{S}^3 & k > 0 \end{cases}$
- $\Gamma$  = a group of holonomies (isometries) of  $M$
- the 3-space for FLRW cosmology is  $M/\Gamma$
- each of the 3 ways of thinking of  $M/\Gamma$  has advantages and disadvantages

# Cosmic topology: definitions

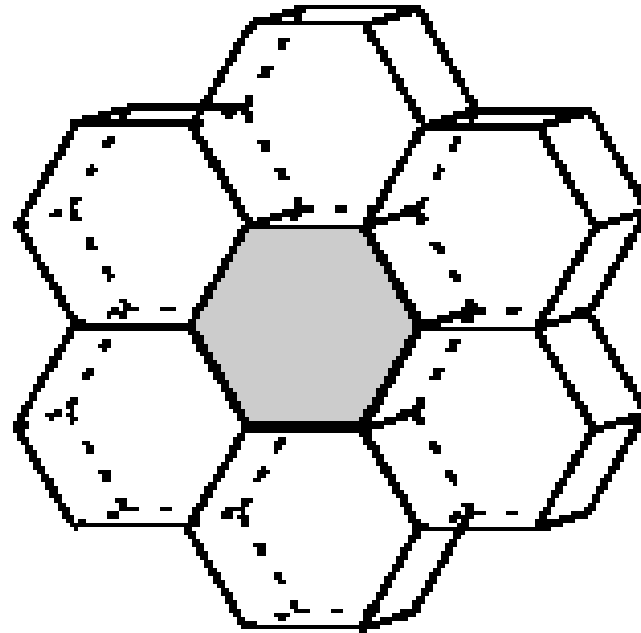
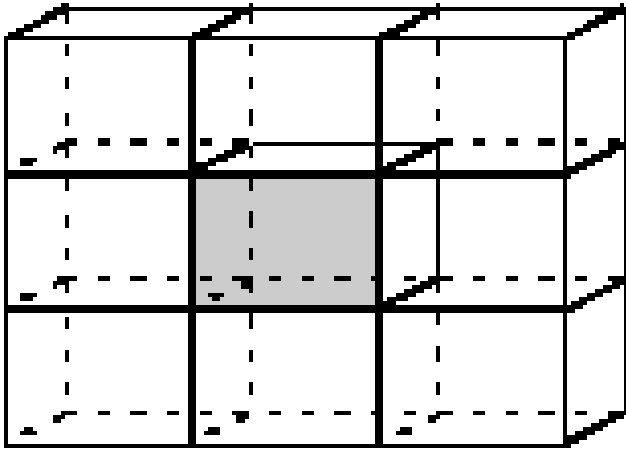
- 3 intuitive methods
- $M$  = universal covering space = “apparent space”
- $M = \begin{cases} \mathbb{H}^3 & k < 0 \\ \mathbb{R}^3 & k = 0 \\ \mathbb{S}^3 & k > 0 \end{cases}$
- $\Gamma$  = a group of holonomies (isometries) of  $M$
- the 3-space for FLRW cosmology is  $M/\Gamma$
- each of the 3 ways of thinking of  $M/\Gamma$  has advantages and disadvantages
- fundamental domain (FD) is not unique;



# Cosmic topology: definitions

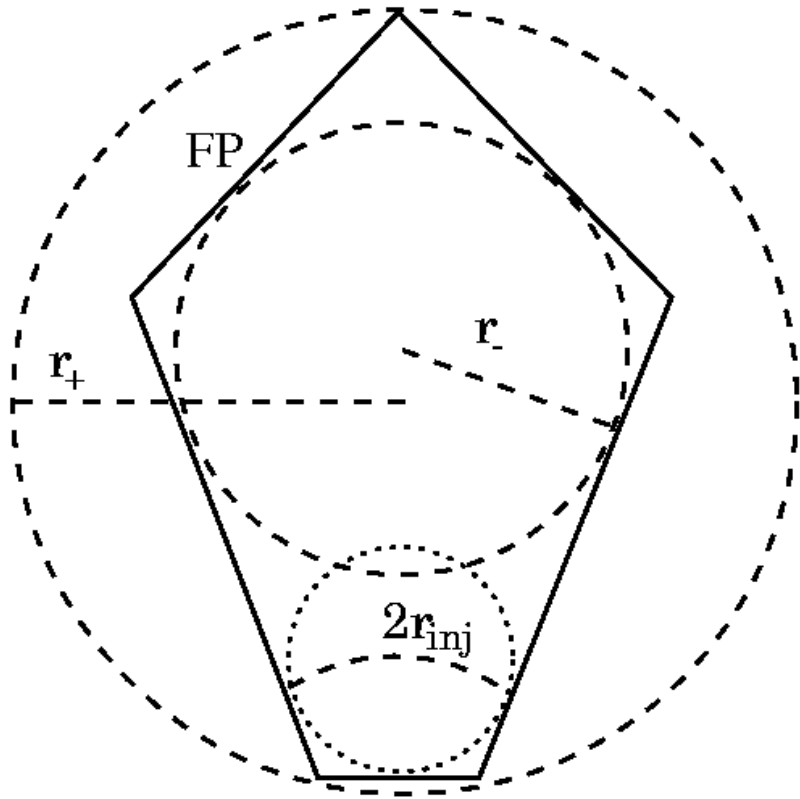
- 3 intuitive methods
- $M$  = universal covering space = “apparent space”
- $M = \begin{cases} \mathbb{H}^3 & k < 0 \\ \mathbb{R}^3 & k = 0 \\ \mathbb{S}^3 & k > 0 \end{cases}$
- $\Gamma$  = a group of holonomies (isometries) of  $M$
- the 3-space for FLRW cosmology is  $M/\Gamma$
- each of the 3 ways of thinking of  $M/\Gamma$  has advantages and disadvantages
- fundamental domain (FD) is not unique; shape of FD may be non-unique

# Cosmic topology: definitions



3D flat examples [arXiv:astro-ph/9901364](https://arxiv.org/abs/astro-ph/9901364)

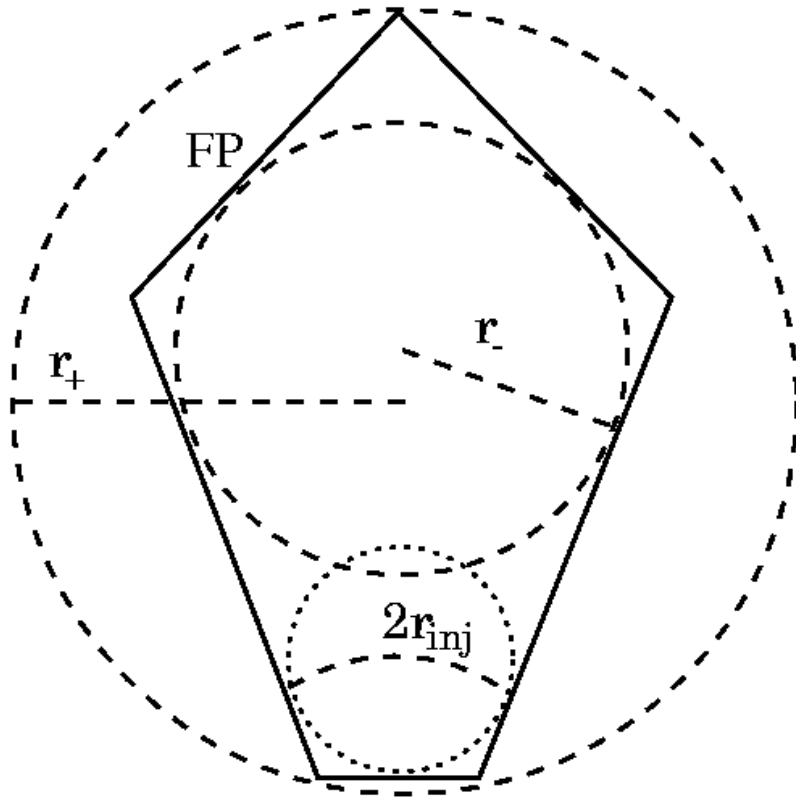
# Cosmic topology: definitions



# Cosmic topology: definitions

size of universe:

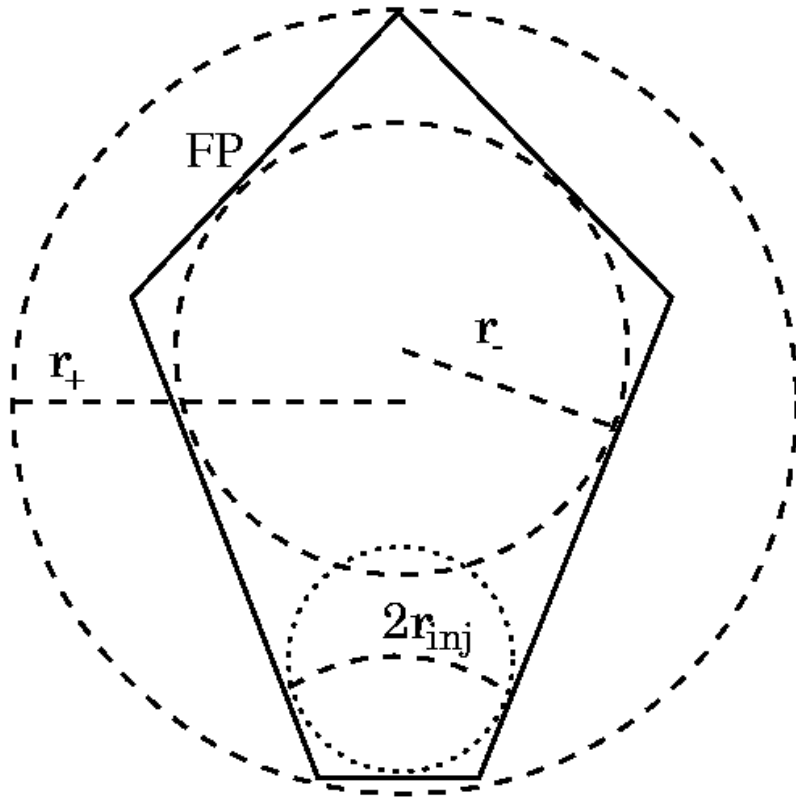
- $r_-$  : biggest sphere *inside* FD



# Cosmic topology: definitions

size of universe:

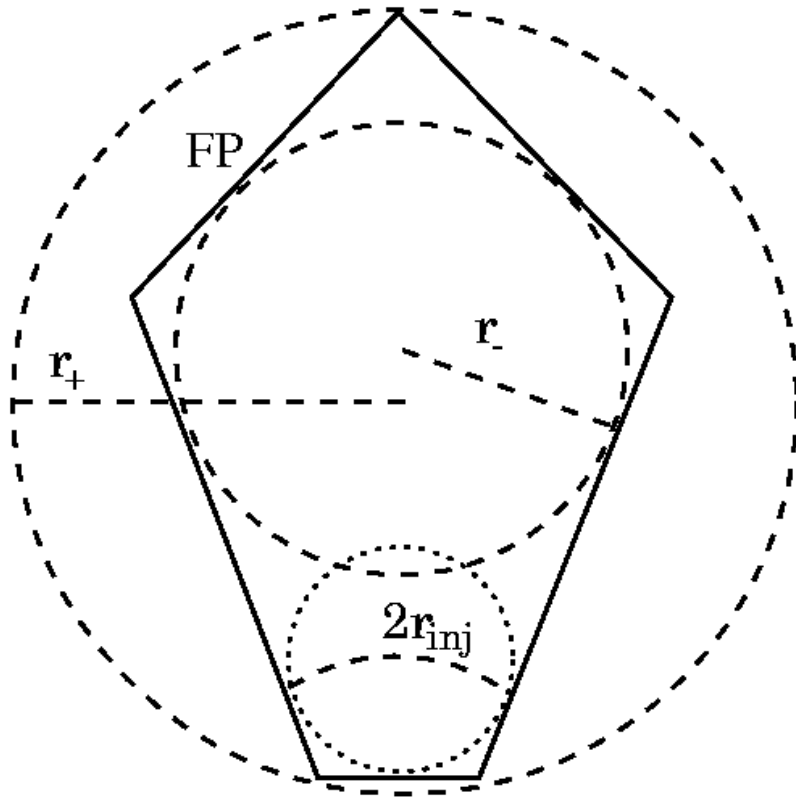
- $r_-$  : biggest sphere *inside* FD
- $r_+$  : smallest sphere *containing* FD



# Cosmic topology: definitions

size of universe:

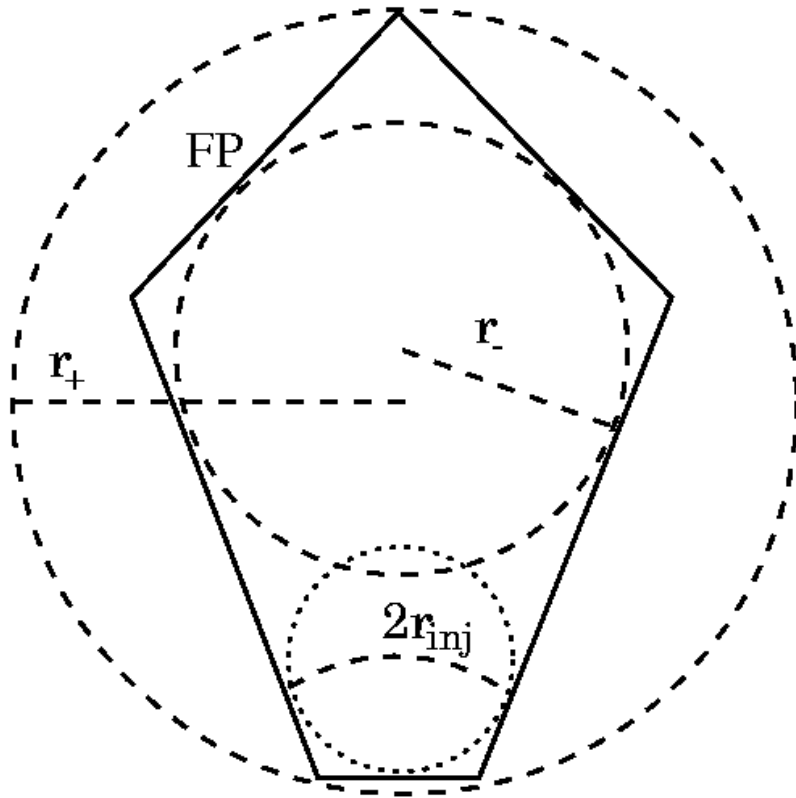
- $r_-$  : biggest sphere *inside* FD
- $r_+$  : smallest sphere *containing* FD
- $2r_{\text{inj}}$  : smallest closed spatial geodesic



# Cosmic topology: definitions

size of universe:

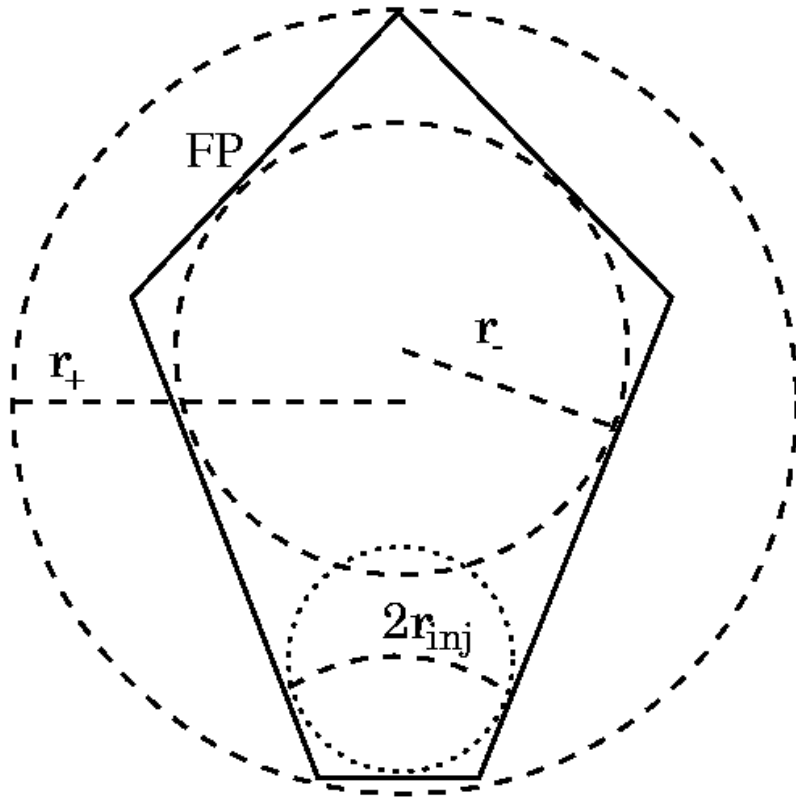
- $r_-$  : biggest sphere *inside* FD
- $r_+$  : smallest sphere *containing* FD
- $2r_{\text{inj}}$  : smallest closed spatial geodesic
- $V_{\text{FD}}^{1/3}$  volume cube root



# Cosmic topology: definitions

size of universe:

- $r_-$  : biggest sphere *inside* FD
- $r_+$  : smallest sphere *containing* FD
- $2r_{\text{inj}}$  : smallest closed spatial geodesic
- $V_{\text{FD}}^{1/3}$  volume cube root
- $r_- \leq r_+$  always

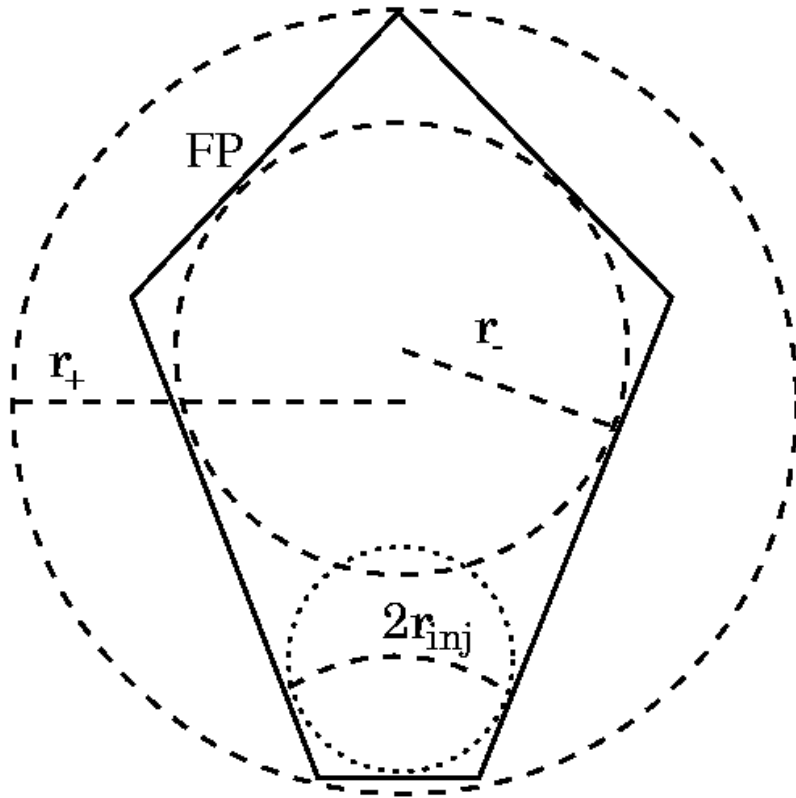




# Cosmic topology: definitions

size of universe:

- $r_-$  : biggest sphere *inside* FD
- $r_+$  : smallest sphere *containing* FD
- $2r_{\text{inj}}$  : smallest closed spatial geodesic
- $V_{\text{FD}}^{1/3}$  volume cube root
- $r_- \leq r_+$  always
- $r_{\text{inj}} < r_-$  or  $r_{\text{inj}} \ll V_{\text{FD}}^{1/3}$  possible





# Families of const $k$ 3-spaces



- flat 3-spaces: 18



# Families of const $k$ 3-spaces



- flat 3-spaces: 18
  - ◆ some finite,

# Families of const $k$ 3-spaces

- flat 3-spaces: 18
  - ◆ some finite, e.g.  $T^3$ ,

# Families of const $k$ 3-spaces

- flat 3-spaces: 18
  - ◆ some finite, e.g.  $T^3$ , some infinite,

# Families of const $k$ 3-spaces

- flat 3-spaces: 18

- ◆ some finite, e.g.  $T^3$ , some infinite, e.g.  $\mathbb{R}^3$

# Families of const $k$ 3-spaces

- flat 3-spaces: 18
  - ◆ some finite, e.g.  $T^3$ , some infinite, e.g.  $\mathbb{R}^3$
- spherical: [arXiv:gr-qc/0106033](https://arxiv.org/abs/gr-qc/0106033) Gausmann et al.

# Families of const $k$ 3-spaces

- flat 3-spaces: 18
  - ◆ some finite, e.g.  $T^3$ , some infinite, e.g.  $\mathbb{R}^3$
- spherical: [arXiv:gr-qc/0106033](https://arxiv.org/abs/gr-qc/0106033) Gausmann et al.
  - ◆  $S^3/\Gamma$ —countably infinite set of 3-spaces



# Families of const $k$ 3-spaces

- flat 3-spaces: 18
  - ◆ some finite, e.g.  $T^3$ , some infinite, e.g.  $\mathbb{R}^3$
- spherical: [arXiv:gr-qc/0106033](https://arxiv.org/abs/gr-qc/0106033) Gausmann et al.
  - ◆  $S^3/\Gamma$ —countably infinite set of 3-spaces
  - ◆ all finite (“compact”)

# Families of const $k$ 3-spaces

- flat 3-spaces: 18
  - ◆ some finite, e.g.  $T^3$ , some infinite, e.g.  $\mathbb{R}^3$
- spherical: [arXiv:gr-qc/0106033](https://arxiv.org/abs/gr-qc/0106033) Gausmann et al.
  - ◆  $S^3/\Gamma$ —countably infinite set of 3-spaces
  - ◆ all finite (“compact”)  $\Rightarrow \Gamma =$  finite group

# Families of const $k$ 3-spaces

- flat 3-spaces: 18
  - ◆ some finite, e.g.  $T^3$ , some infinite, e.g.  $\mathbb{R}^3$
- spherical: [arXiv:gr-qc/0106033](https://arxiv.org/abs/gr-qc/0106033) Gausmann et al.
  - ◆  $S^3/\Gamma$ —countably infinite set of 3-spaces
  - ◆ all finite (“compact”)  $\Rightarrow \Gamma =$  finite group
  - ◆ completely classified (Threlfall & Seifert 1930)

# Families of const $k$ 3-spaces

- flat 3-spaces: 18
  - ◆ some finite, e.g.  $T^3$ , some infinite, e.g.  $\mathbb{R}^3$
- spherical: [arXiv:gr-qc/0106033](https://arxiv.org/abs/gr-qc/0106033) Gausmann et al.
  - ◆  $S^3/\Gamma$ —countably infinite set of 3-spaces
  - ◆ all finite (“compact”)  $\Rightarrow \Gamma =$  finite group
  - ◆ completely classified (Threlfall & Seifert 1930)
  - ◆  $S^3, S^3/\mathbb{Z}_2,$

# Families of const $k$ 3-spaces

- flat 3-spaces: 18
  - ◆ some finite, e.g.  $T^3$ , some infinite, e.g.  $\mathbb{R}^3$
- spherical: [arXiv:gr-qc/0106033](https://arxiv.org/abs/gr-qc/0106033) Gausmann et al.
  - ◆  $S^3/\Gamma$ —countably infinite set of 3-spaces
  - ◆ all finite (“compact”)  $\Rightarrow \Gamma =$  finite group
  - ◆ completely classified (Threlfall & Seifert 1930)
  - ◆  $S^3$ ,  $S^3/\mathbb{Z}_2$ , lens spaces ,

# Families of const $k$ 3-spaces

- flat 3-spaces: 18
  - ◆ some finite, e.g.  $T^3$ , some infinite, e.g.  $\mathbb{R}^3$
- spherical: [arXiv:gr-qc/0106033](https://arxiv.org/abs/gr-qc/0106033) Gausmann et al.
  - ◆  $S^3/\Gamma$ —countably infinite set of 3-spaces
  - ◆ all finite (“compact”)  $\Rightarrow \Gamma =$  finite group
  - ◆ completely classified (Threlfall & Seifert 1930)
  - ◆  $S^3$ ,  $S^3/\mathbb{Z}_2$ , lens spaces, “well-proportioned” spaces, ...

# Families of const $k$ 3-spaces

- flat 3-spaces: 18
  - ◆ some finite, e.g.  $T^3$ , some infinite, e.g.  $\mathbb{R}^3$
- spherical: [arXiv:gr-qc/0106033](https://arxiv.org/abs/gr-qc/0106033) Gausmann et al.
  - ◆  $S^3/\Gamma$ —countably infinite set of 3-spaces
  - ◆ all finite (“compact”)  $\Rightarrow \Gamma =$  finite group
  - ◆ completely classified (Threlfall & Seifert 1930)
  - ◆  $S^3$ ,  $S^3/\mathbb{Z}_2$ , lens spaces, “well-proportioned” spaces, ...
  - ◆ w:Poincaré Conjecture “Every simply connected, closed 3-manifold is homeomorphic to the 3-sphere.”  
w:Grigori Perelman, [arXiv:math/0211159](https://arxiv.org/abs/math/0211159) + [arXiv:math/0303109](https://arxiv.org/abs/math/0303109) + [arXiv:math/0307245](https://arxiv.org/abs/math/0307245)

# Families of const $k$ 3-spaces

- flat 3-spaces: 18
  - ◆ some finite, e.g.  $T^3$ , some infinite, e.g.  $\mathbb{R}^3$
- spherical: [arXiv:gr-qc/0106033](https://arxiv.org/abs/gr-qc/0106033) Gausmann et al.
  - ◆  $S^3/\Gamma$ —countably infinite set of 3-spaces
  - ◆ all finite (“compact”)  $\Rightarrow \Gamma =$  finite group
  - ◆ completely classified (Threlfall & Seifert 1930)
  - ◆  $S^3$ ,  $S^3/\mathbb{Z}_2$ , lens spaces, “well-proportioned” spaces, ...
- hyperbolic:
  - ◆ countably infinite superset known



# Families of const $k$ 3-spaces

- flat 3-spaces: 18
  - ◆ some finite, e.g.  $T^3$ , some infinite, e.g.  $\mathbb{R}^3$
- spherical: [arXiv:gr-qc/0106033](https://arxiv.org/abs/gr-qc/0106033) Gausmann et al.
  - ◆  $S^3/\Gamma$ —countably infinite set of 3-spaces
  - ◆ all finite (“compact”)  $\Rightarrow \Gamma =$  finite group
  - ◆ completely classified (Threlfall & Seifert 1930)
  - ◆  $S^3$ ,  $S^3/\mathbb{Z}_2$ , lens spaces, “well-proportioned” spaces, ...
- hyperbolic:
  - ◆ countably infinite superset known
  - ◆ some finite, some infinite

# Families of const $k$ 3-spaces

- flat 3-spaces: 18
  - ◆ some finite, e.g.  $T^3$ , some infinite, e.g.  $\mathbb{R}^3$
- spherical: [arXiv:gr-qc/0106033](https://arxiv.org/abs/gr-qc/0106033) Gausmann et al.
  - ◆  $S^3/\Gamma$ —countably infinite set of 3-spaces
  - ◆ all finite (“compact”)  $\Rightarrow \Gamma =$  finite group
  - ◆ completely classified (Threlfall & Seifert 1930)
  - ◆  $S^3$ ,  $S^3/\mathbb{Z}_2$ , lens spaces, “well-proportioned” spaces, ...
- hyperbolic:
  - ◆ countably infinite superset known
  - ◆ some finite, some infinite
  - ◆ active research area, e.g. [arXiv:0705.4325](https://arxiv.org/abs/0705.4325) min. vol.

# Cosmic topol: top. accel.

- cosmic topology theory:

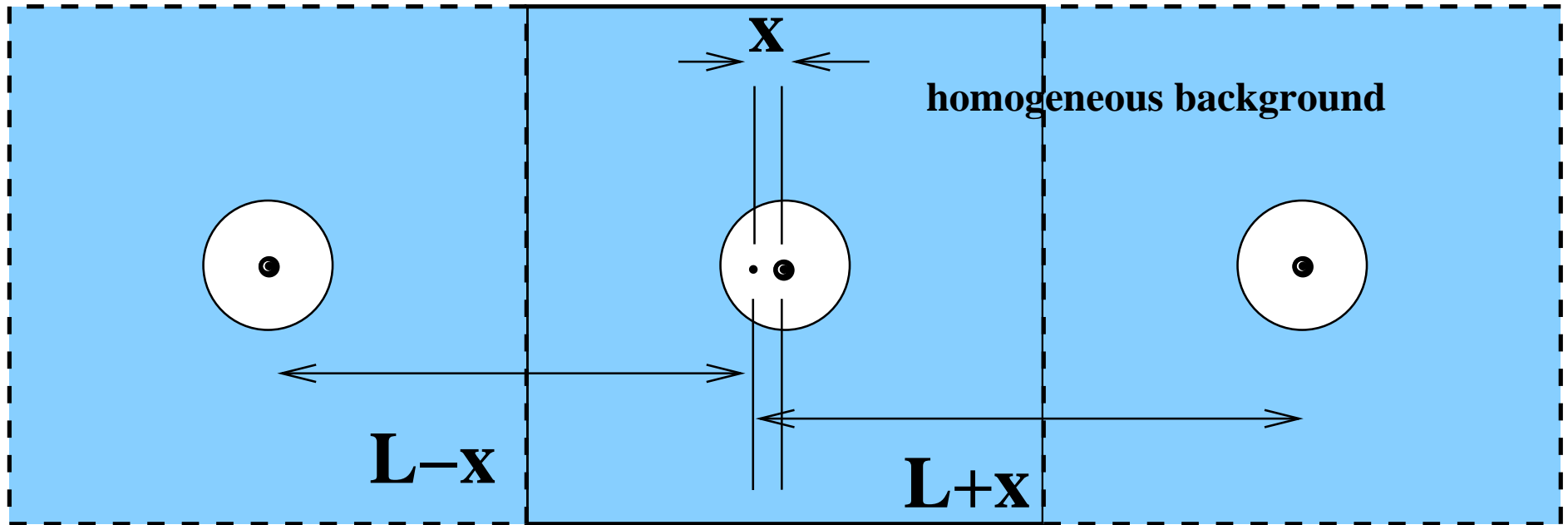
# Cosmic topol: top. accel.

- cosmic topology theory:
- either very theoretical (quantum gravity)

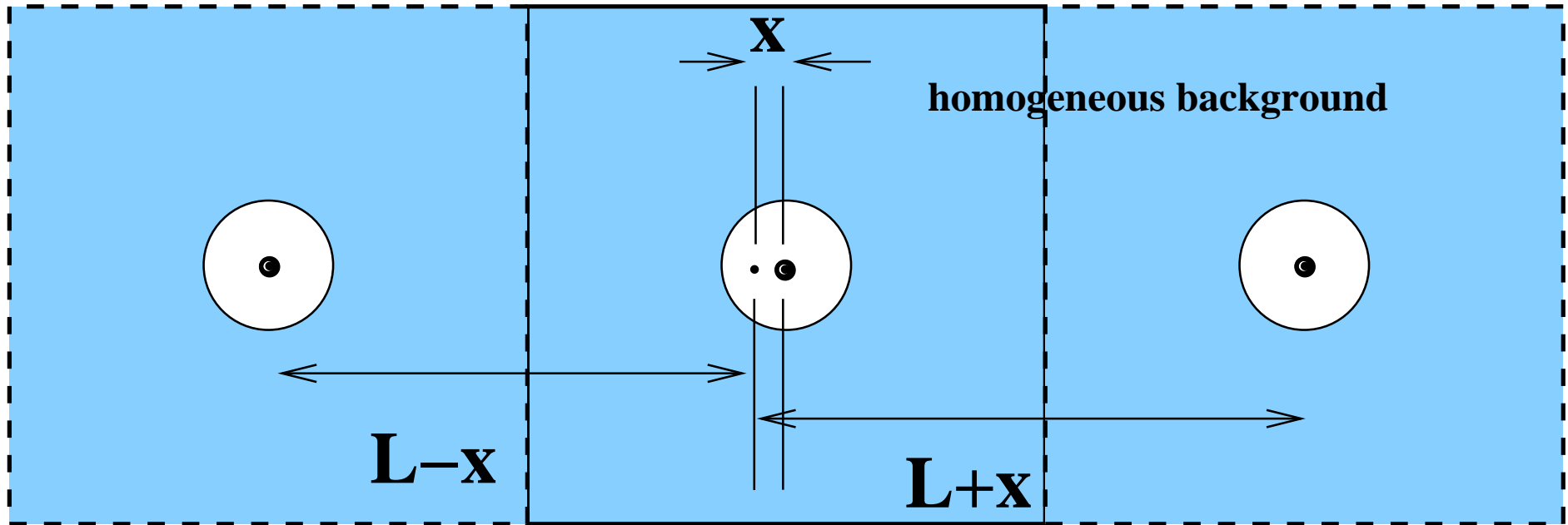
# Cosmic topol: top. accel.

- cosmic topology theory:
- either very theoretical (quantum gravity)
- or very simple (topological acceleration)

# Cosmic topol: top. accel.

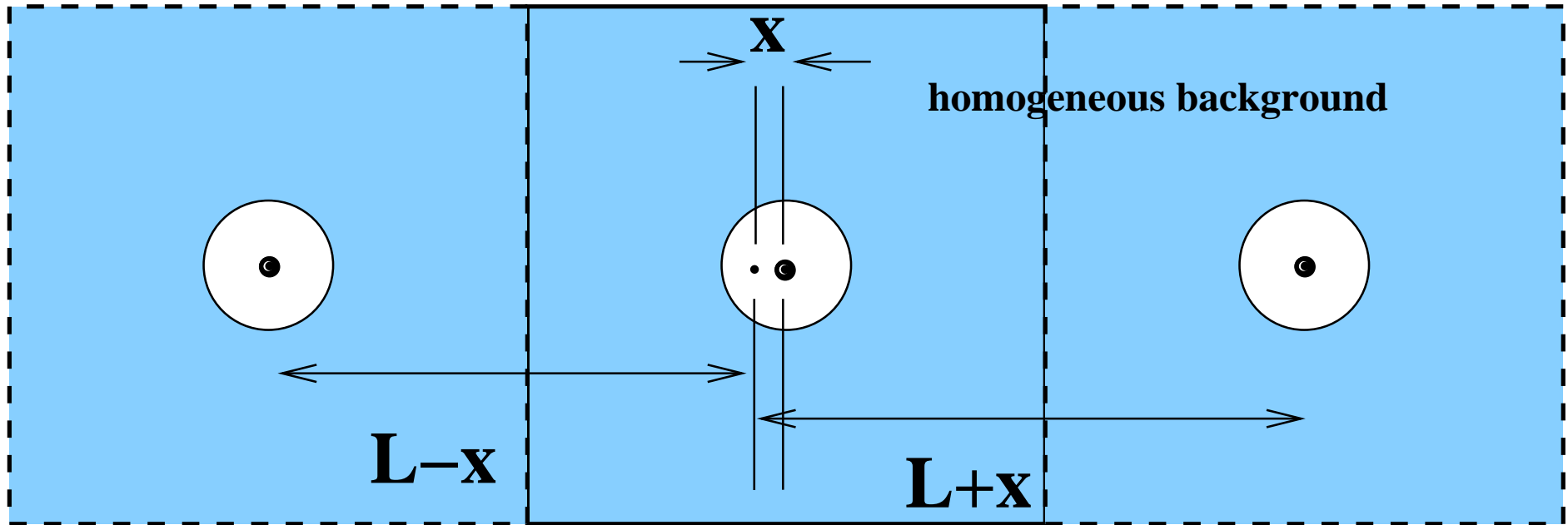


# Cosmic topol: top. accel.



$$\ddot{x} \approx -G \frac{m}{x^2} + Gm \left[ \frac{1}{(L-x)^2} - \frac{1}{(L+x)^2} \right]$$

# Cosmic topol: top. accel.

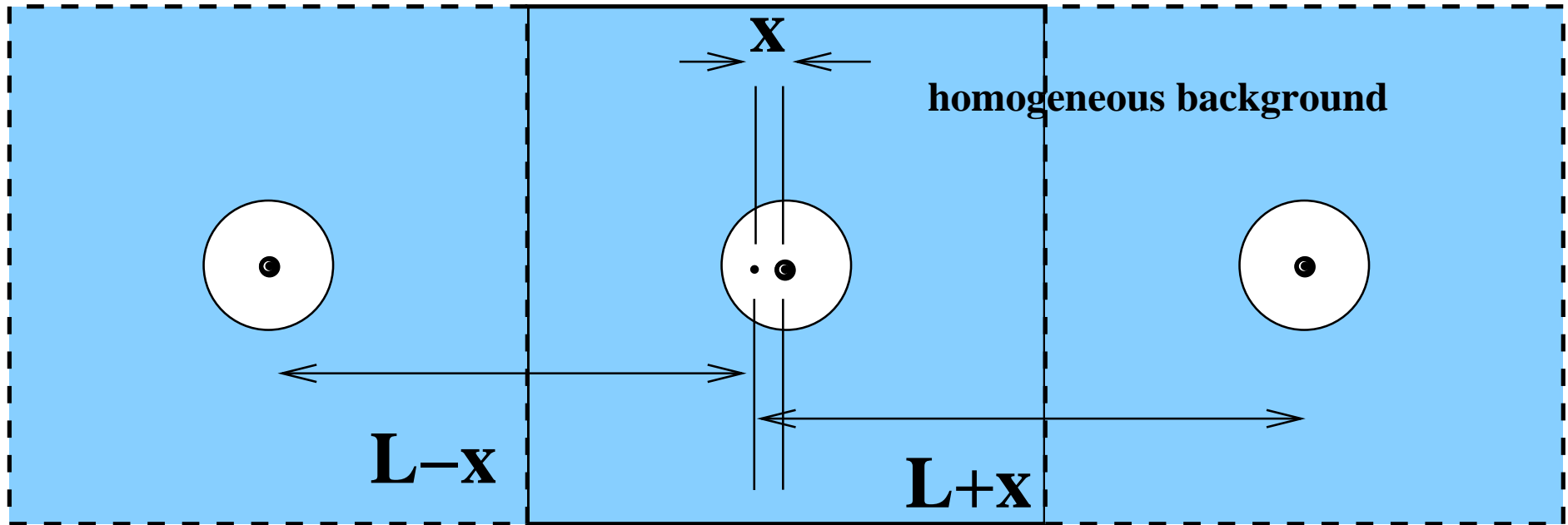


$$\ddot{x} \approx -G \frac{m}{x^2} + Gm \left[ \frac{1}{(L-x)^2} - \frac{1}{(L+x)^2} \right]$$

$$\approx -G \frac{m}{x^2} + \frac{4Gm}{L^2} \frac{x}{L}$$



# Cosmic topol: top. accel.

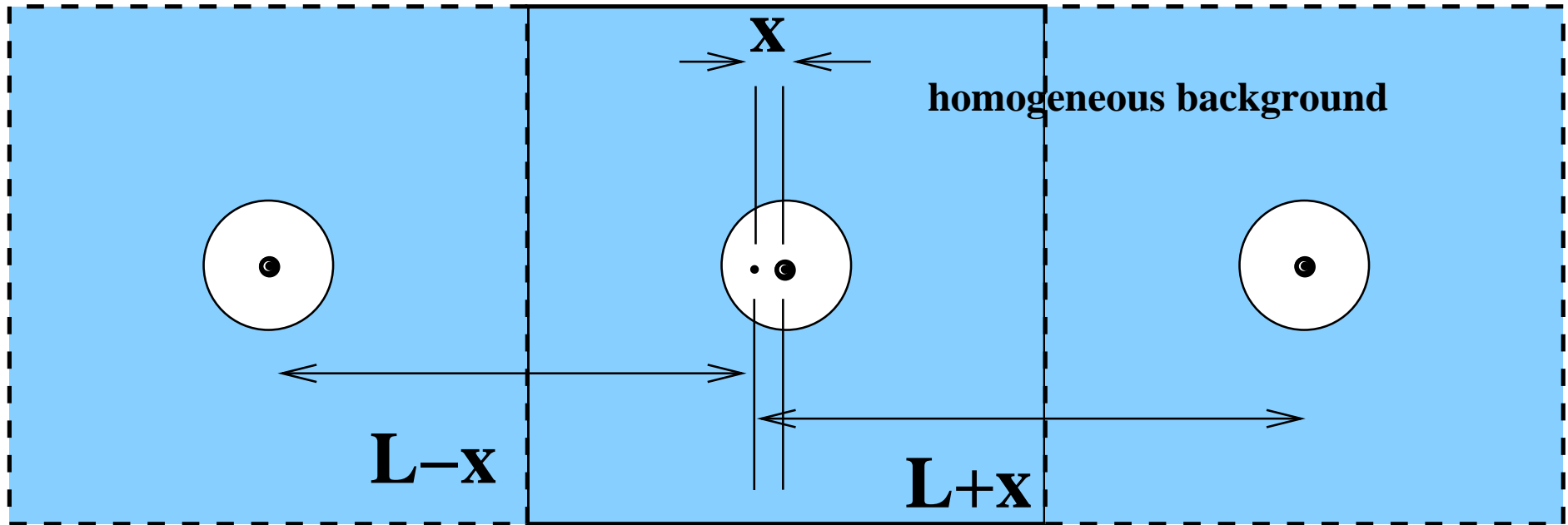


$$\ddot{x} \approx -G \frac{m}{x^2} + Gm \left[ \frac{1}{(L-x)^2} - \frac{1}{(L+x)^2} \right]$$

$$\approx -G \frac{m}{x^2} + \frac{4Gm}{L^2} \frac{x}{L}$$

$$\ddot{x}_{\text{resid}} \propto (x/L)^1 + \dots$$

# Cosmic topol: top. accel.



$$\ddot{x} \approx -G \frac{m}{x^2} + Gm \left[ \frac{1}{(L-x)^2} - \frac{1}{(L+x)^2} \right]$$

$$\approx -G \frac{m}{x^2} + \frac{4Gm}{L^2} \frac{x}{L}$$

$$\ddot{x}_{\text{resid}} \propto (x/L)^1 + \dots$$

topological acceleration— [arXiv:astro-ph/0602159](https://arxiv.org/abs/0602159)

# Cosmic topol: obs. strategies

empirical strategies: [arXiv:astro-ph/0010189](https://arxiv.org/abs/astro-ph/0010189)

# Cosmic topol: obs. strategies

empirical strategies: [arXiv:astro-ph/0010189](https://arxiv.org/abs/astro-ph/0010189)

A. multiple topological images:

# Cosmic topol: obs. strategies

empirical strategies: [arXiv:astro-ph/0010189](https://arxiv.org/abs/astro-ph/0010189)

A. multiple topological images:

A.i 3D (grav collapsed objects):

# Cosmic topol: obs. strategies

empirical strategies: [arXiv:astro-ph/0010189](https://arxiv.org/abs/astro-ph/0010189)

A. multiple topological images:

A.i 3D (grav collapsed objects):

A.i.1 local isometries—collect “type I pairs” or “local pairs”

# Cosmic topol: obs. strategies

empirical strategies: [arXiv:astro-ph/0010189](https://arxiv.org/abs/astro-ph/0010189)

A. multiple topological images:

A.i 3D (grav collapsed objects):

A.i.1 local isometries—collect “type I pairs” or “local pairs”

A.i.2 cosmic crystallography—collect “type II pairs” or  
“holonomy pairs”

# Cosmic topol: obs. strategies

empirical strategies: [arXiv:astro-ph/0010189](https://arxiv.org/abs/astro-ph/0010189)

A. multiple topological images:

A.i 3D (grav collapsed objects):

A.i.1 local isometries—collect “type I pairs” or “local pairs”

A.i.2 cosmic crystallography—collect “type II pairs” or  
“holonomy pairs”

A.i.3 characteristics of individual objects



# Cosmic topol: obs. strategies

empirical strategies: [arXiv:astro-ph/0010189](https://arxiv.org/abs/astro-ph/0010189)

A. multiple topological images:

A.i 3D (grav collapsed objects):

A.i.1 local isometries—collect “type I pairs” or “local pairs”

A.i.2 cosmic crystallography—collect “type II pairs” or  
“holonomy pairs”

A.i.3 characteristics of individual objects—(PESEL method)

# Cosmic topol: obs. strategies

A. multiple topological images:

# Cosmic topol: obs. strategies

A. multiple topological images:

A.ii 2D (e.g. microwave background = CMB):

# Cosmic topol: obs. strategies

A. multiple topological images:

A.ii 2D (e.g. microwave background = CMB):

A.ii.1 identified circles principle

# Cosmic topol: obs. strategies

A. multiple topological images:

A.ii 2D (e.g. microwave background = CMB):

A.ii.1 identified circles principle

A.ii.2 matched discs corollary

# Cosmic topol: obs. strategies

A. multiple topological images:

A.ii 2D (e.g. microwave background = CMB):

A.ii.1 identified circles principle

A.ii.2 matched discs corollary

A.ii.3 patterns of spots

# Cosmic topol: obs. strategies

A. multiple topological images:

A.ii 2D (e.g. microwave background = CMB):

A.ii.1 identified circles principle

A.ii.2 matched discs corollary

A.ii.3 patterns of spots

A.ii.4 perturbation statistics assumptions

# Cosmic topol: obs. strategies

A. multiple topological images:

A.ii 2D (e.g. microwave background = CMB):

A.ii.1 identified circles principle

A.ii.2 matched discs corollary

A.ii.3 patterns of spots

A.ii.4 perturbation statistics assumptions

B. other:



# Cosmic topol: obs. strategies

A. multiple topological images:

A.ii 2D (e.g. microwave background = CMB):

A.ii.1 identified circles principle

A.ii.2 matched discs corollary

A.ii.3 patterns of spots

A.ii.4 perturbation statistics assumptions

B. other:

B.i cosmic strings

# Cosmic topol: obs. strategies

A. multiple topological images:

A.ii 2D (e.g. microwave background = CMB):

A.ii.1 identified circles principle

A.ii.2 matched discs corollary

A.ii.3 patterns of spots

A.ii.4 perturbation statistics assumptions

B. other:

B.i cosmic strings

B.ii topological acceleration

# 3D strategies

- TODO...

# 2D strategies

■ TODO...